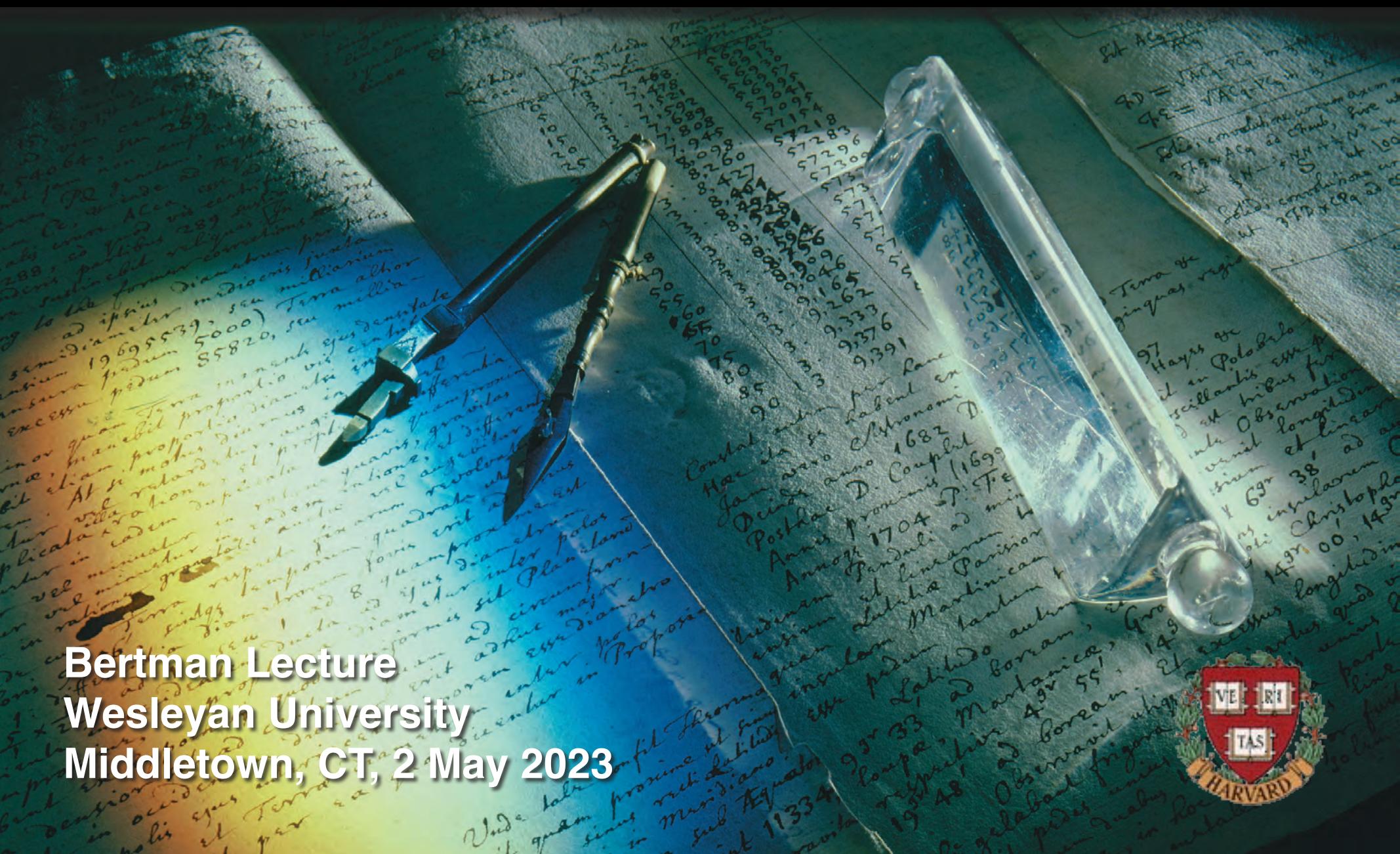


The surprising world where optical properties approach zero



Bertman Lecture
 Wesleyan University
 Middletown, CT, 2 May 2023

The surprising world where optical properties approach zero



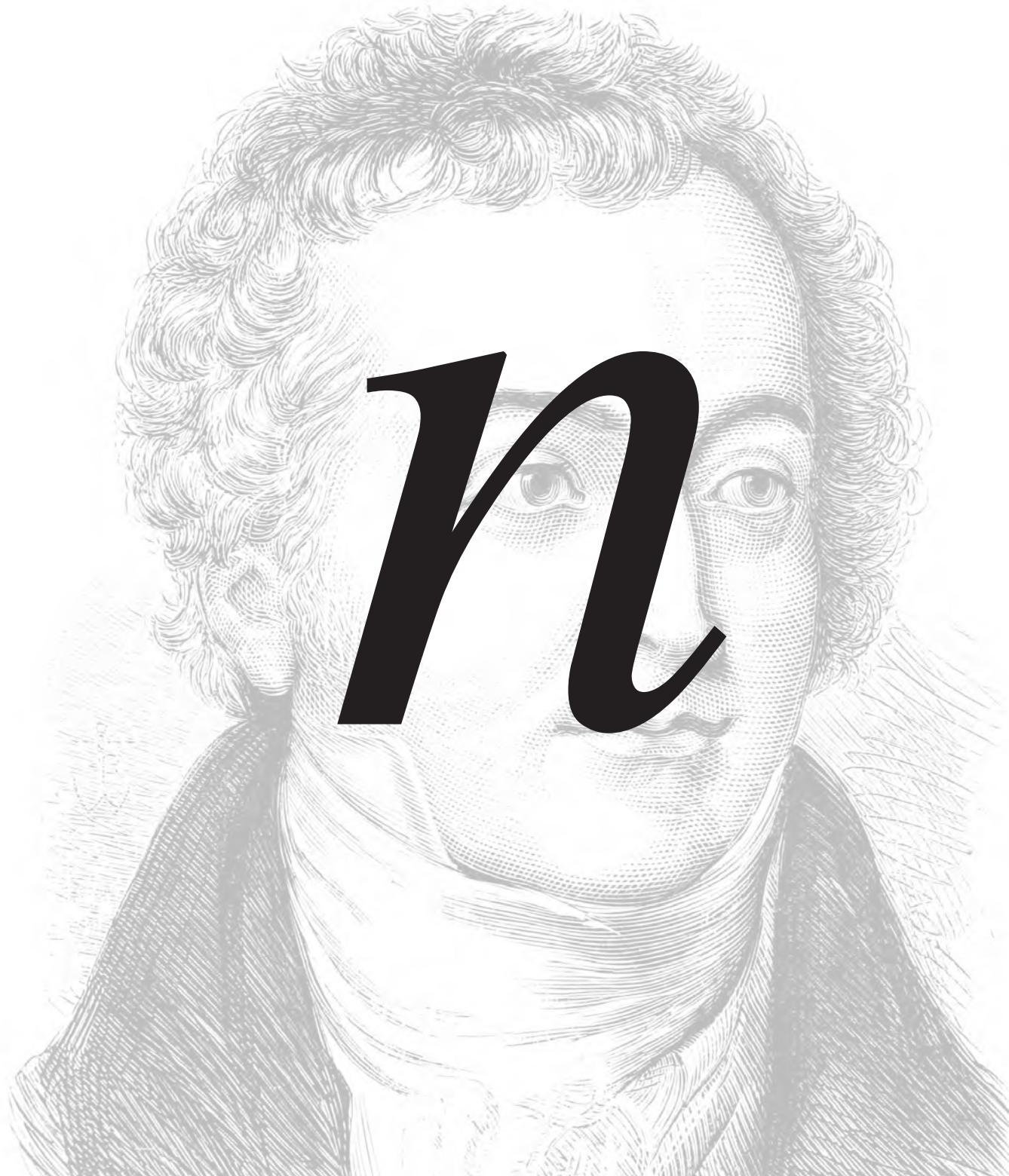
@eric_mazur

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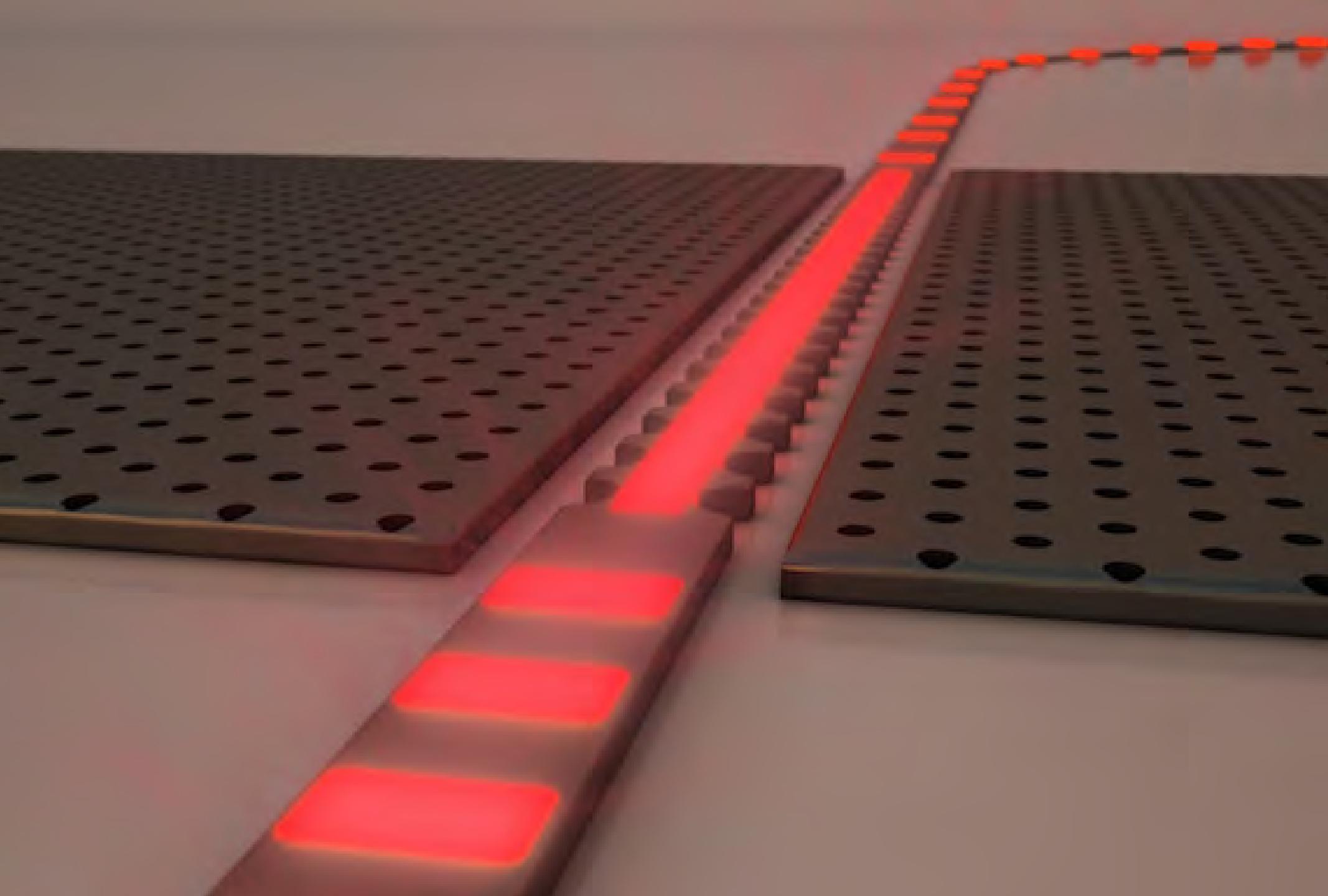
n



n



1 index



1 index

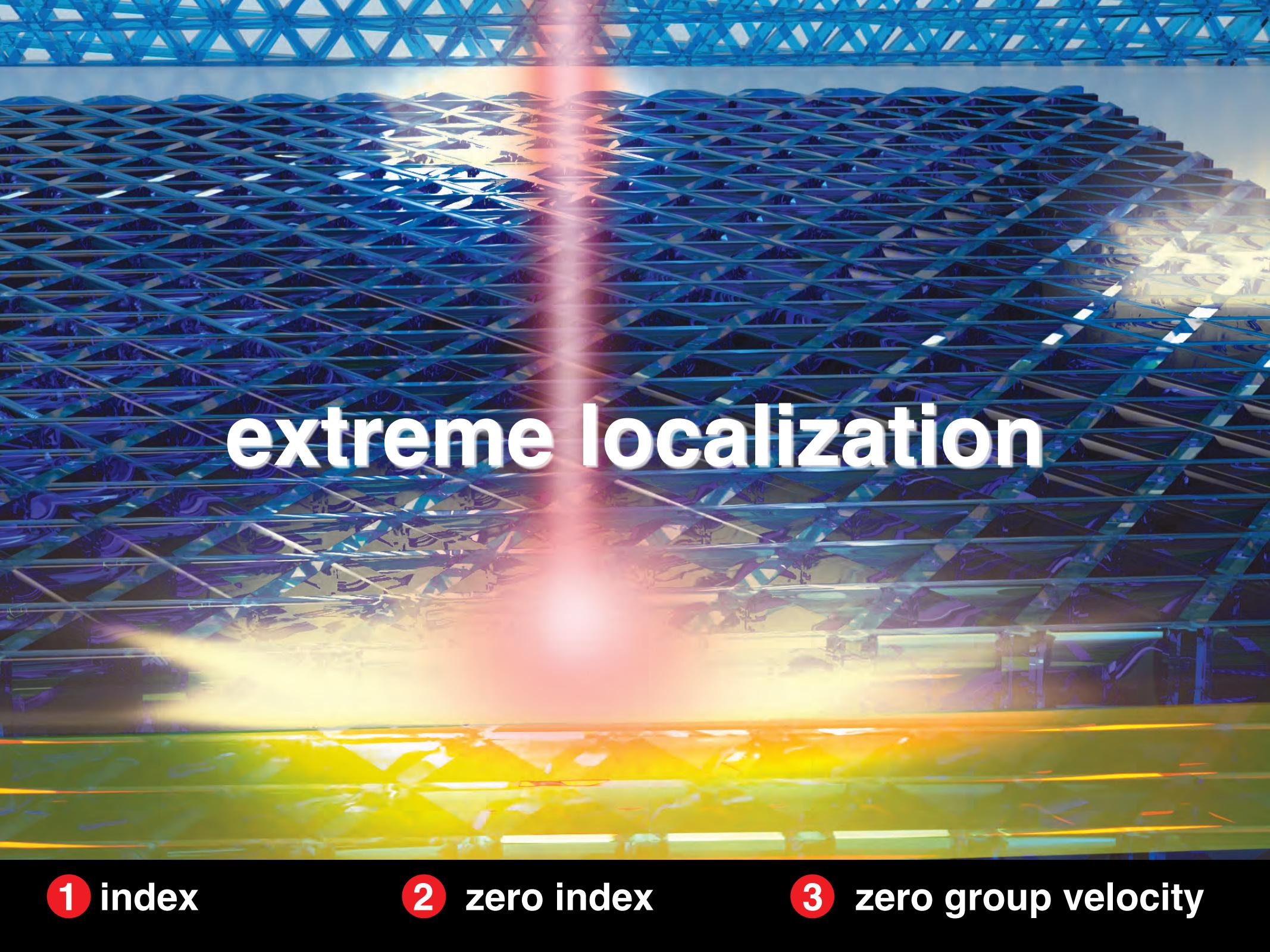
2 zero index

A photograph of a red laser beam passing through a metal grating. The beam is initially narrow and parallel, but as it moves along the grating, it spreads out into a series of bright, horizontal bands. The background is dark, and the laser beam is the primary light source.

extreme spreading

1 index

2 zero index

The background features a complex, glowing blue lattice structure composed of many thin, intersecting rods. A bright, vertical beam of light passes through the center of the lattice, transitioning from red at the top to yellow at the bottom. The overall effect is a high-energy, futuristic visualization.

extreme localization

1 index

2 zero index

3 zero group velocity

Propagation of EM wave

Propagation of EM wave

governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu\epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

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In dispersive media $n = n(\omega)$.

Index of refraction

$$n = \sqrt{\epsilon\mu}$$

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So $n(\omega)$ determined by response of material to external fields

Index of refraction

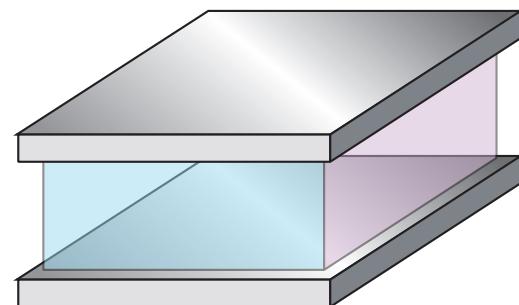
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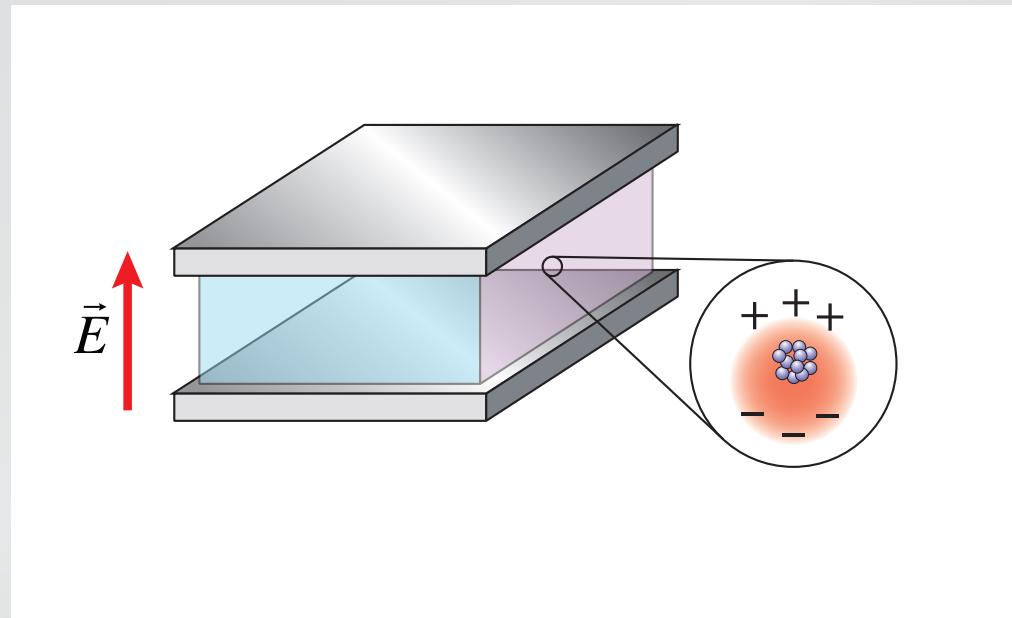
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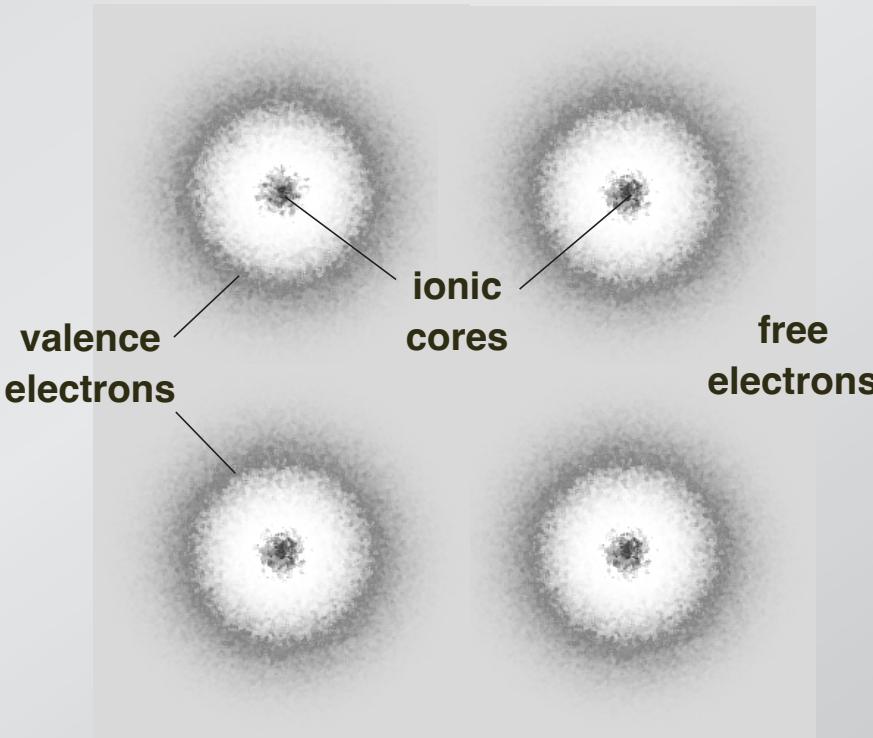


$\epsilon(\omega)$ measure of attenuation of electric field

Index of refraction

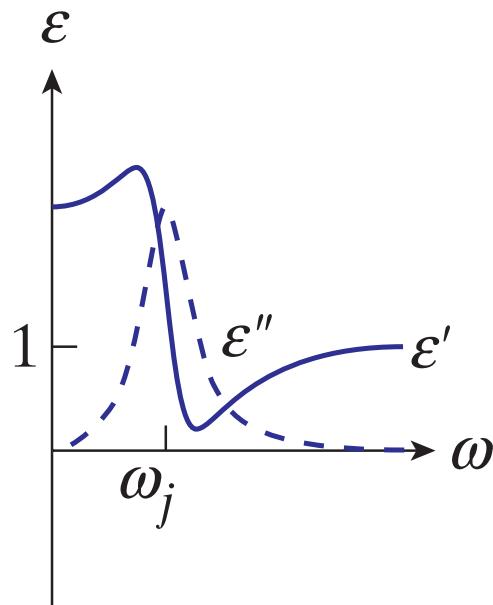
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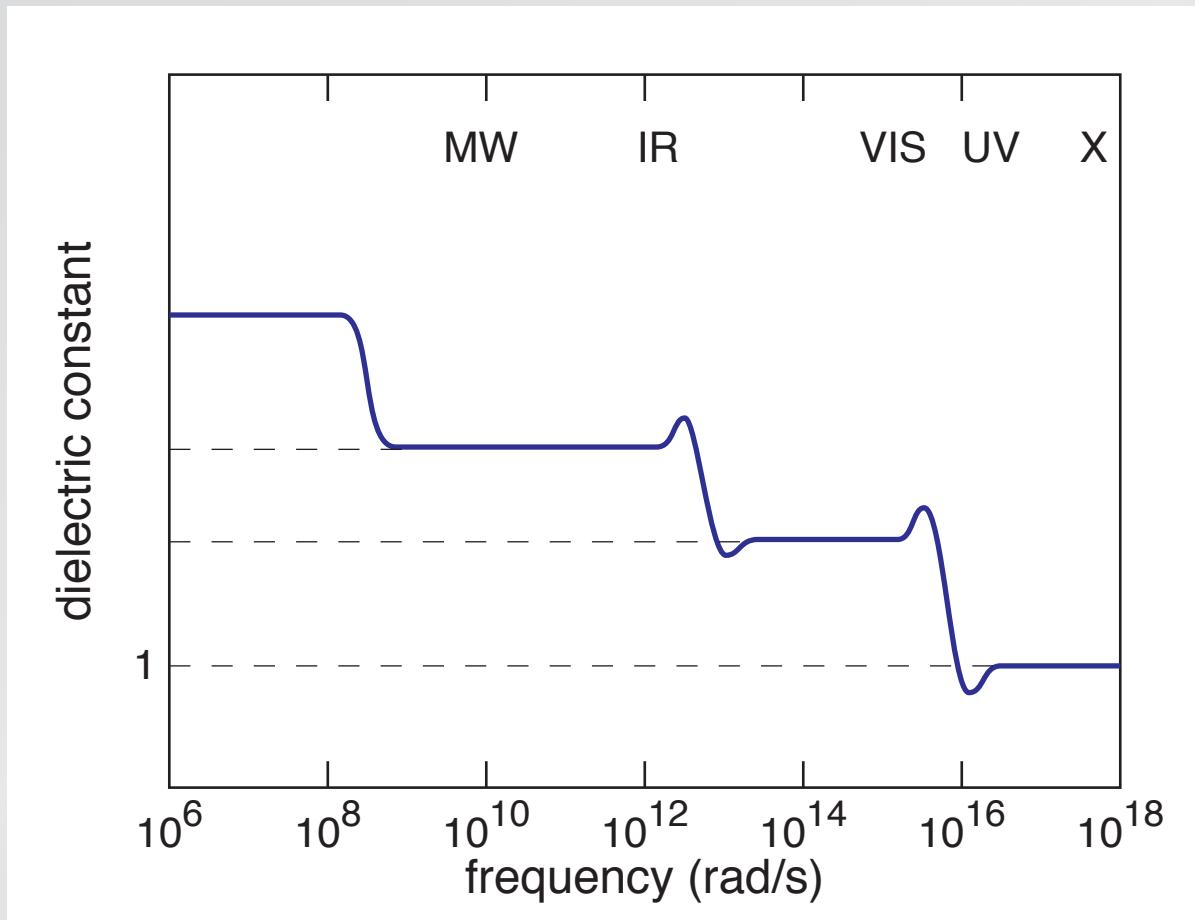


Dielectric constant

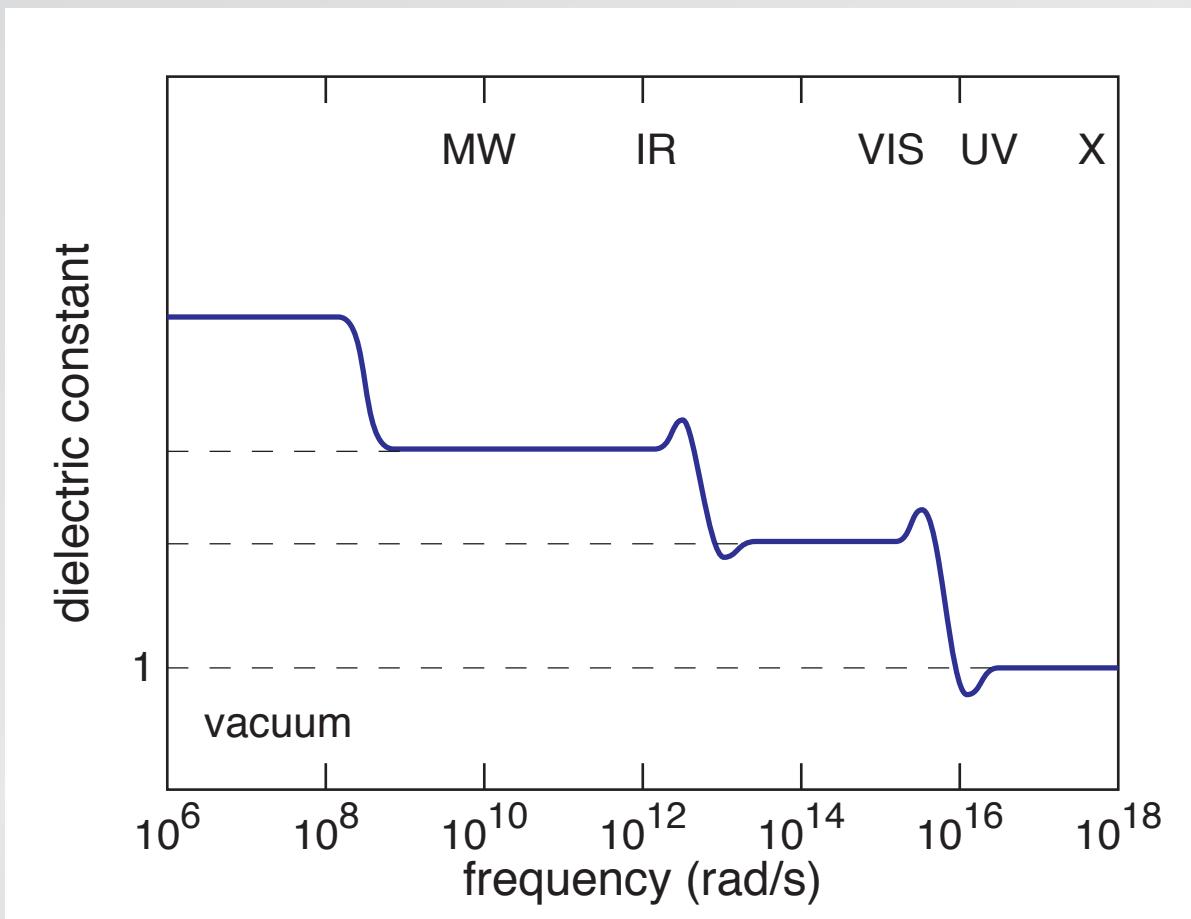
Lorentz oscillator



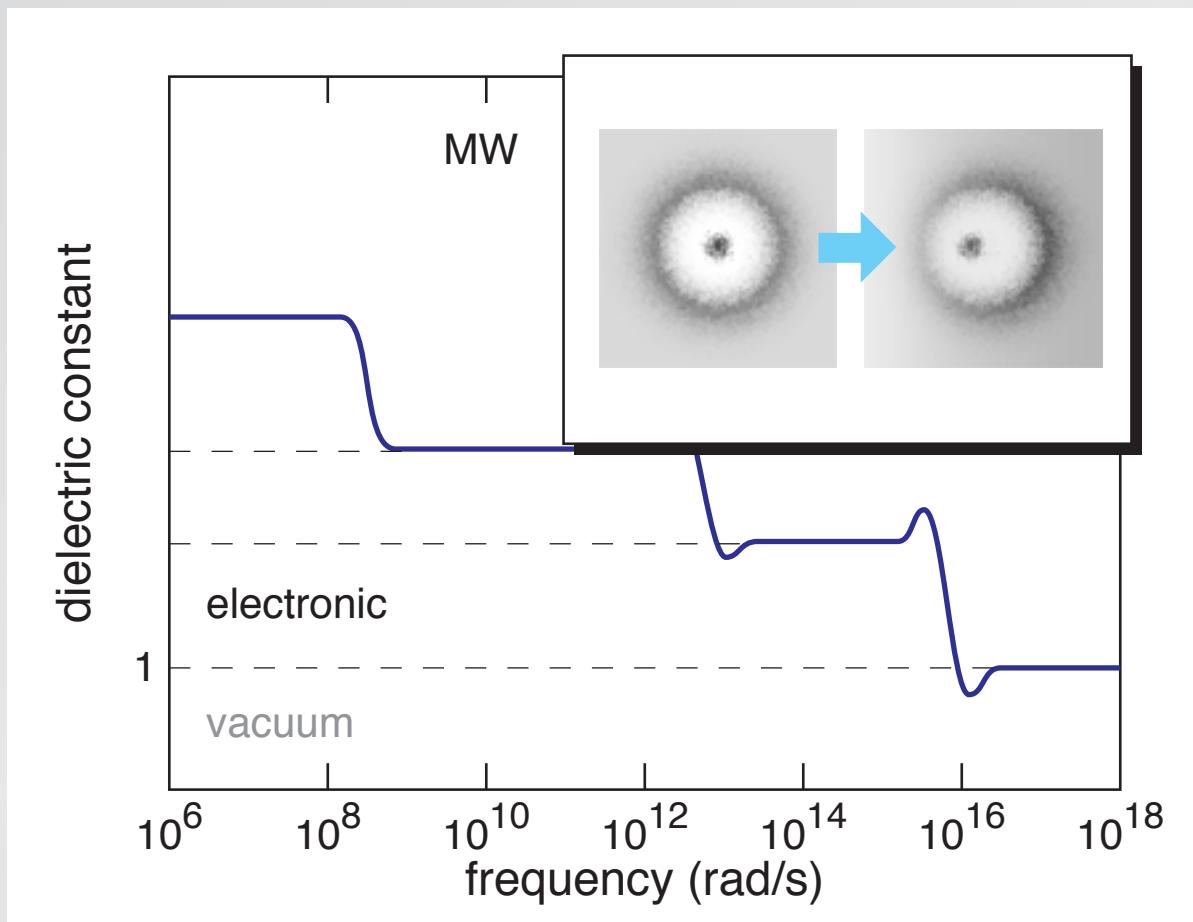
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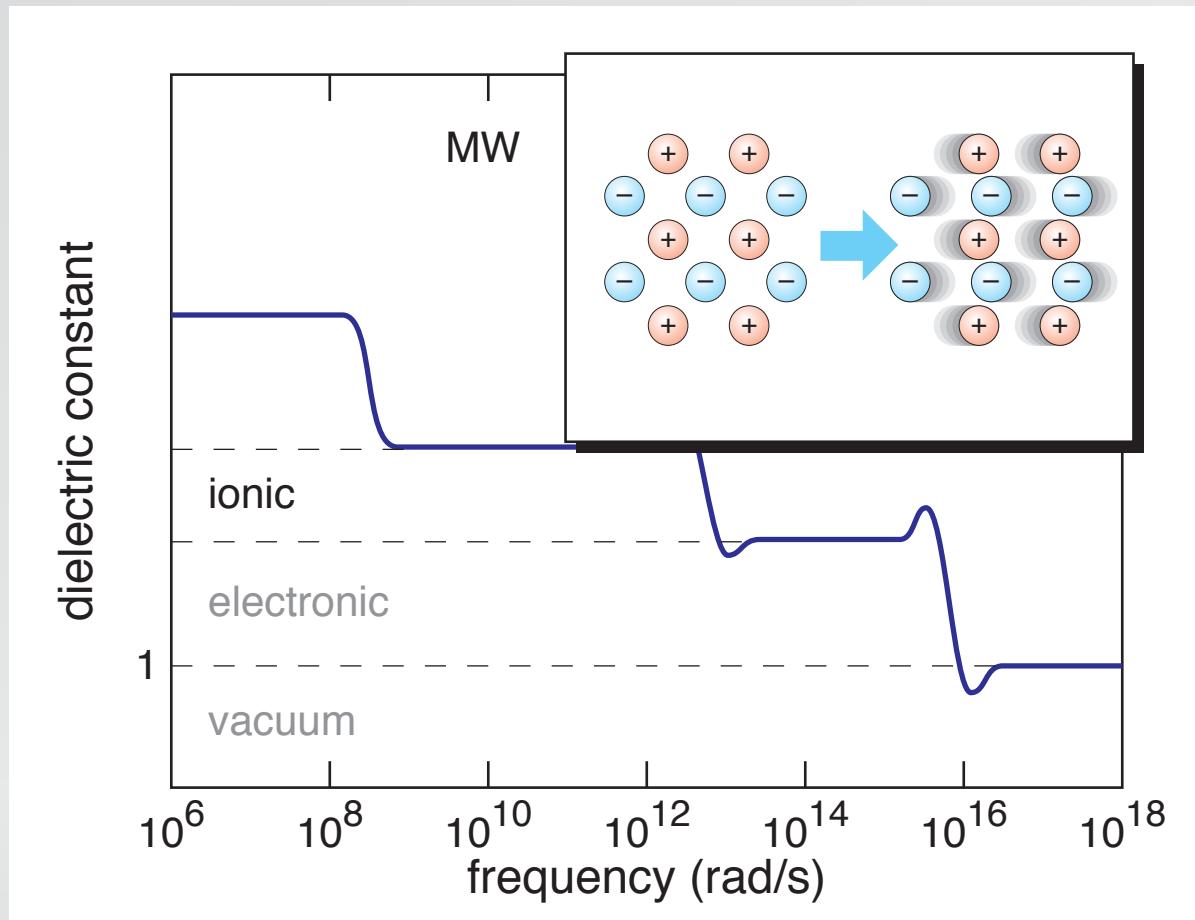
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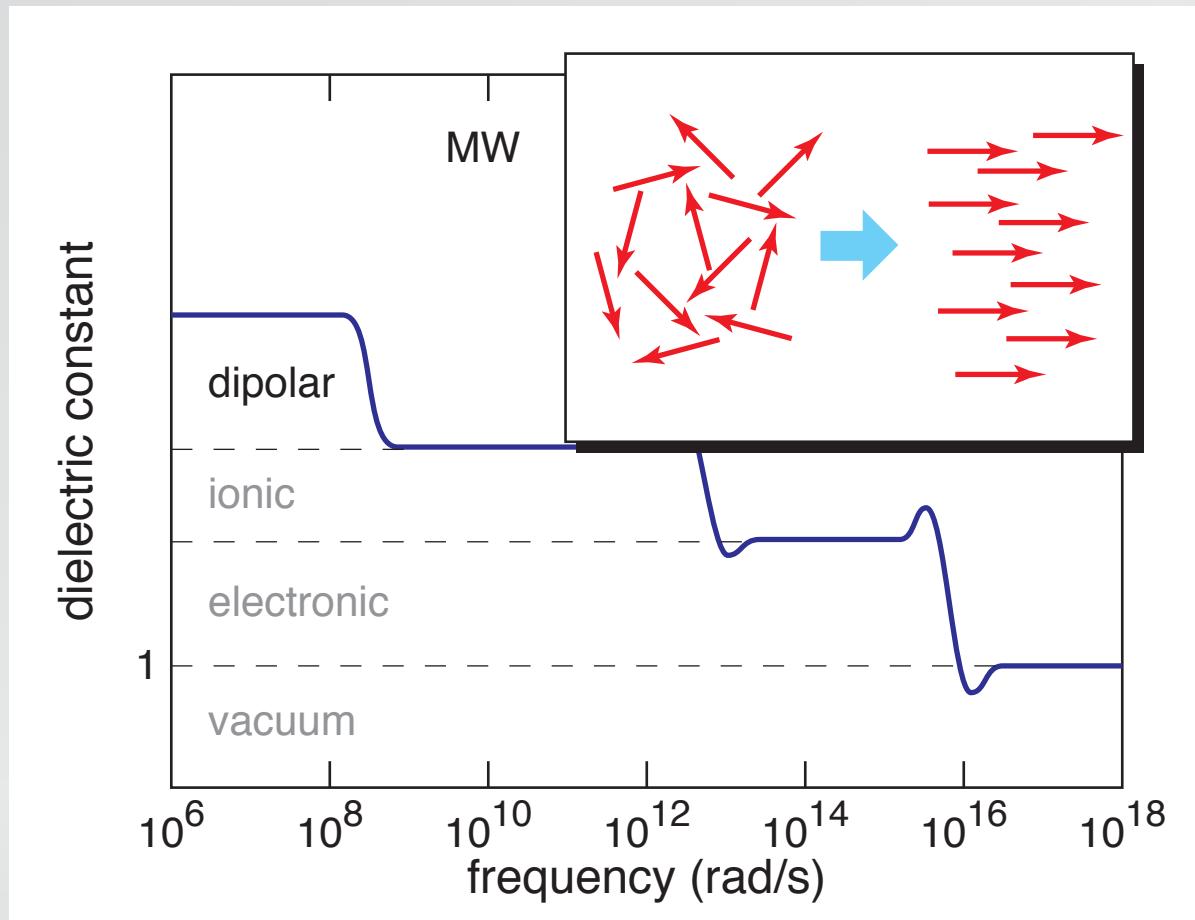
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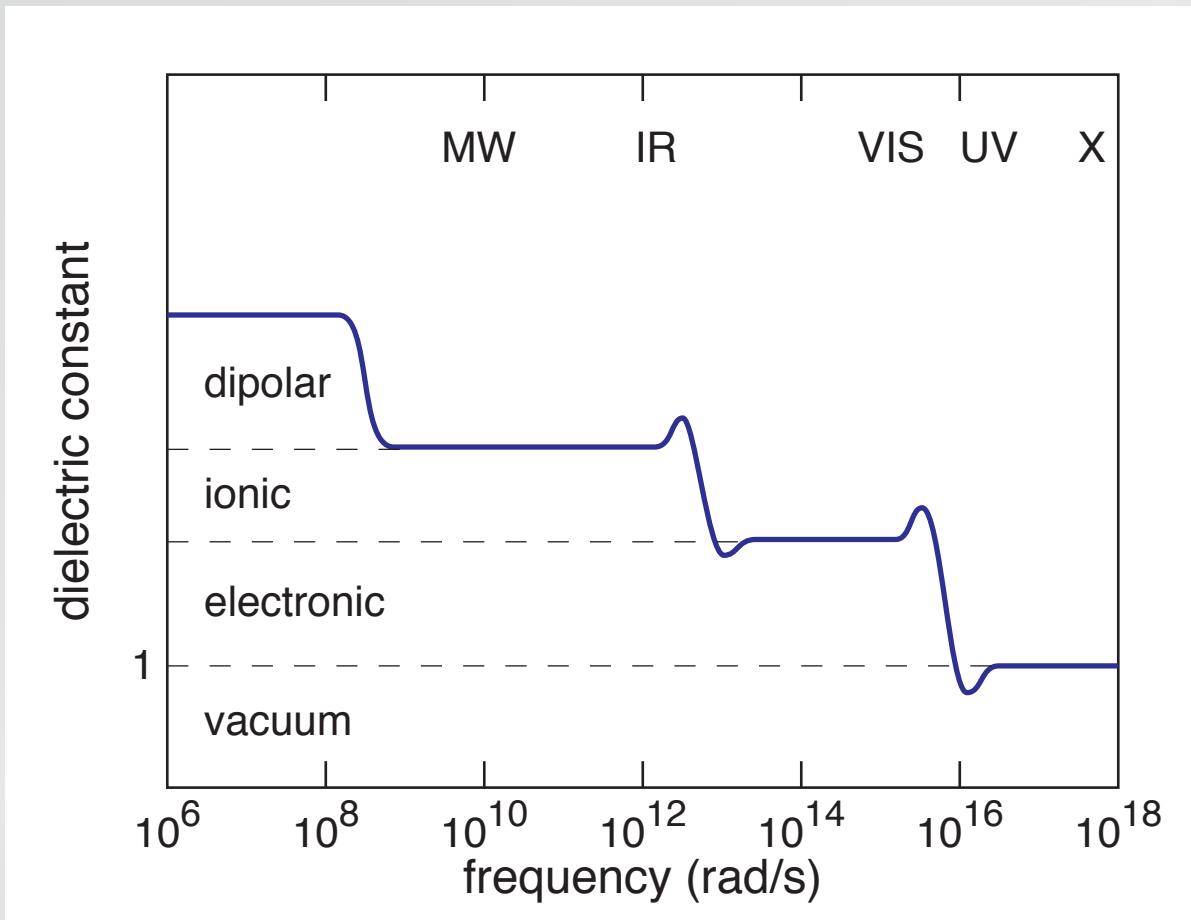
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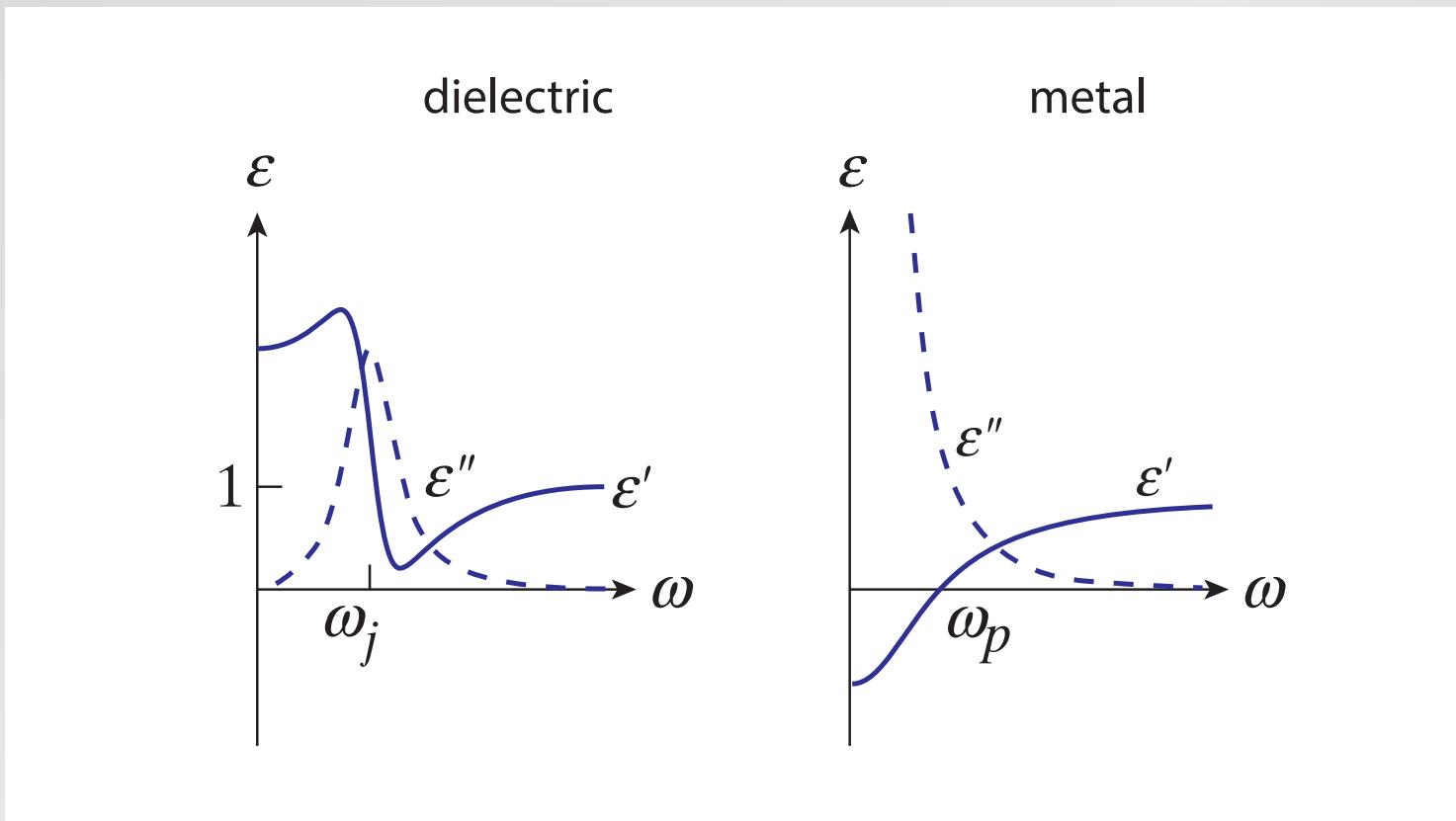
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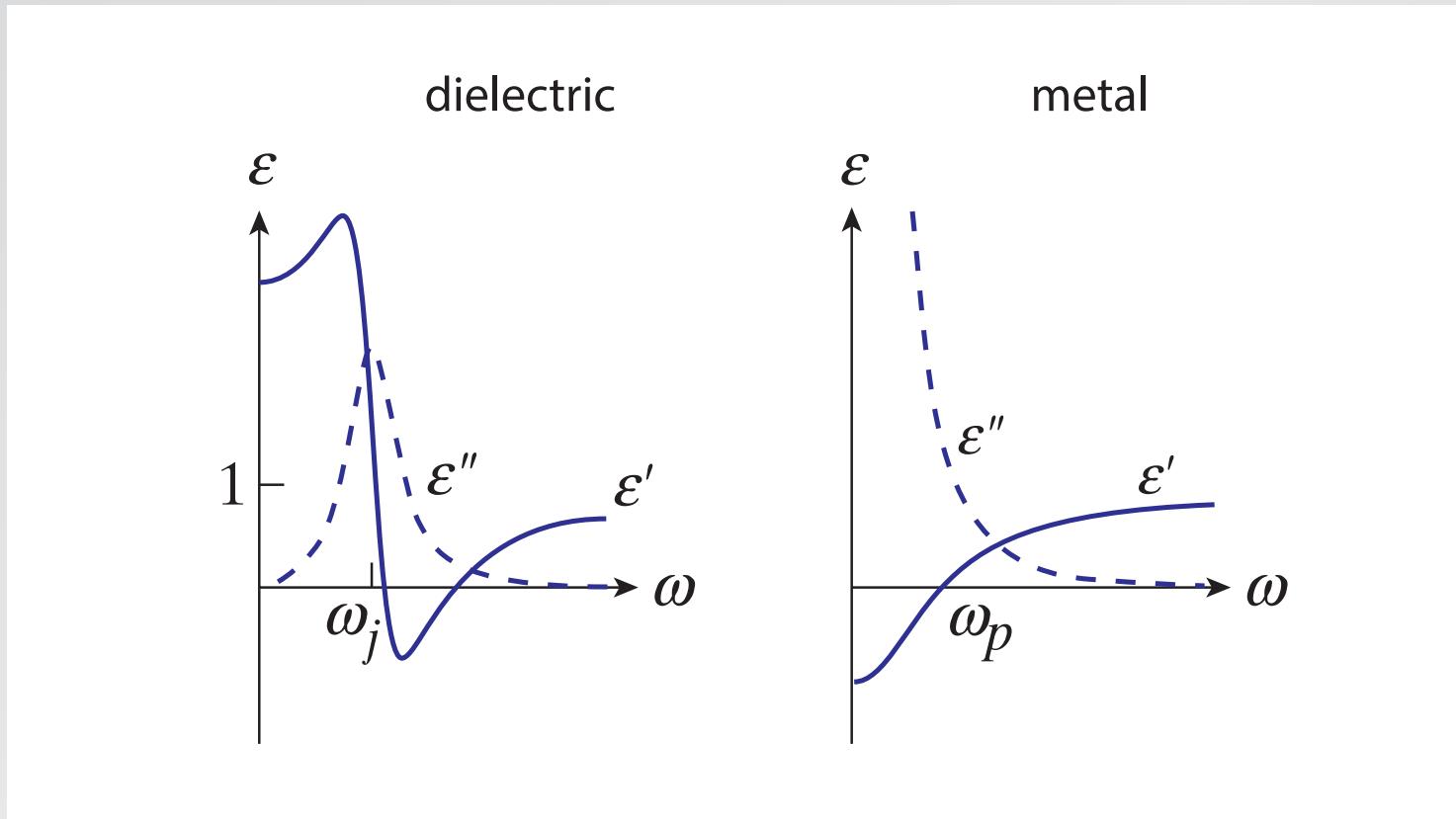


Lorentz and Drude models



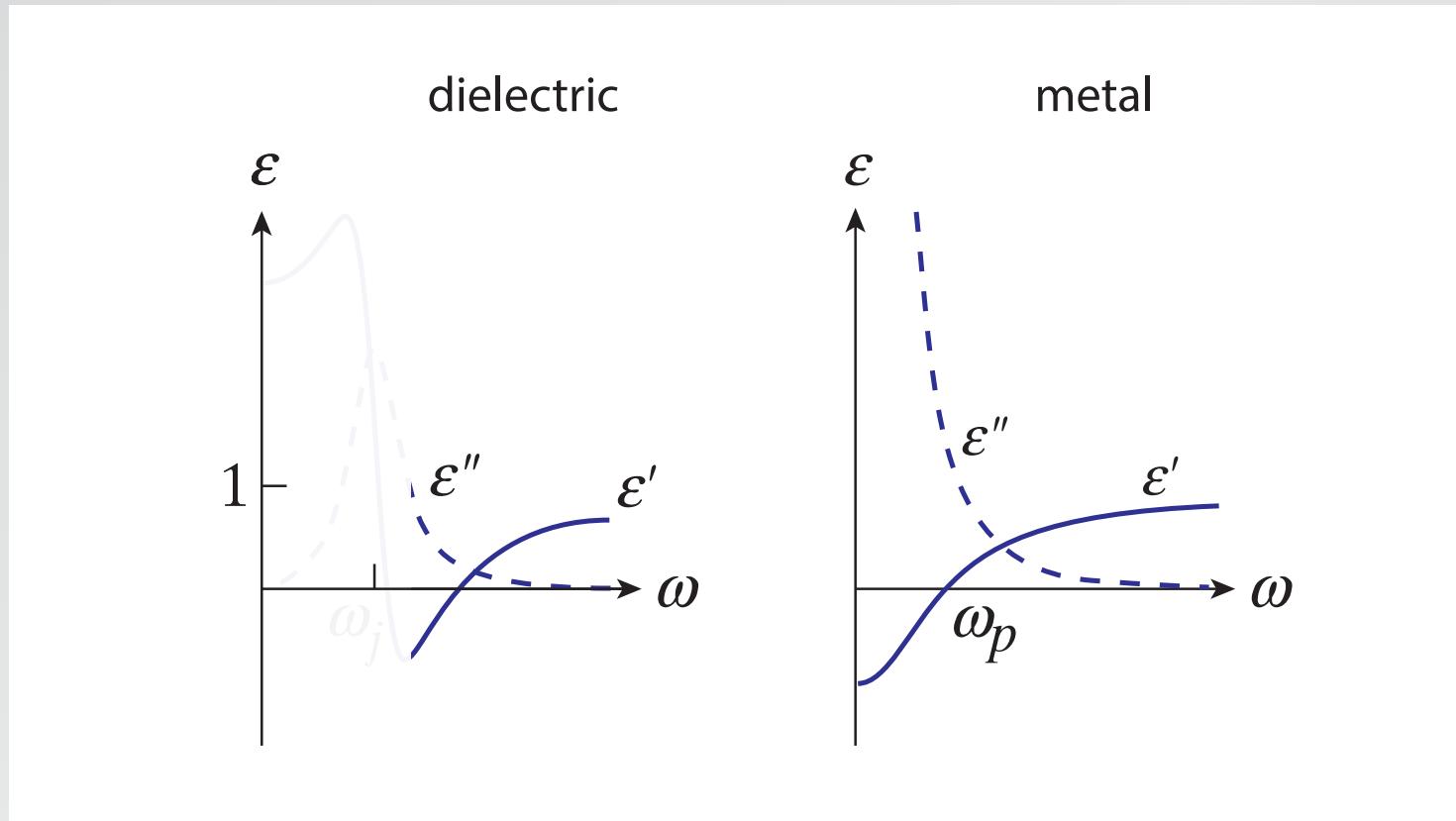
Lorentz and Drude models

for a strong (dielectric) resonance ε can become negative



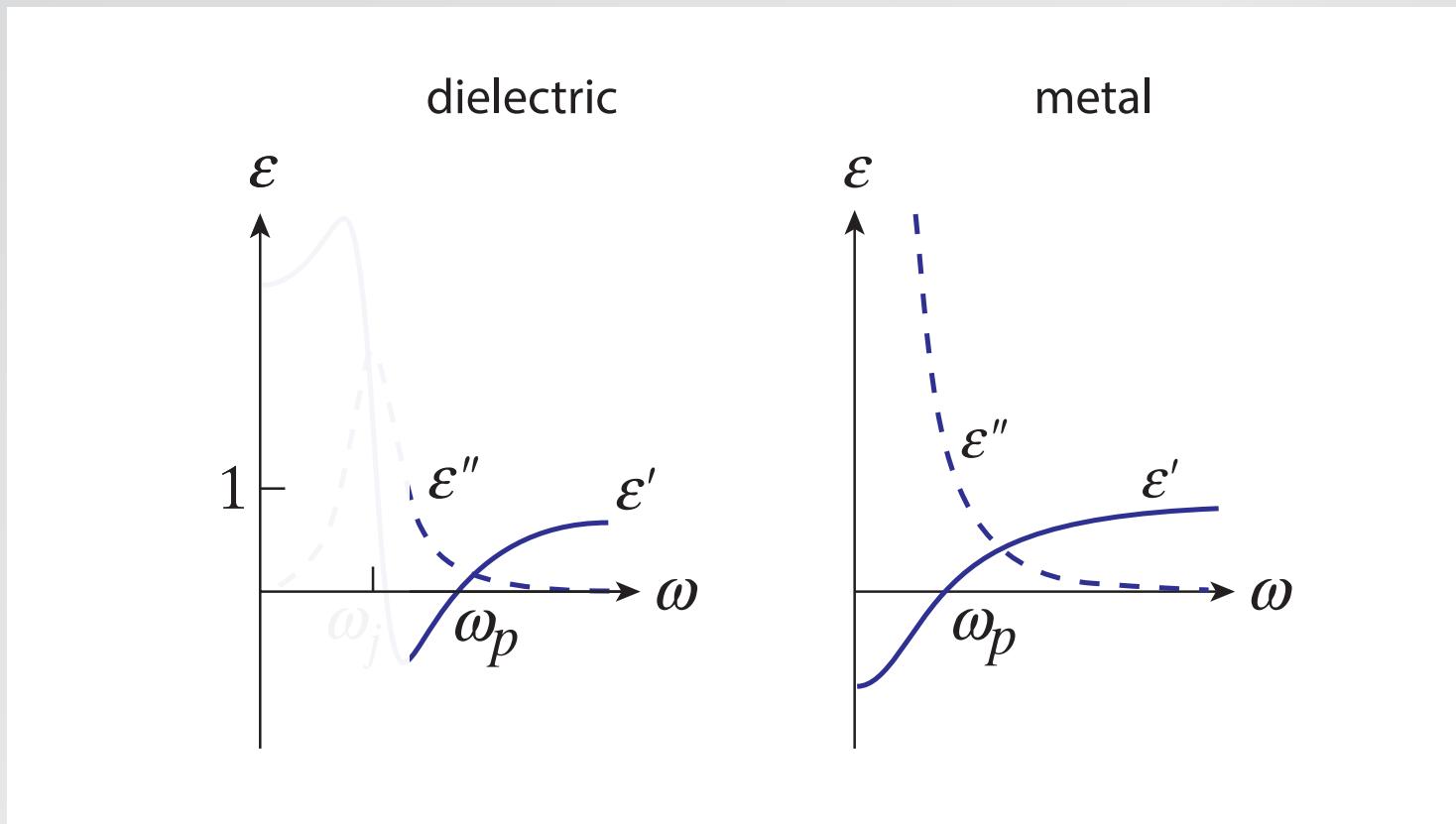
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valence electrons in dielectric then behave like a plasma



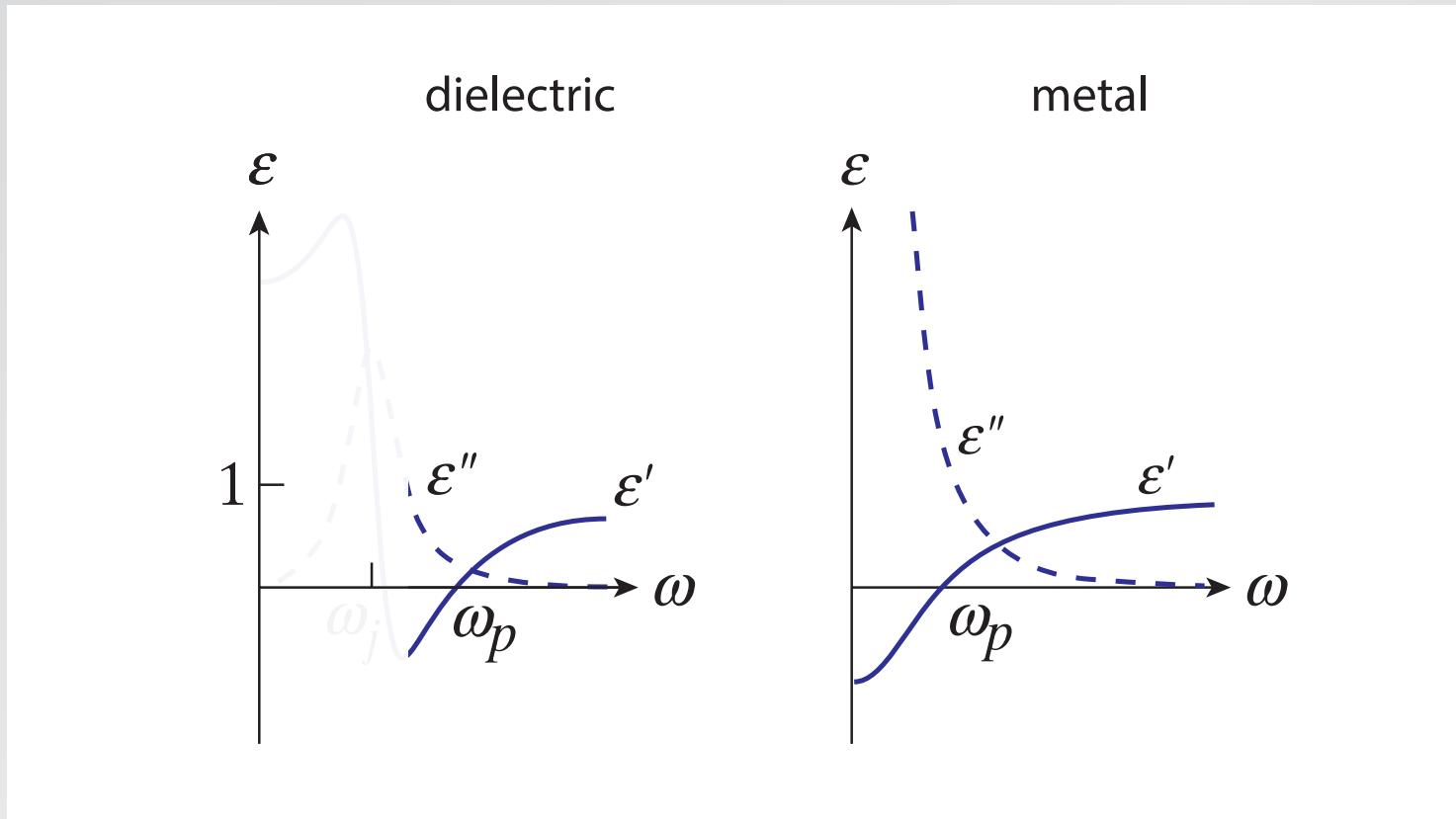
Lorentz and Drude models

with plasma frequency above the resonance



Lorentz and Drude models

(and far below the UV region)



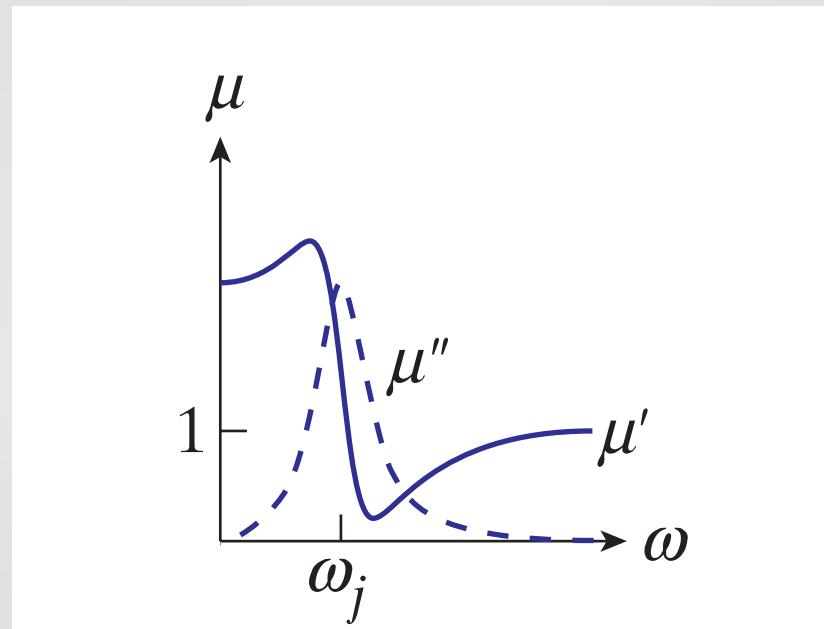
Index also determined by magnetic response

$$n = \sqrt{\epsilon\mu}$$

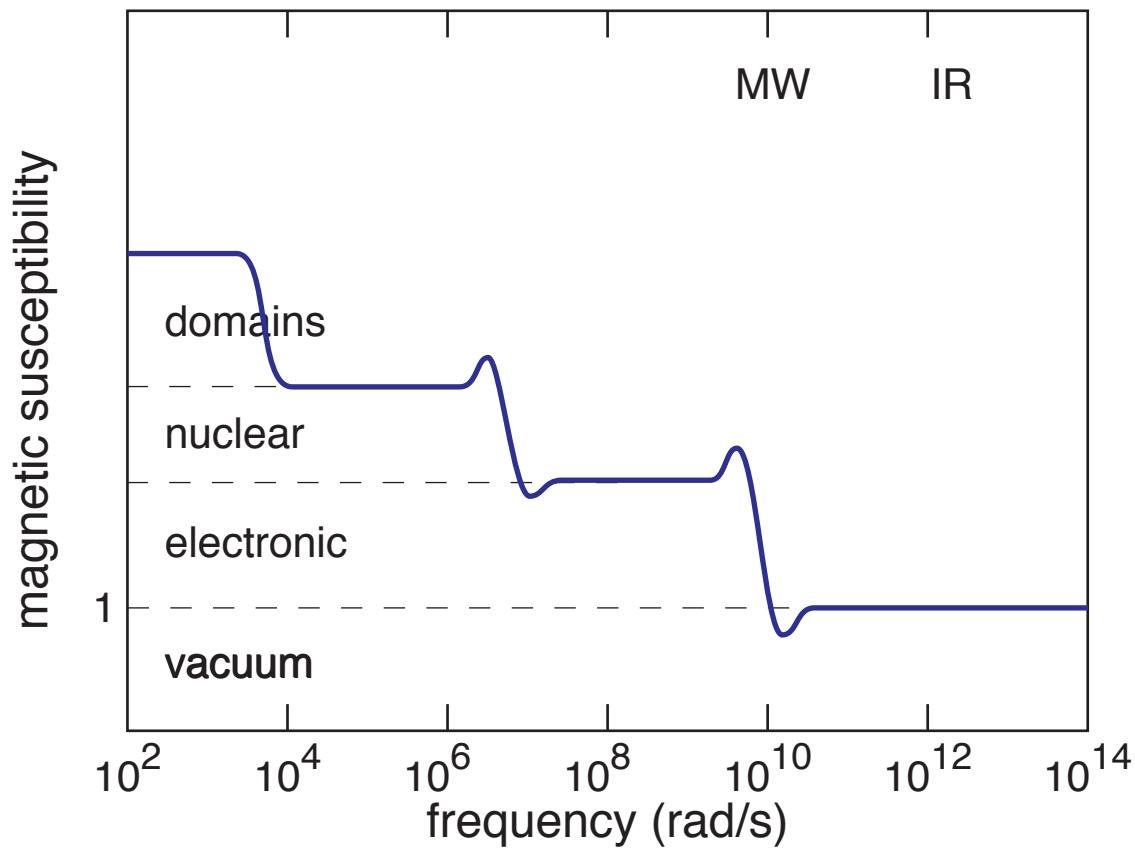
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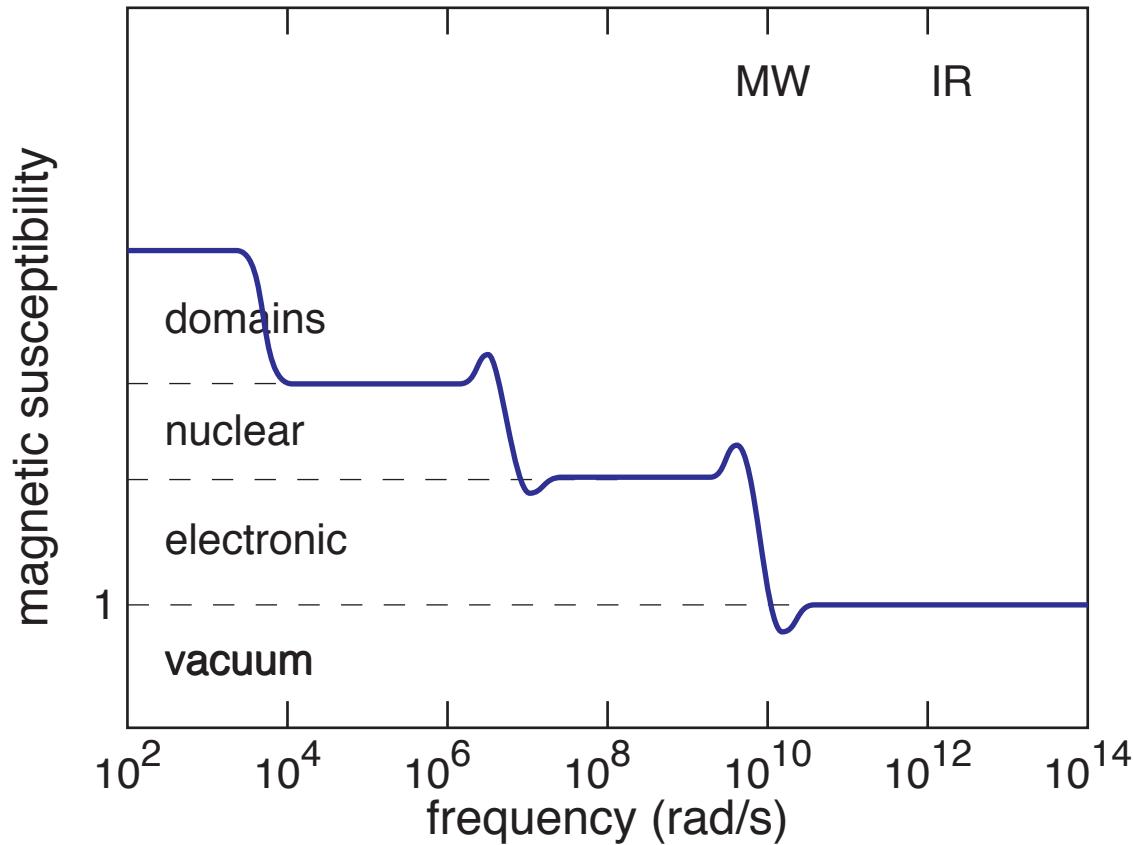


Magnetic response



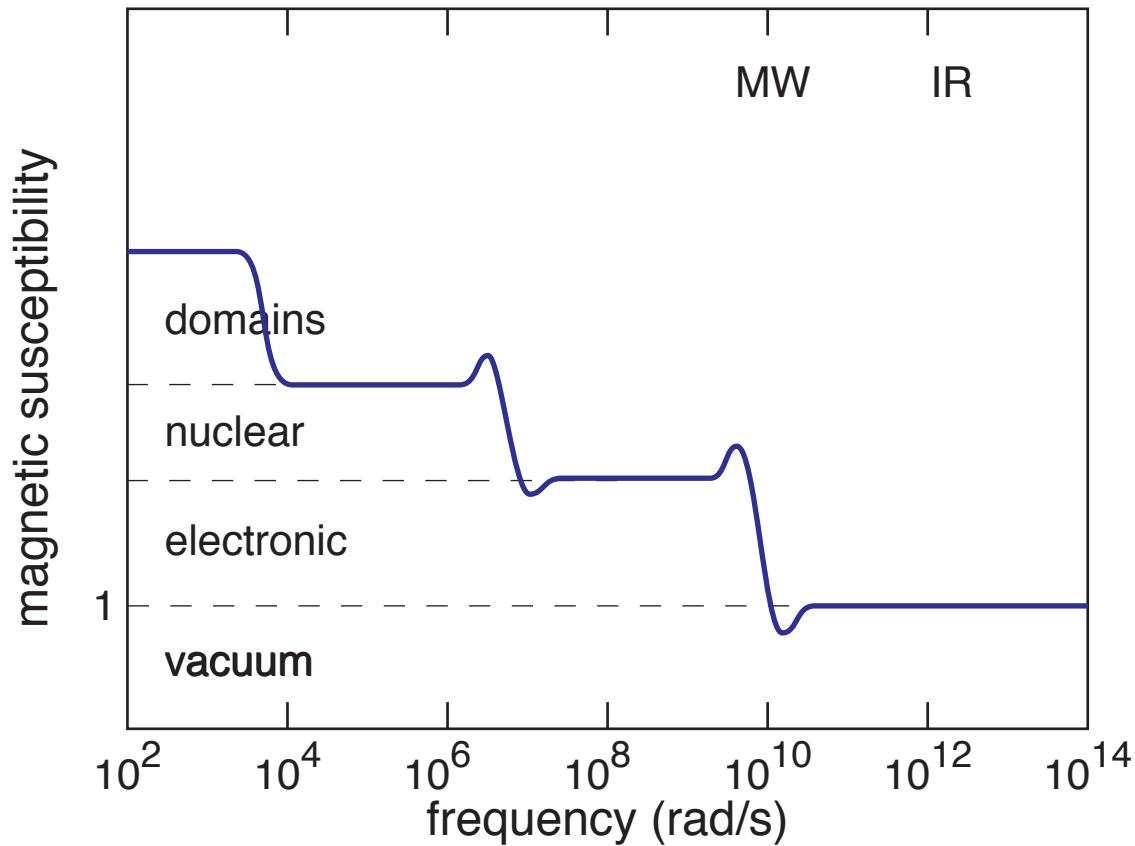
Magnetic response

but magnetic resonances occur below optical frequencies



Magnetic response

so, in optical regime, $\mu \approx 1$



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Both ϵ and μ are complex and their real parts can be negative.

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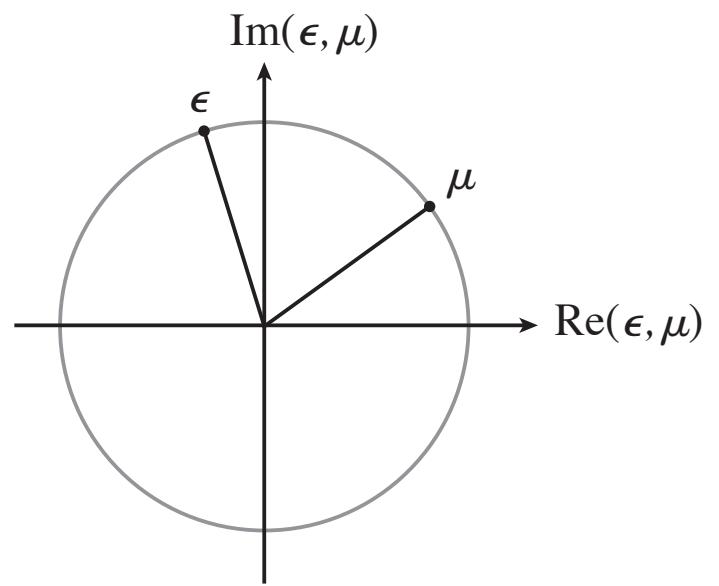
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Write complex quantities as

$$\varepsilon = |\varepsilon| e^{i\theta} \quad \mu = |\mu| e^{i\phi}$$

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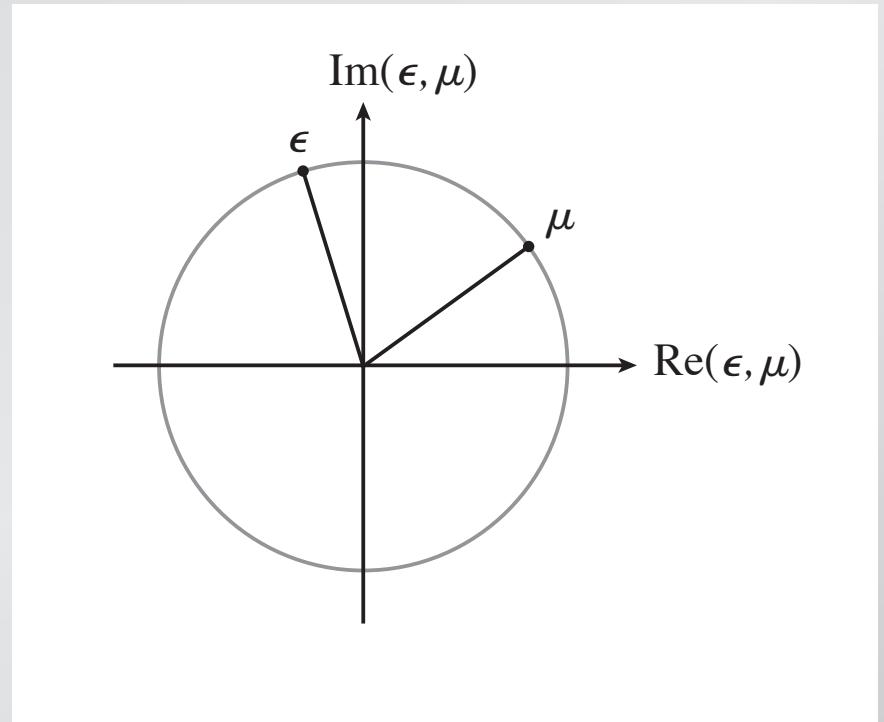


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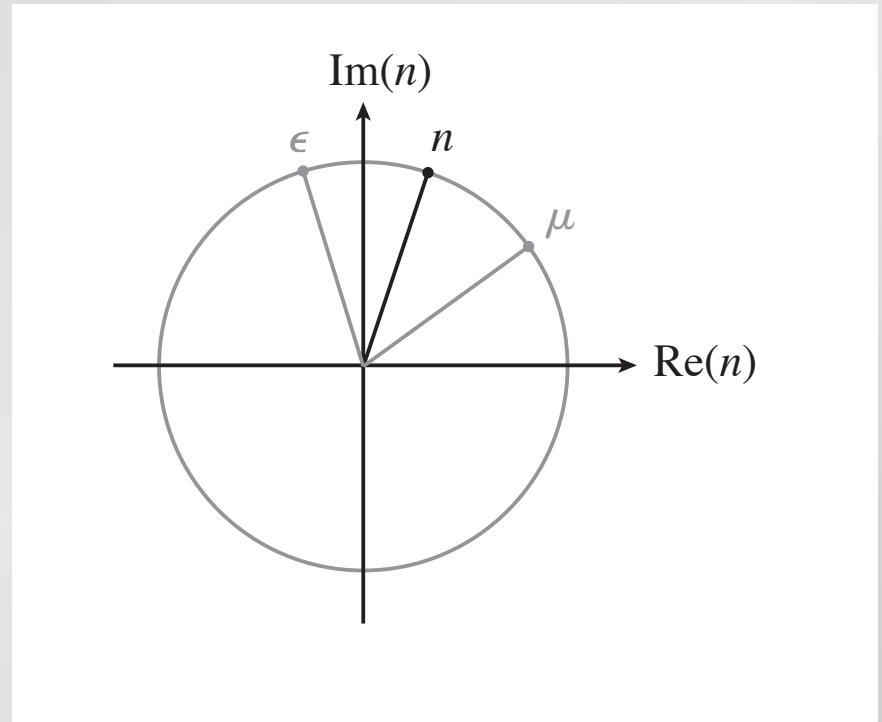


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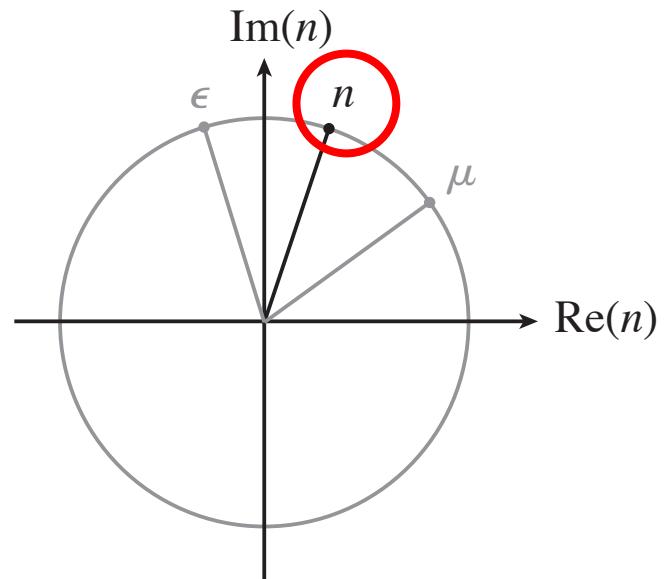
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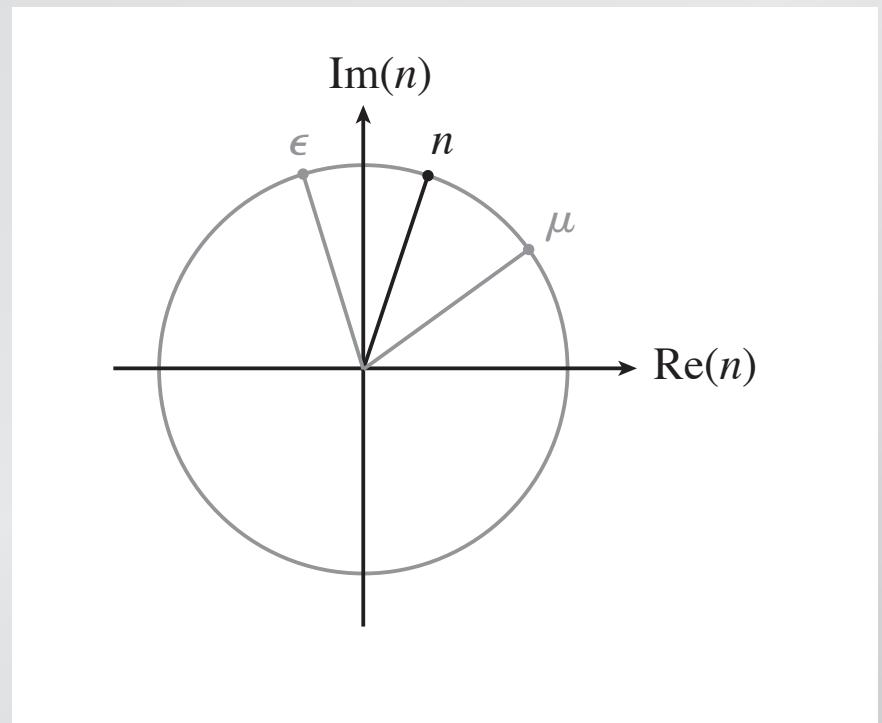
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Q: Is this only possible value?

1. yes
2. no, there's one more
3. there are many more
4. it depends



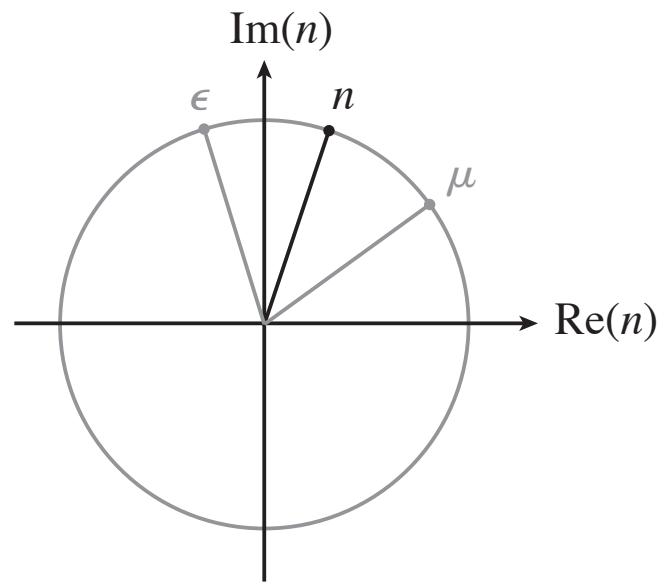
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Can add 2π to exponent

$$e^{+i(\theta+\phi)} = e^{+i[\theta+\phi+2\pi]}$$



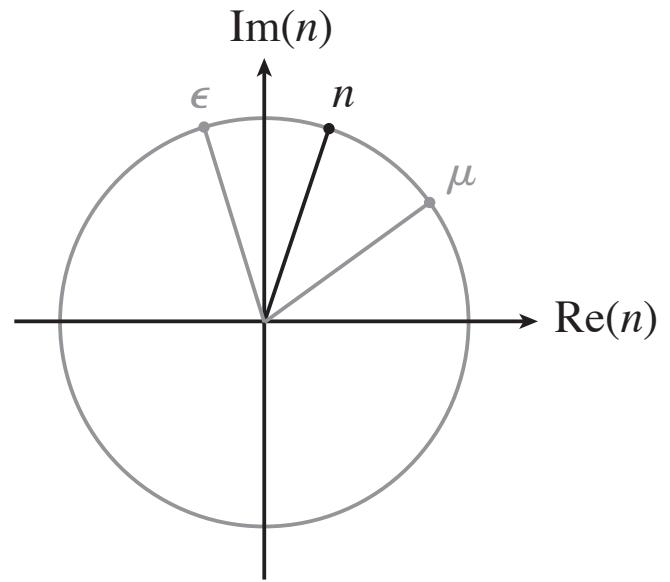
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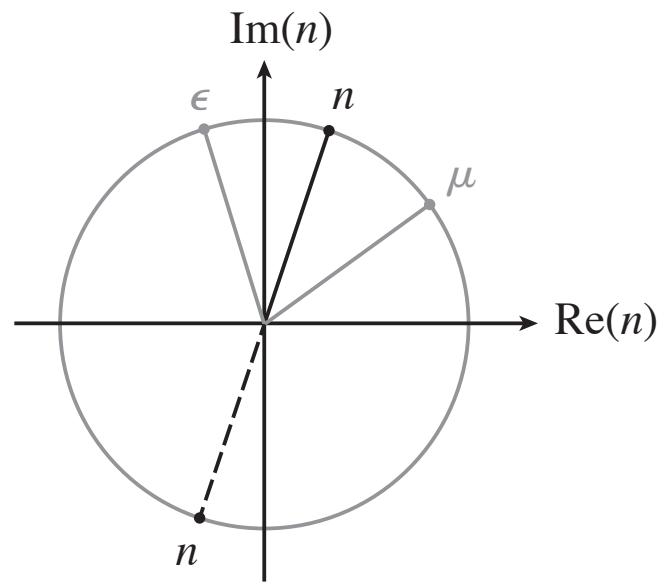
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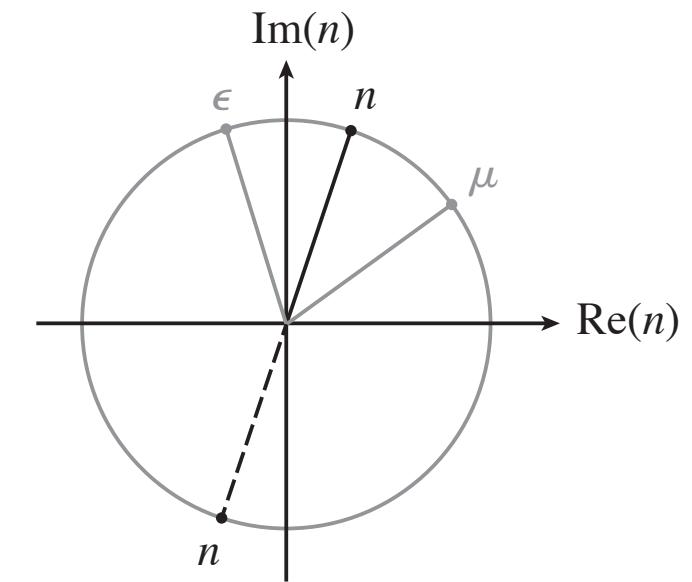
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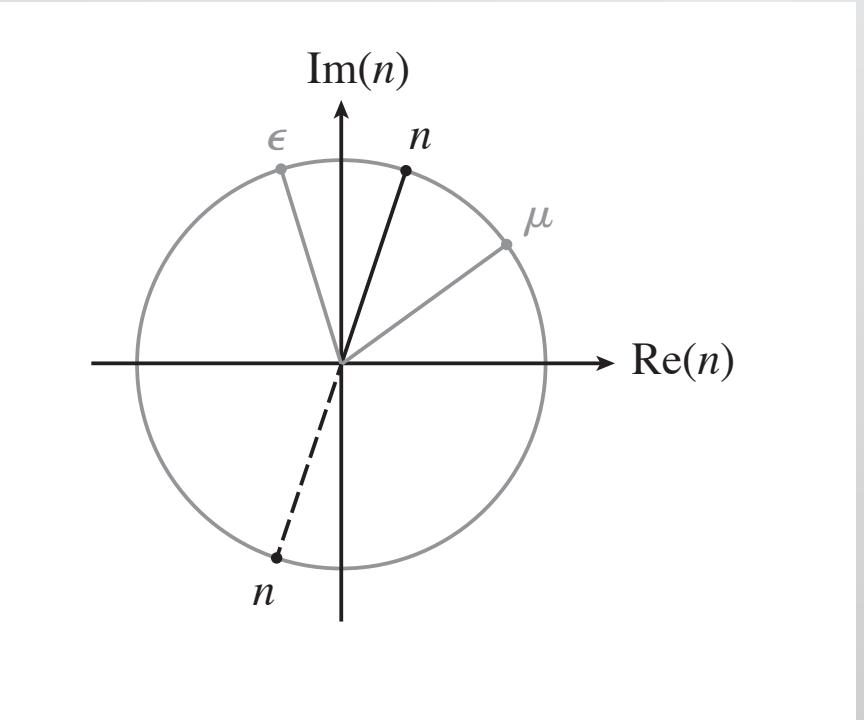
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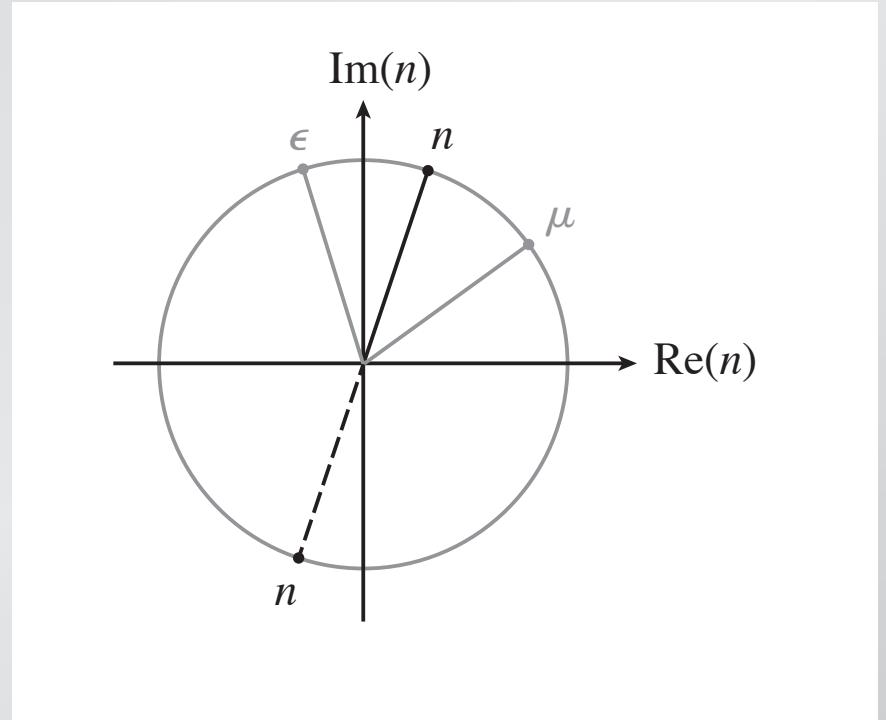
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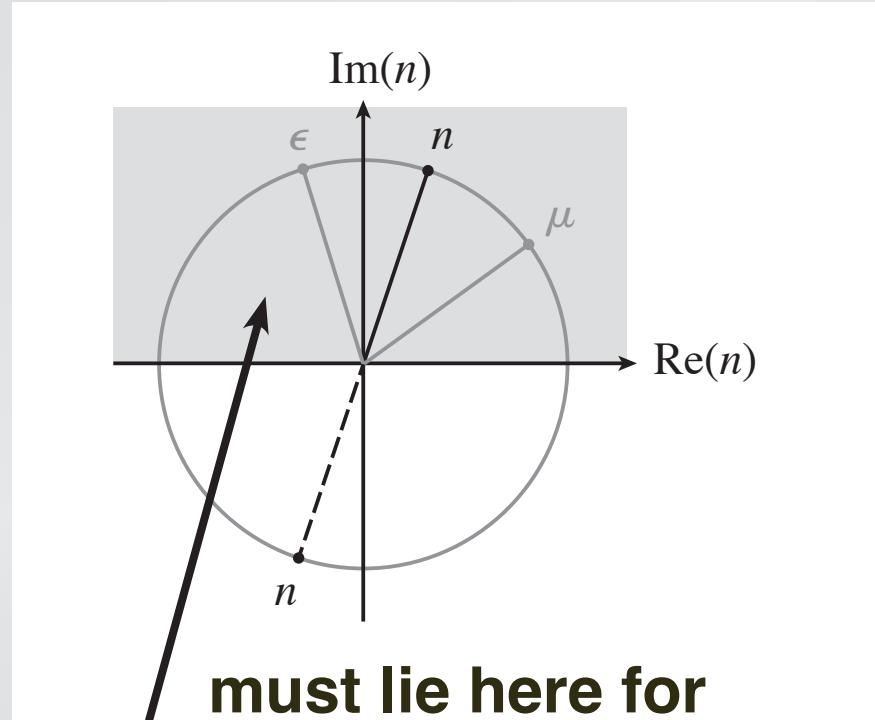
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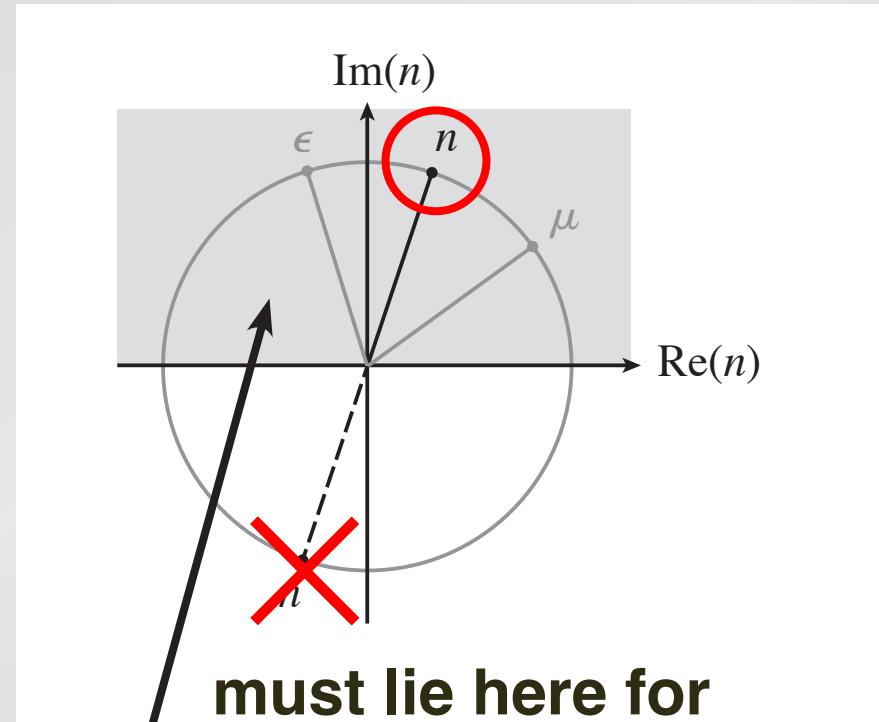
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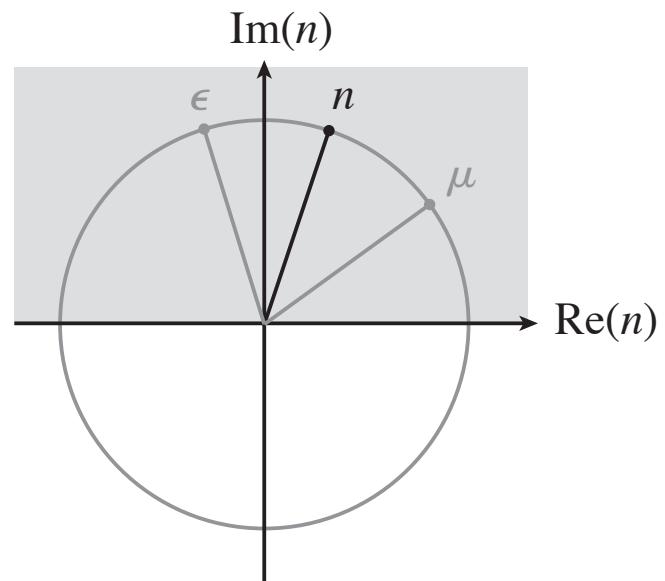
Q: Is this the only possible value?

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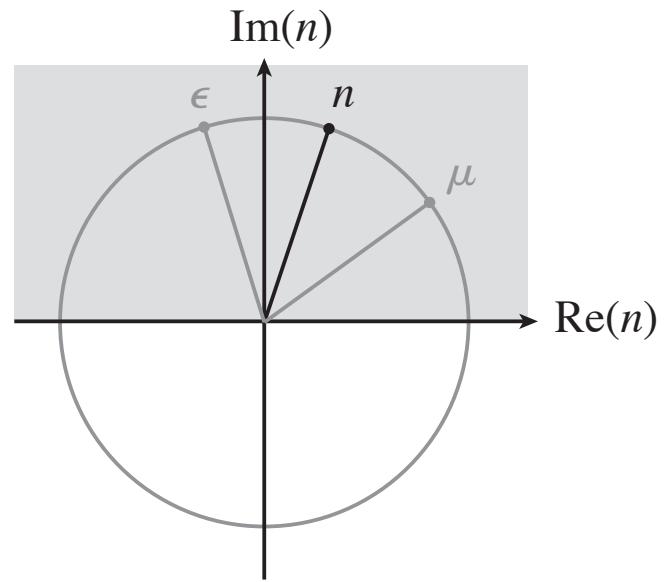
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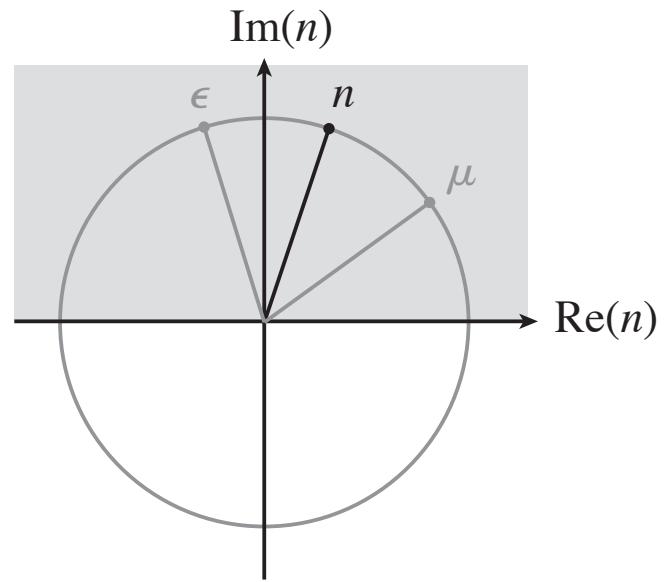
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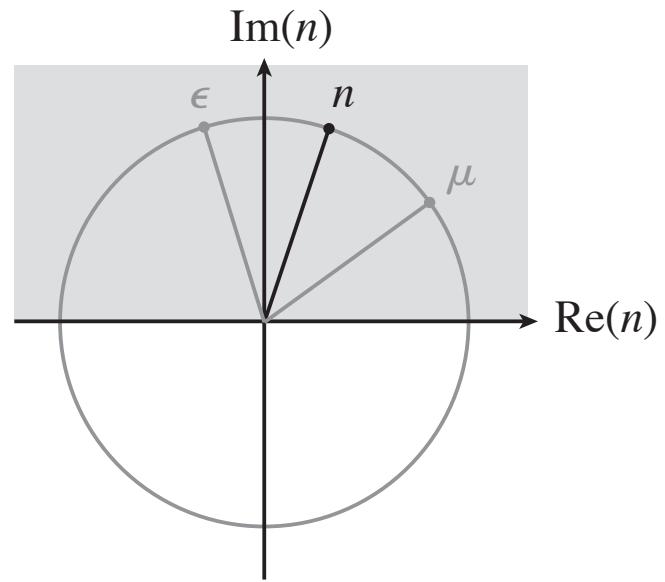
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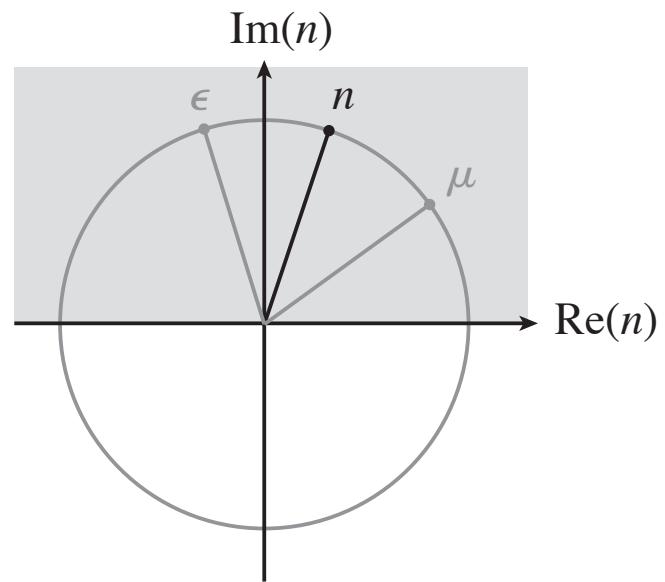
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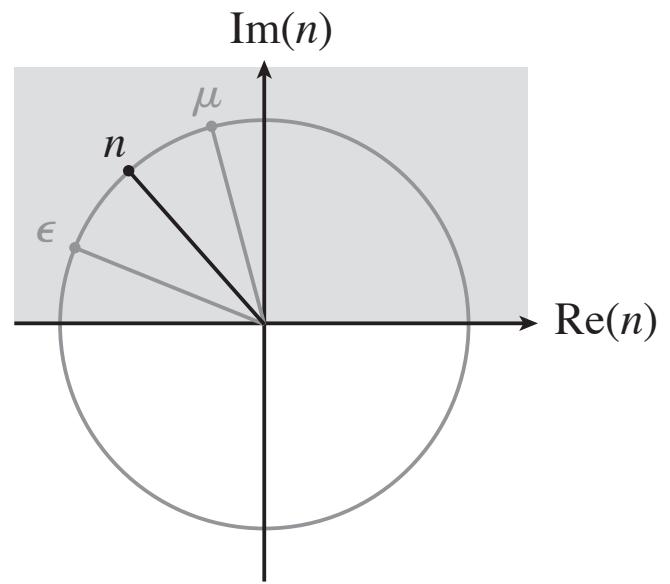
To find n (passive materials):

1. Draw line that bisects ϵ and μ
2. Choose upper branch



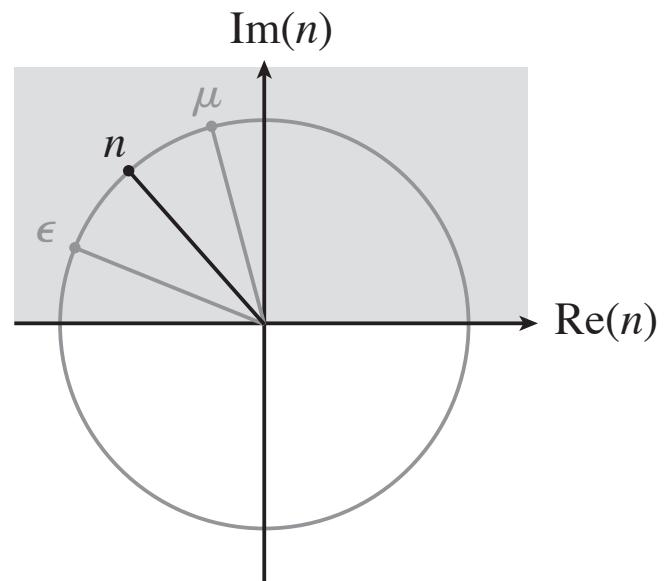
What happens when $\text{Re}\varepsilon$ and/or $\text{Re}\mu$ is negative?

**For certain values of ε and μ
we can get a *negative* $\text{Re}(n)$!**



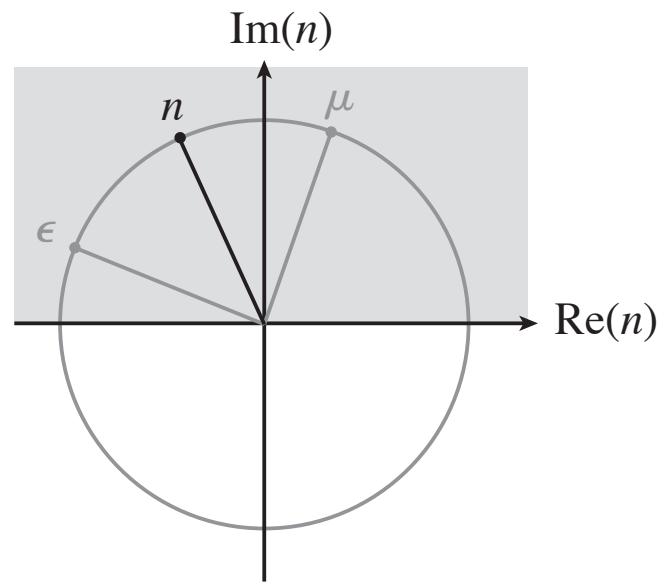
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- 1. yes**
- 2. no**

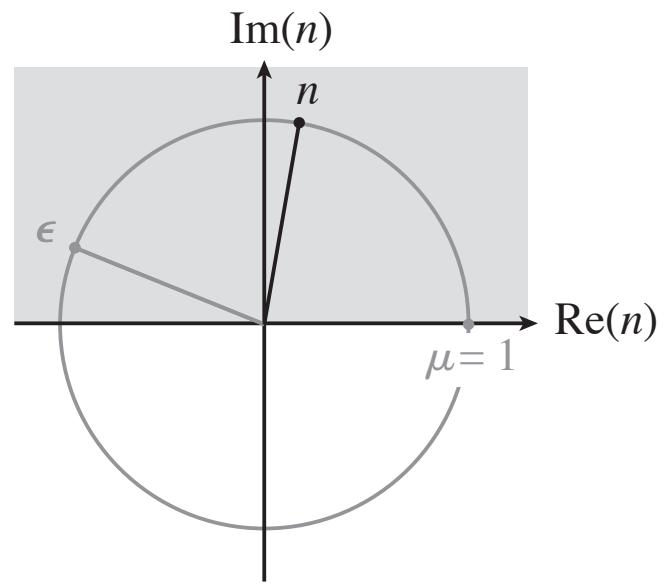


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1. yes
2. no ✓



**However, need magnetic response
to achieve $\text{Re}(n) \leq 0$!**



What happens when $\text{Re}(n) < 0$?

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Remember

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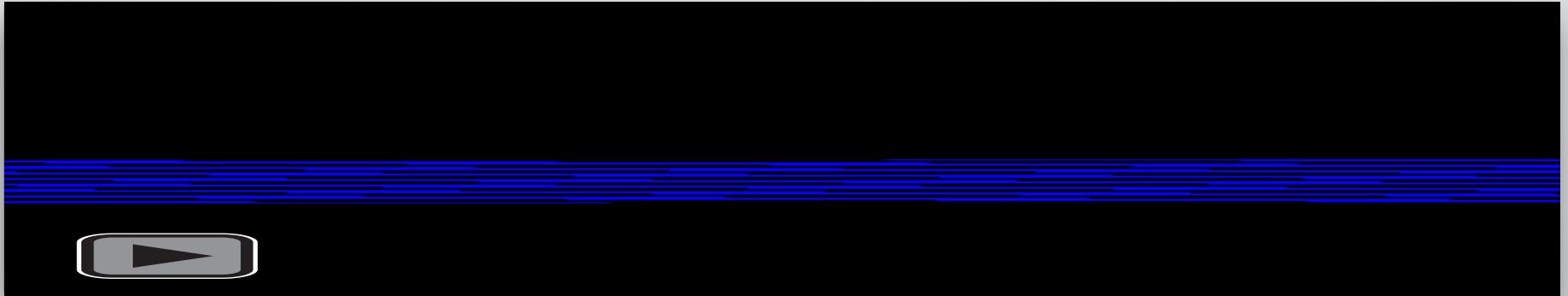
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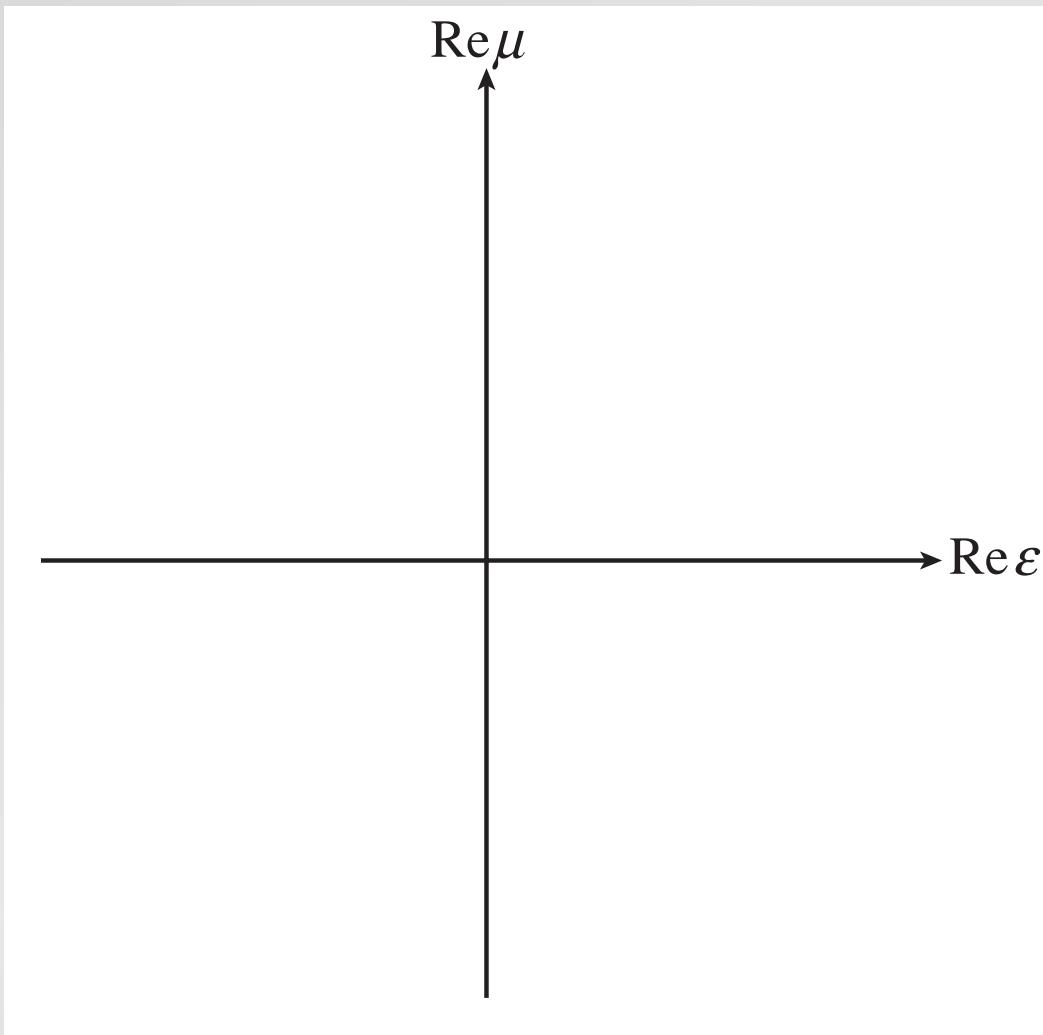
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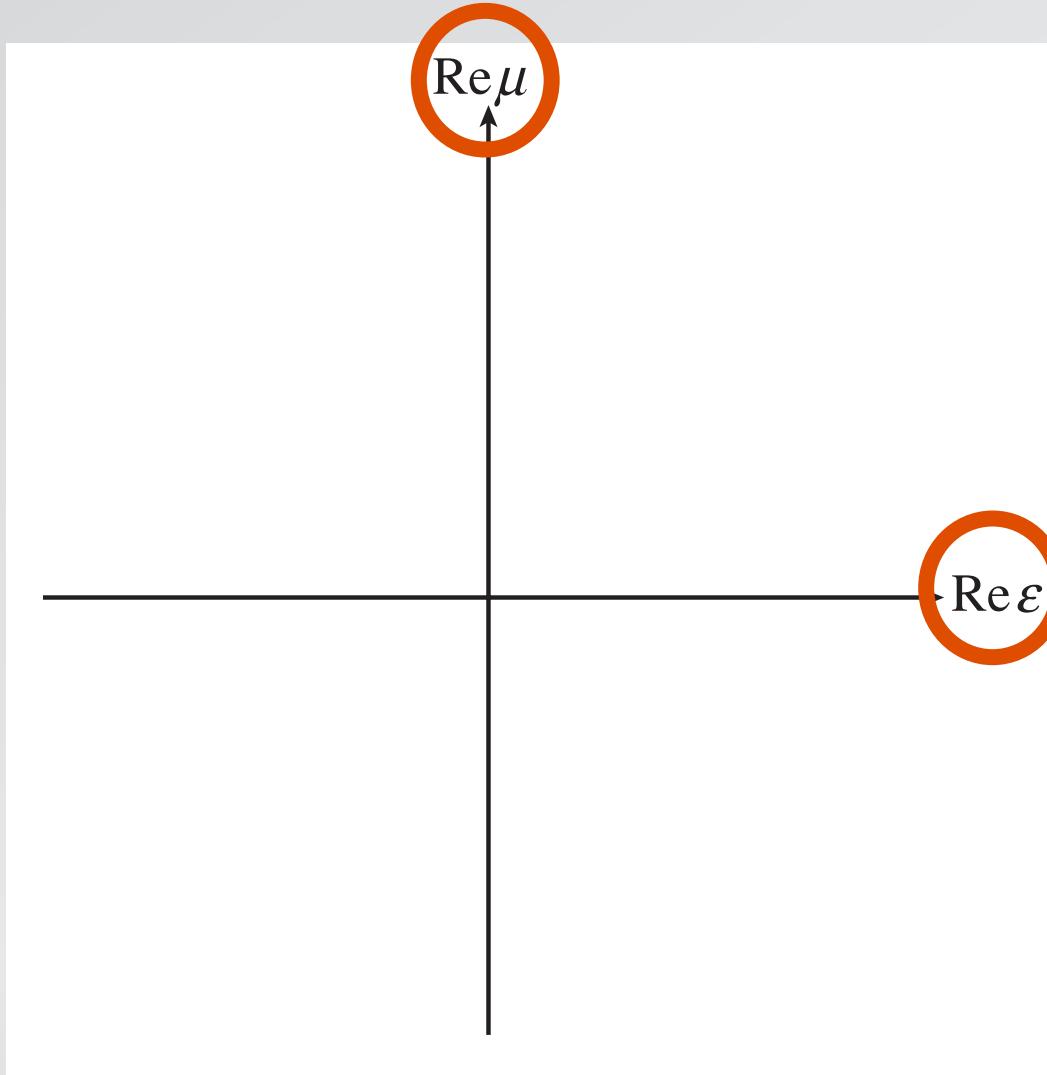


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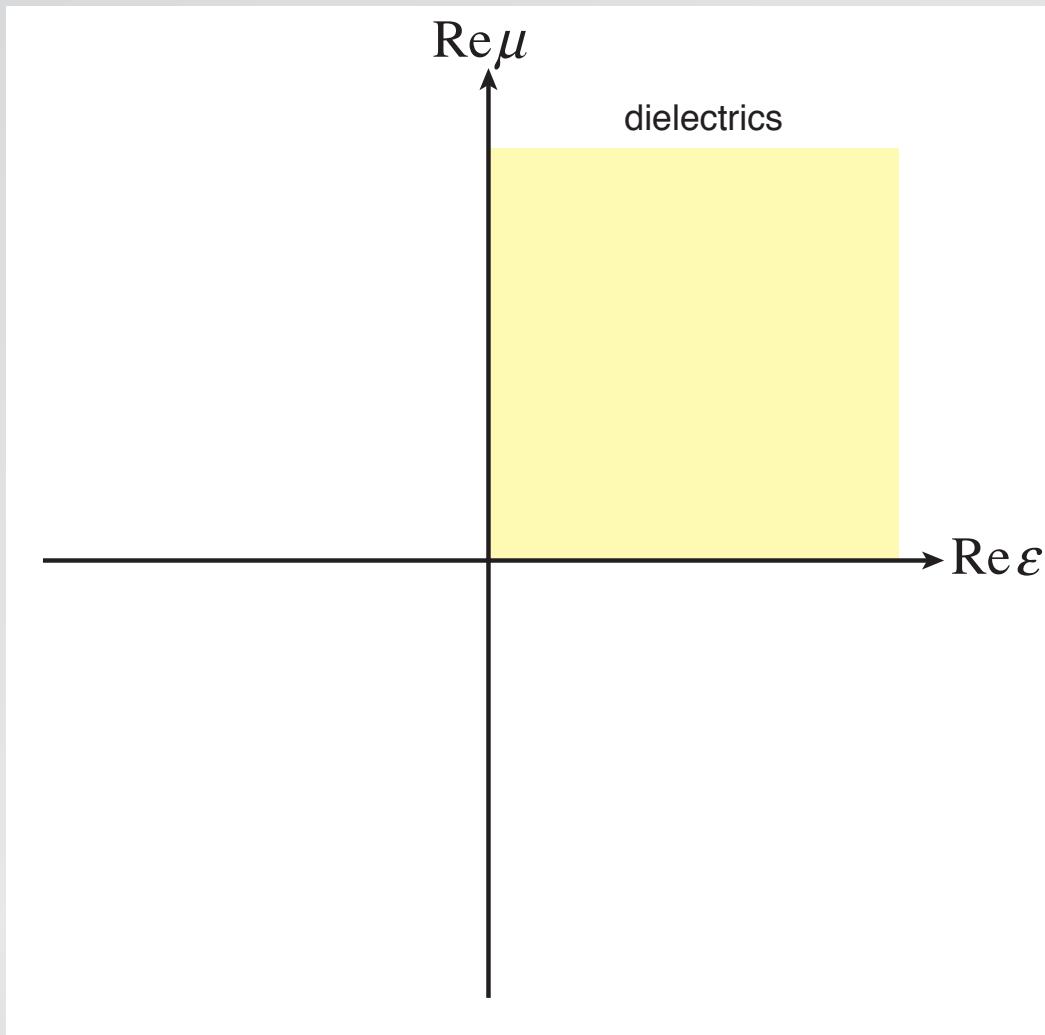
classification of (non-lossy) materials



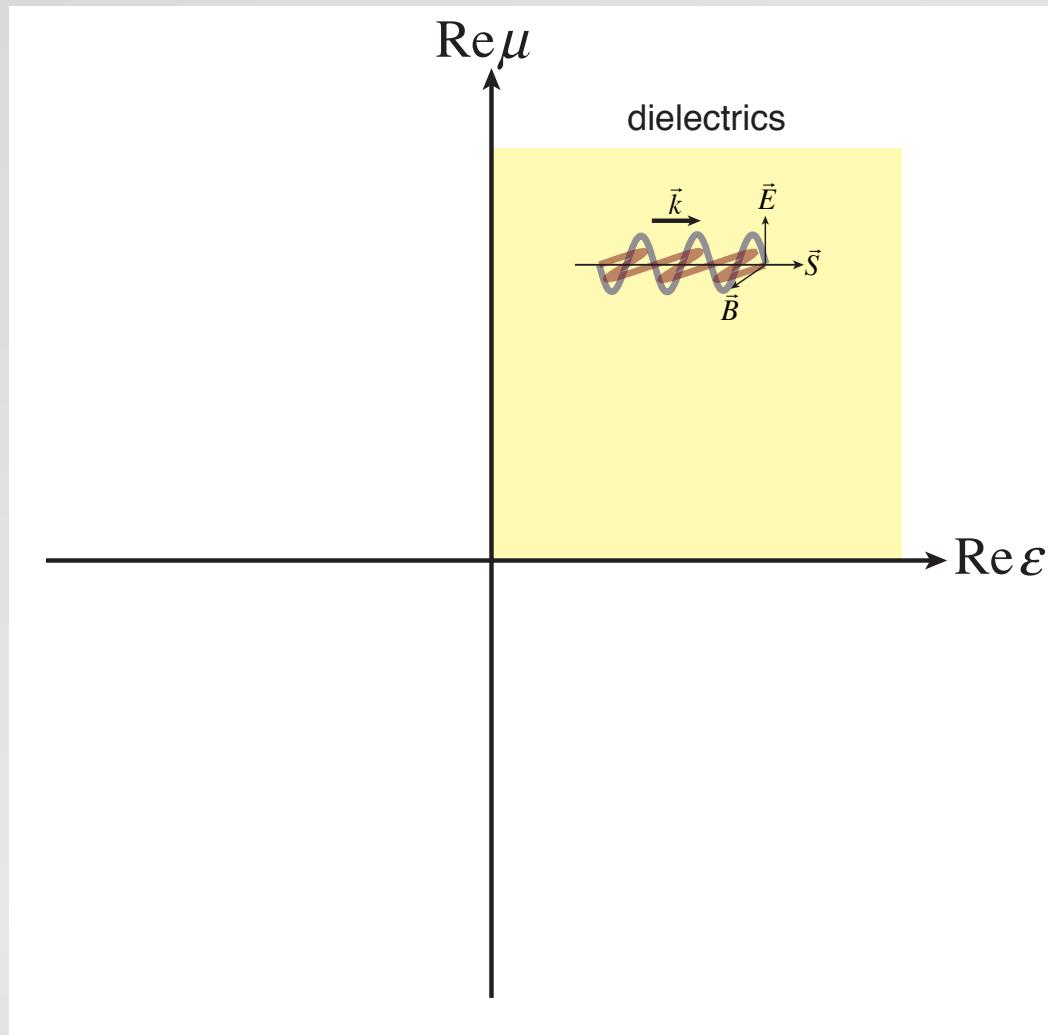
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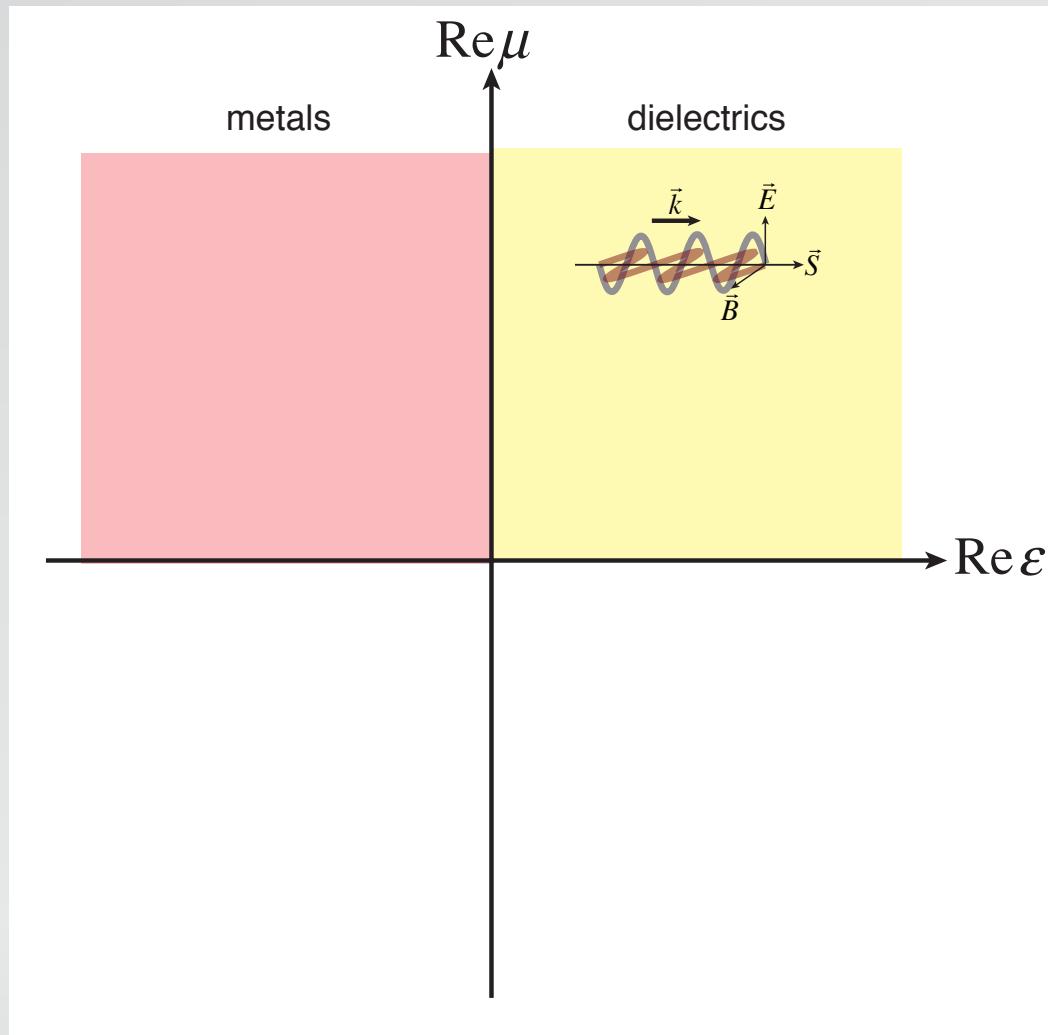
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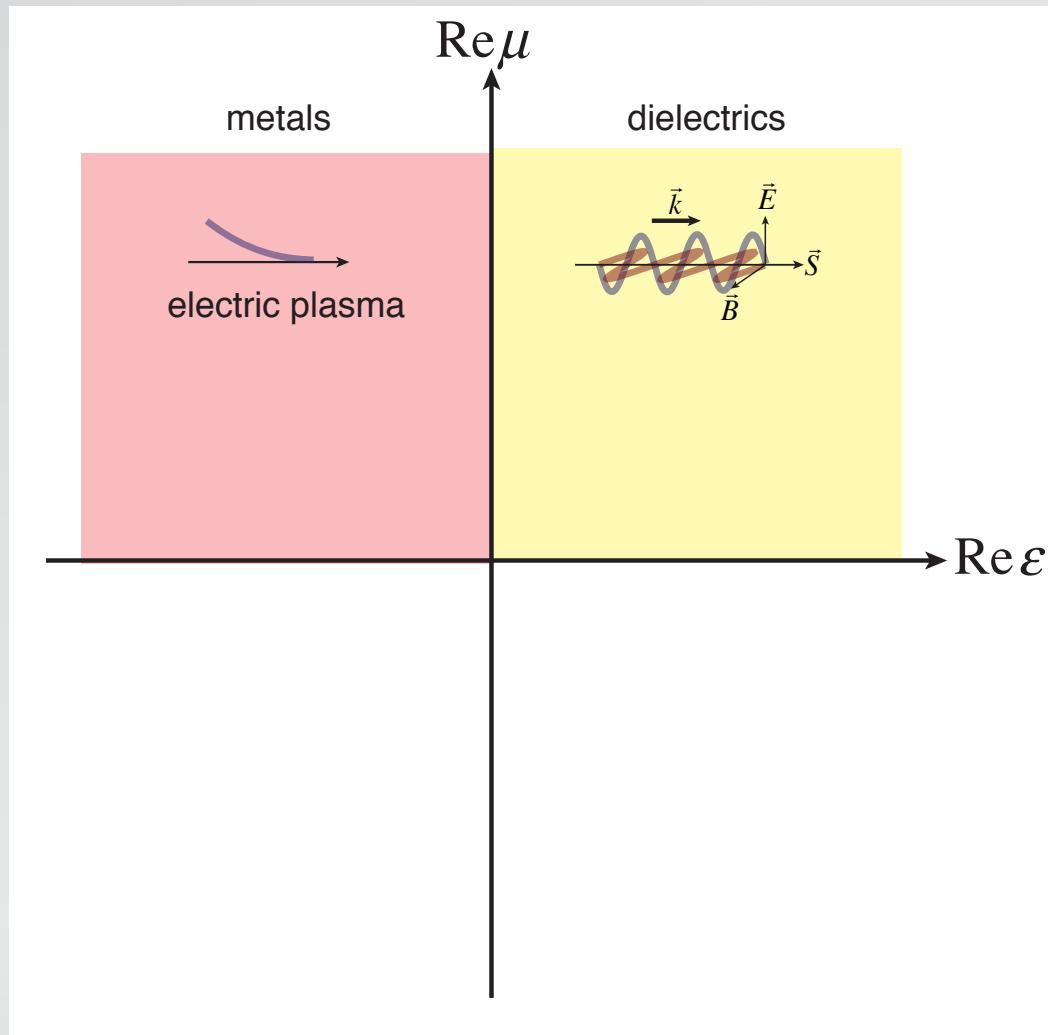
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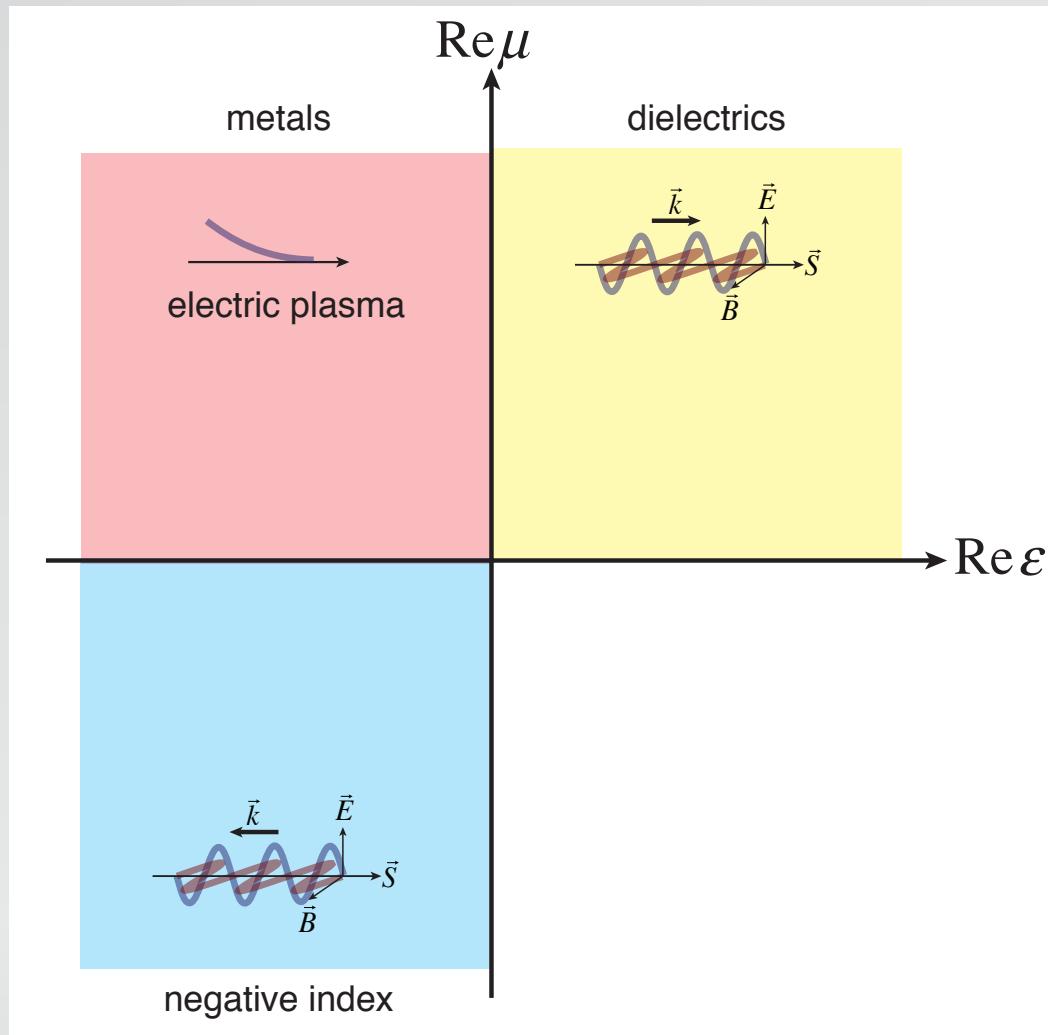
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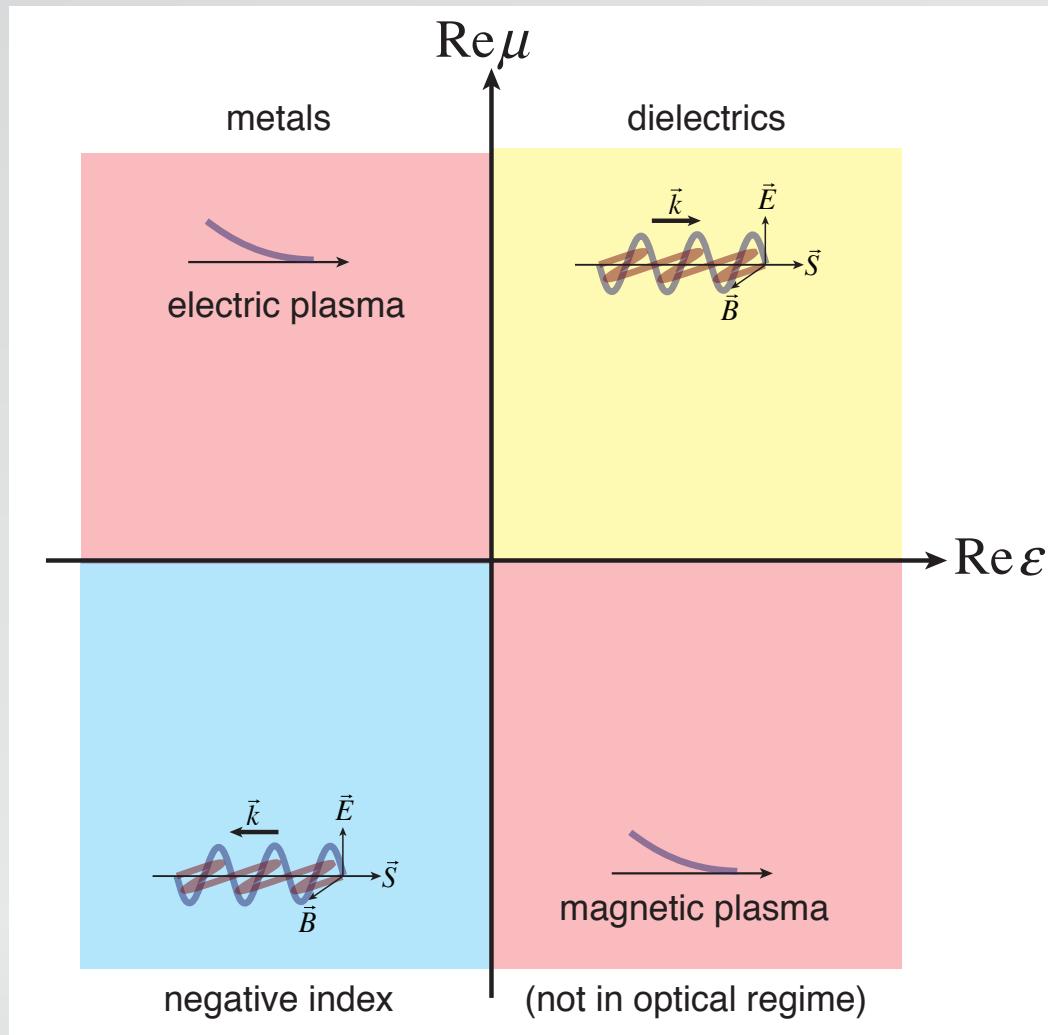
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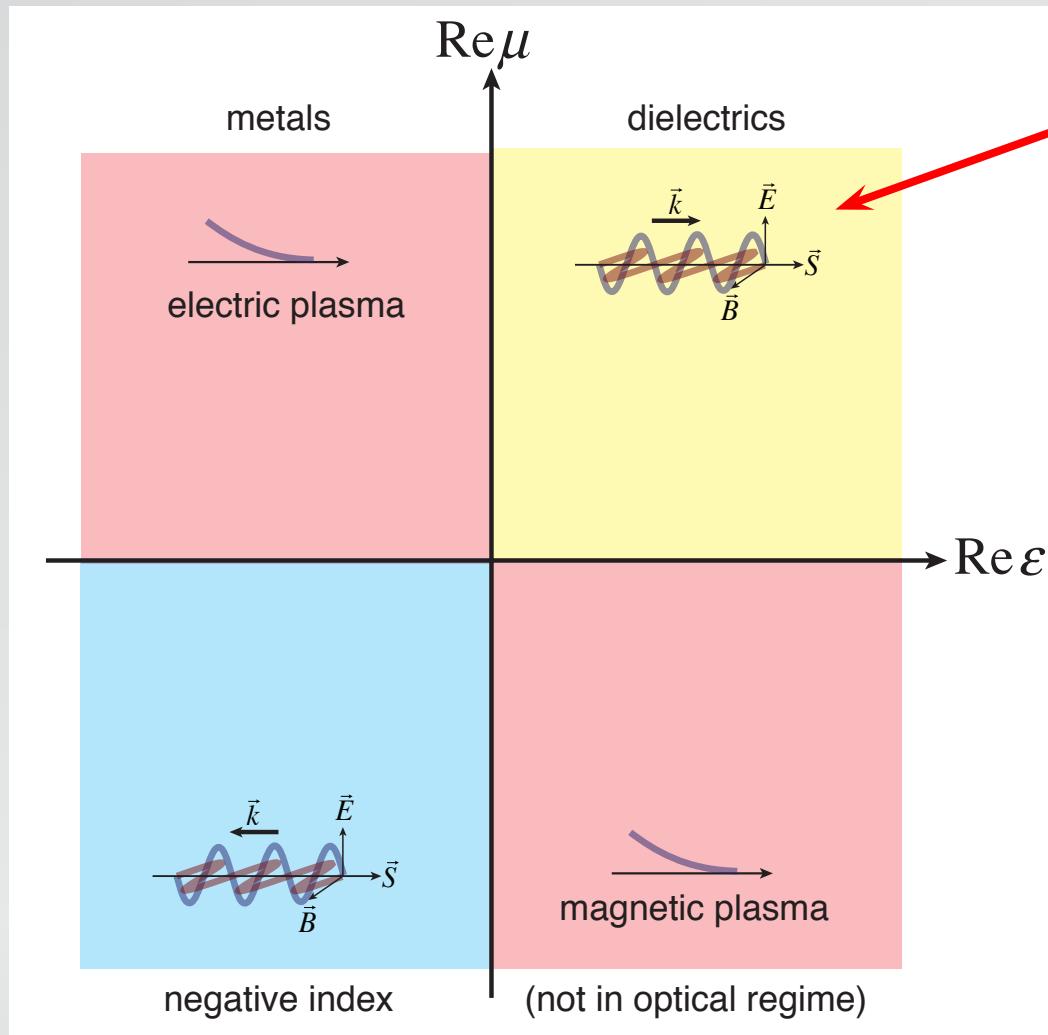
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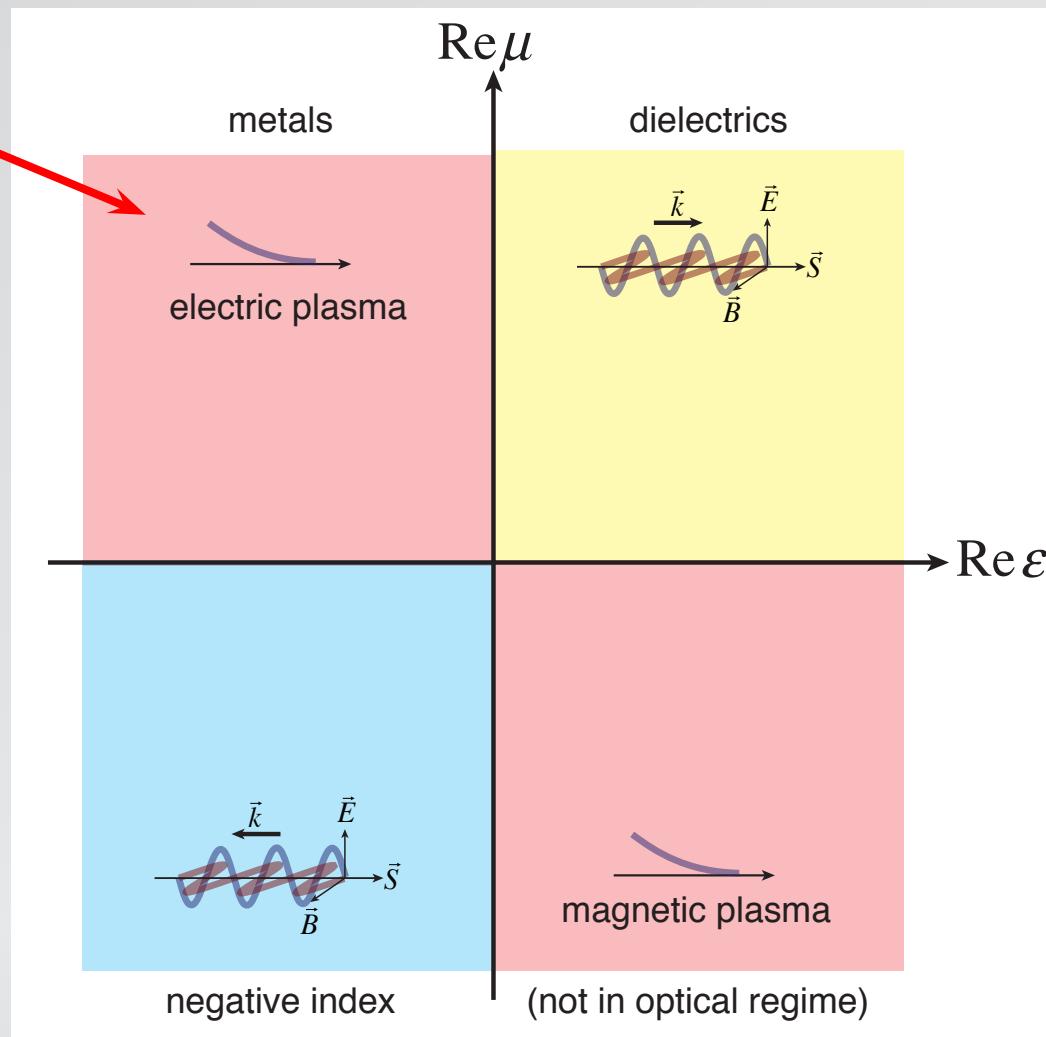
classification of (non-lossy) materials



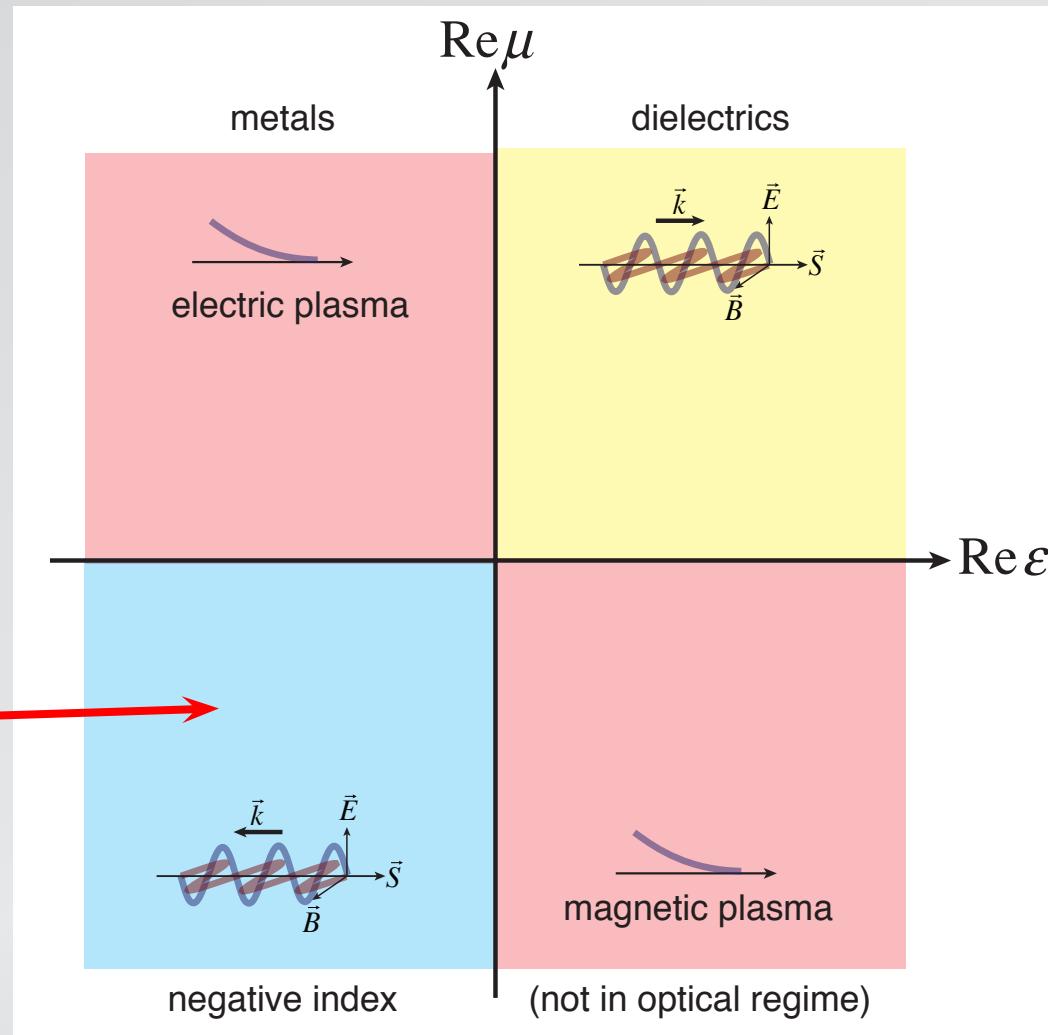
limited by
diffraction

classification of (non-lossy) materials

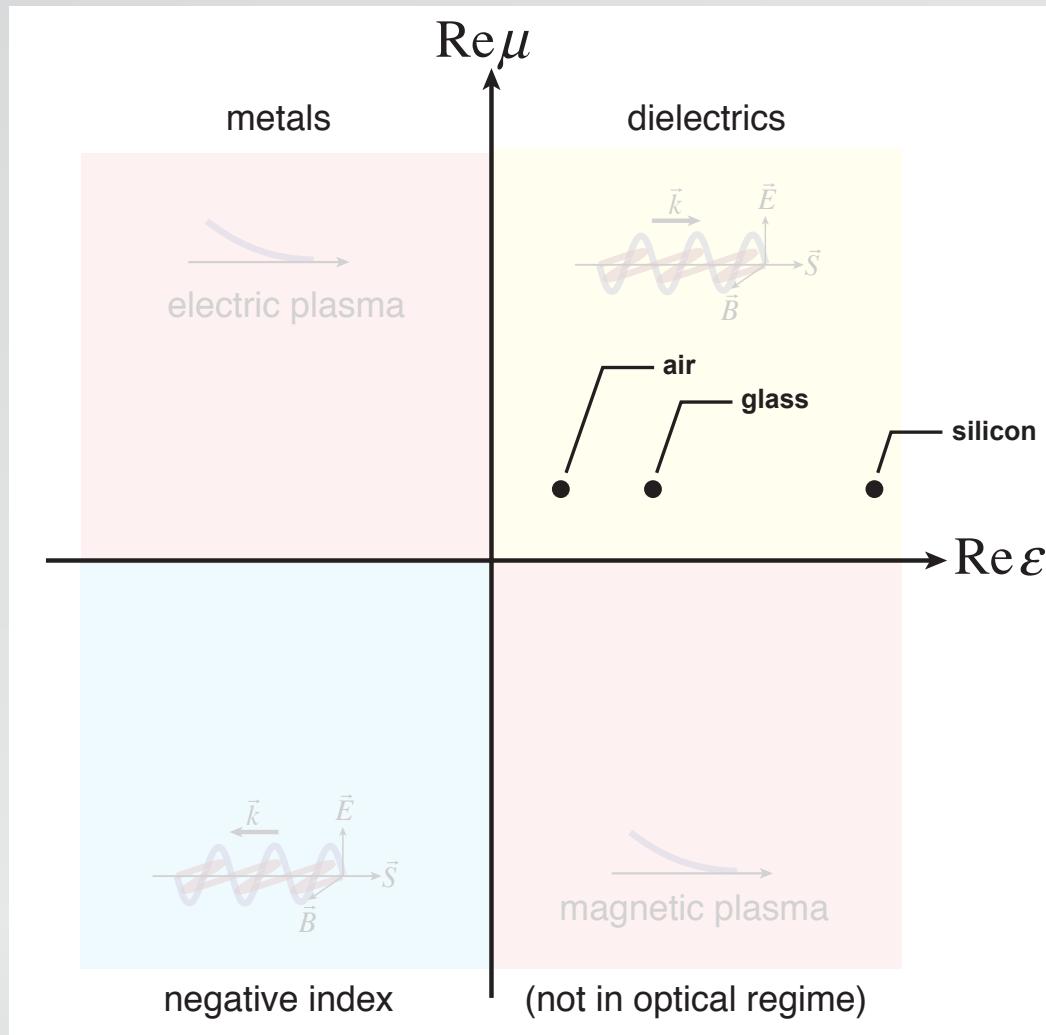
no propagation



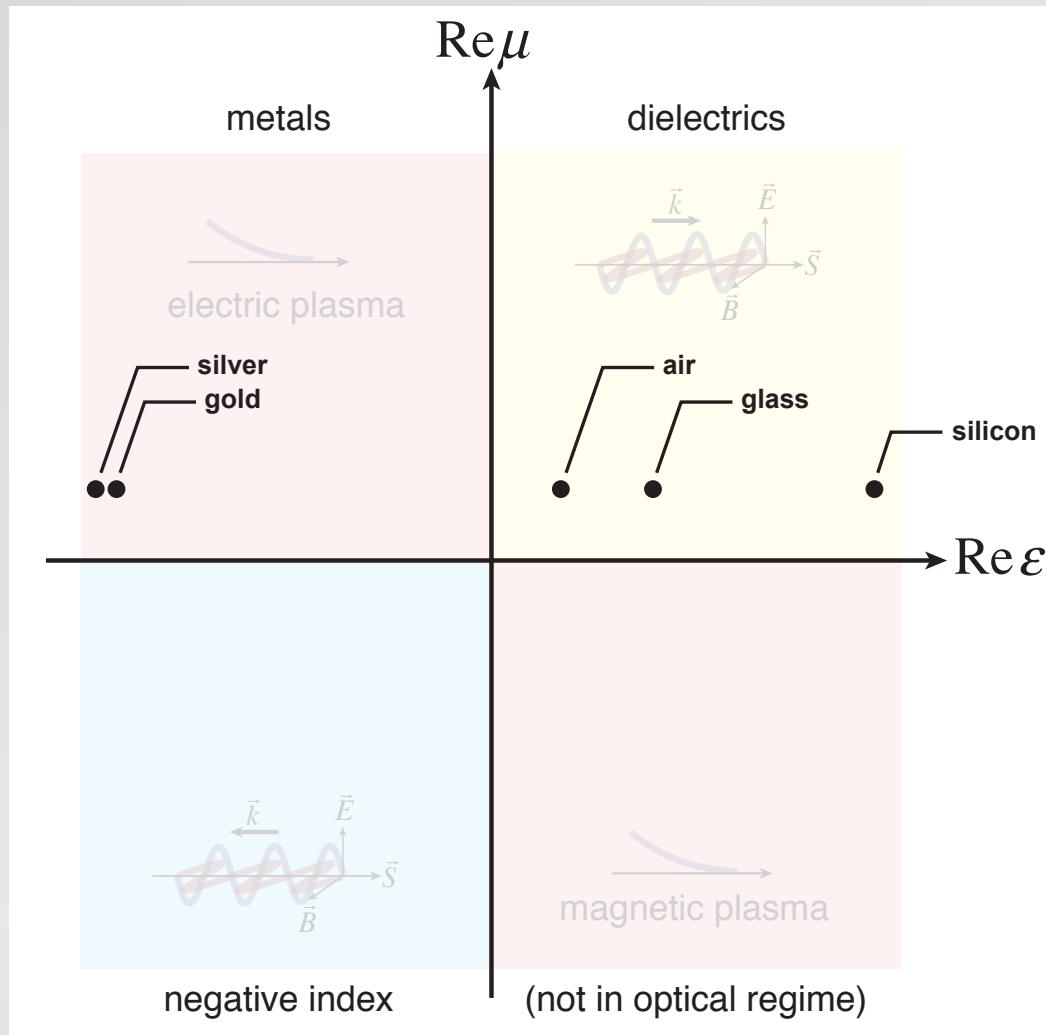
classification of (non-lossy) materials



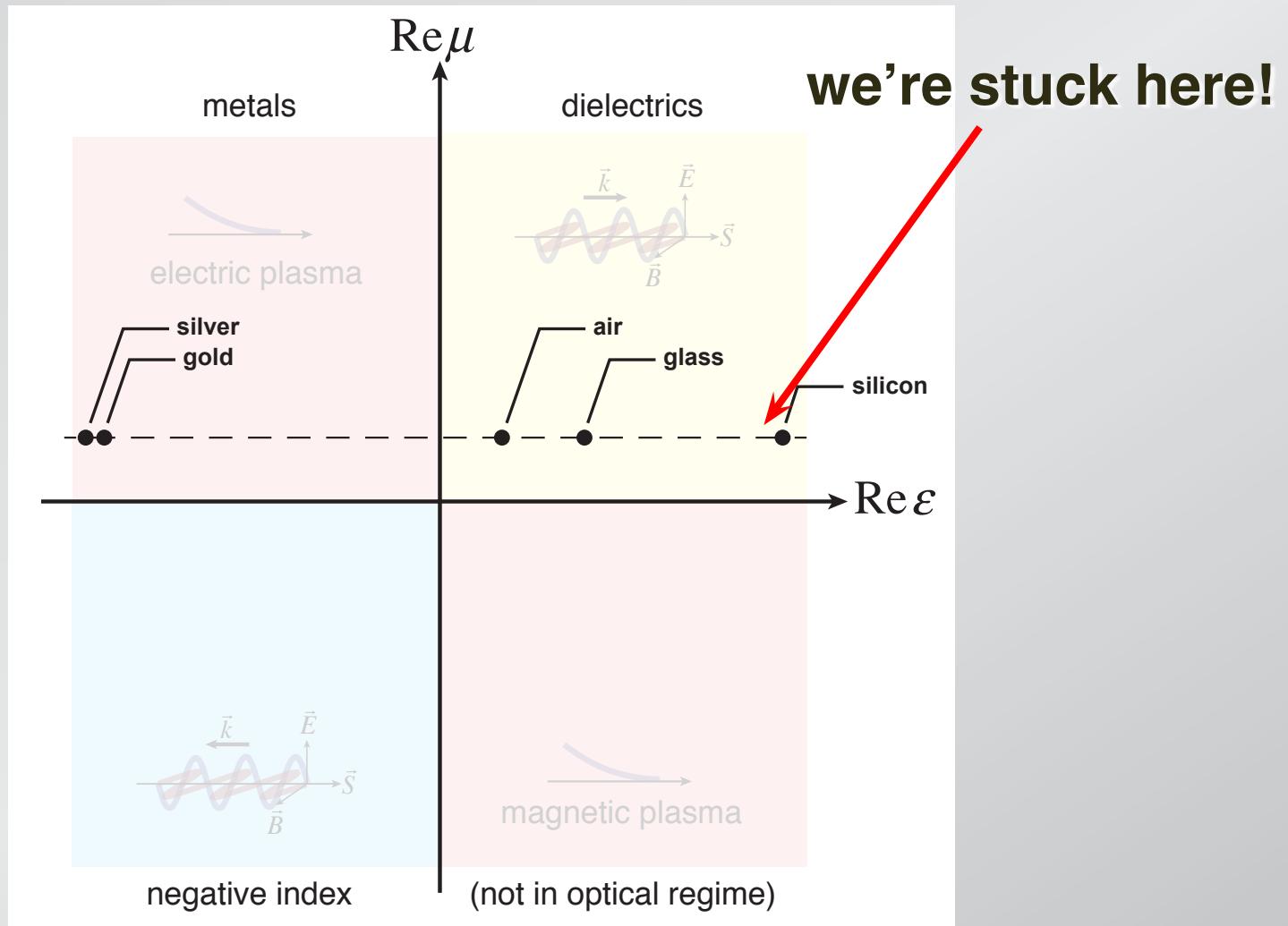
common materials very limited



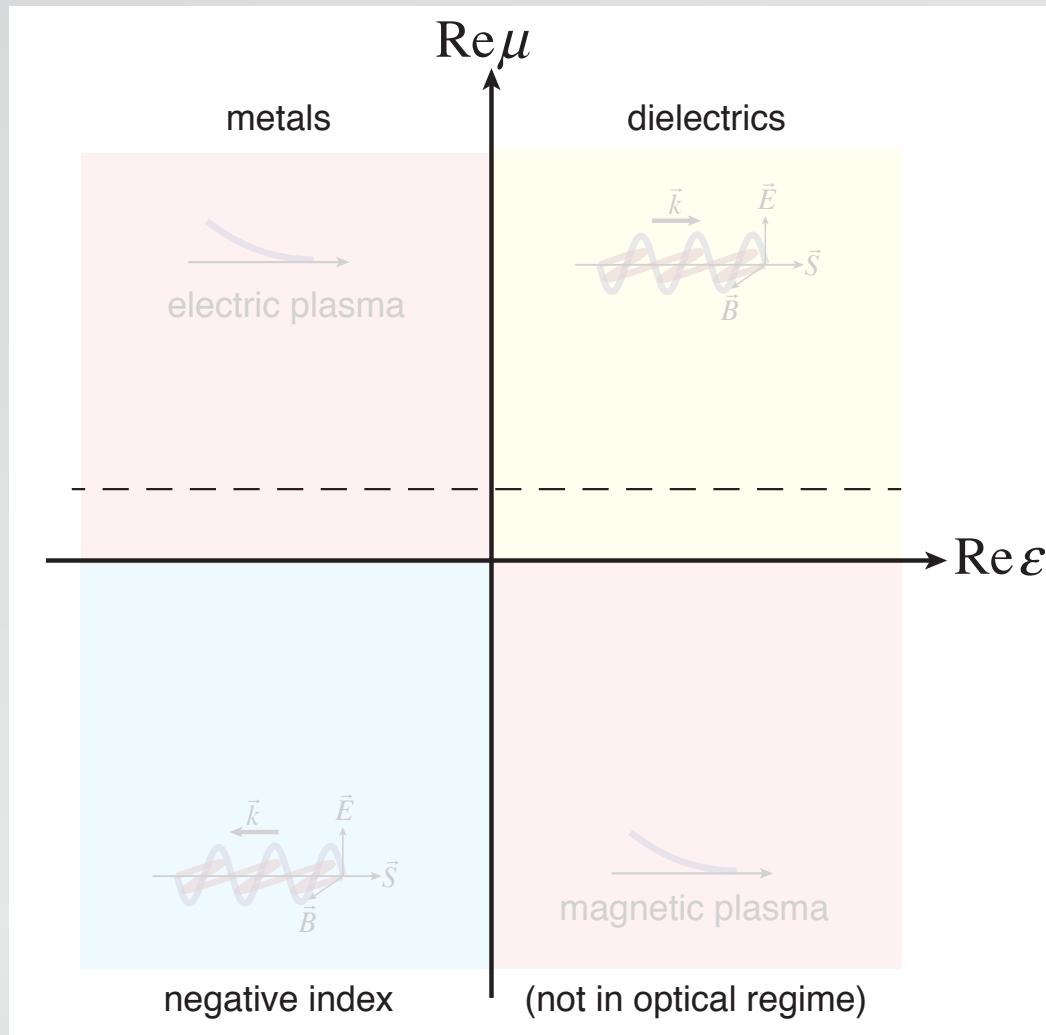
common materials very limited



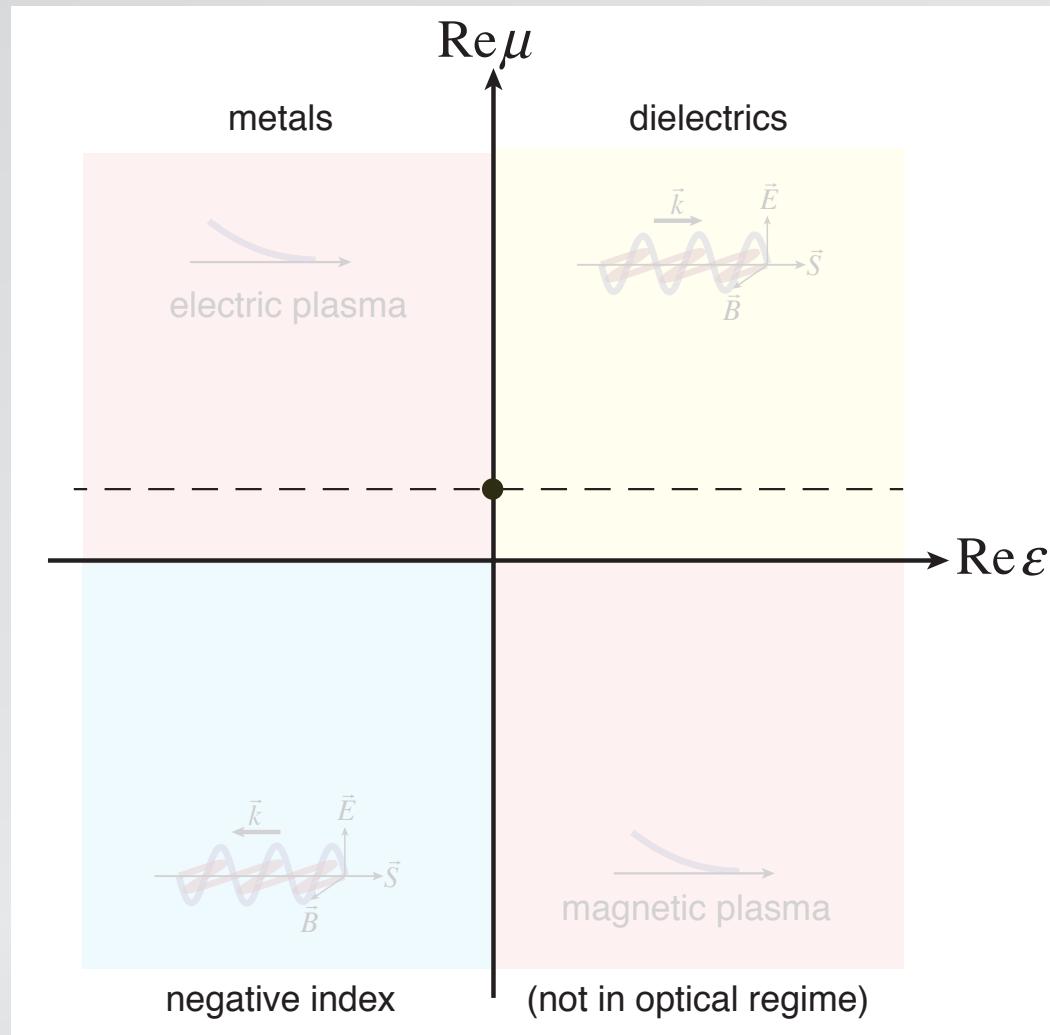
common materials very limited



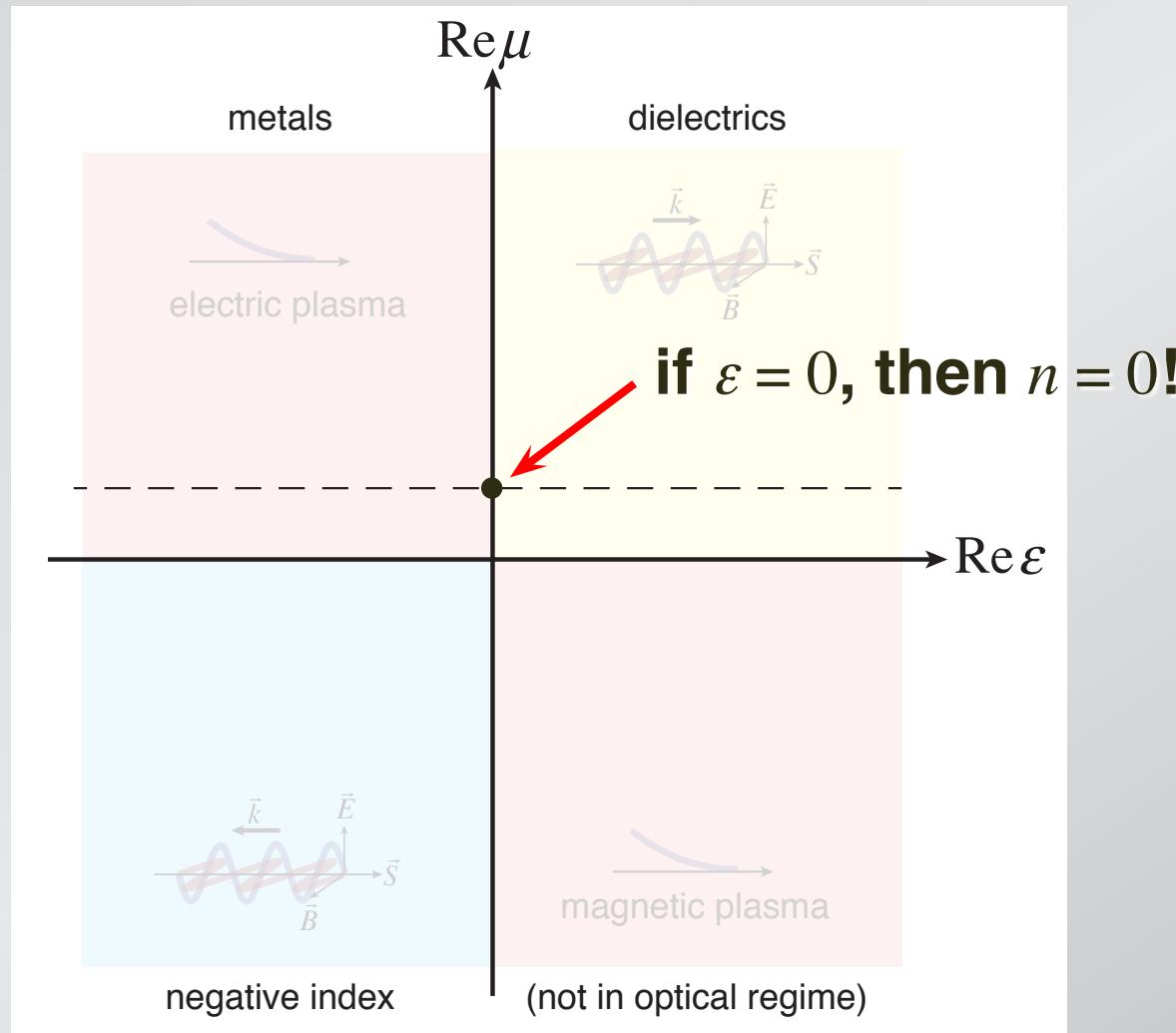
What happens on the axes?



what if we let $\varepsilon = 0$?



what if we let $\varepsilon = 0$?



1 index

2 zero index

Q: If $n = 0$, which of the following is true?

- 1. the frequency goes to zero.**
- 2. the phase velocity becomes infinite.**
- 3. both of the above.**
- 4. neither of the above.**

1 index

2 zero index

wave equation

$$\nabla^2 \vec{E} - \frac{\mu\epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

solution

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)}$$

where

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

1 index

2 zero index

wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon \omega^2 \vec{E}}{c^2 n^2} = 0$$

solution

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)}$$

where

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

1 index

2 zero index

wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon \omega^2 \vec{E}}{c^2 n^2} = 0$$

solution

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)} \rightarrow \vec{E} = \vec{E}_o e^{-i\omega t}$$

where

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

1 index

2 zero index

wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon \omega^2 \vec{E}}{c^2 n^2} = 0$$

solution

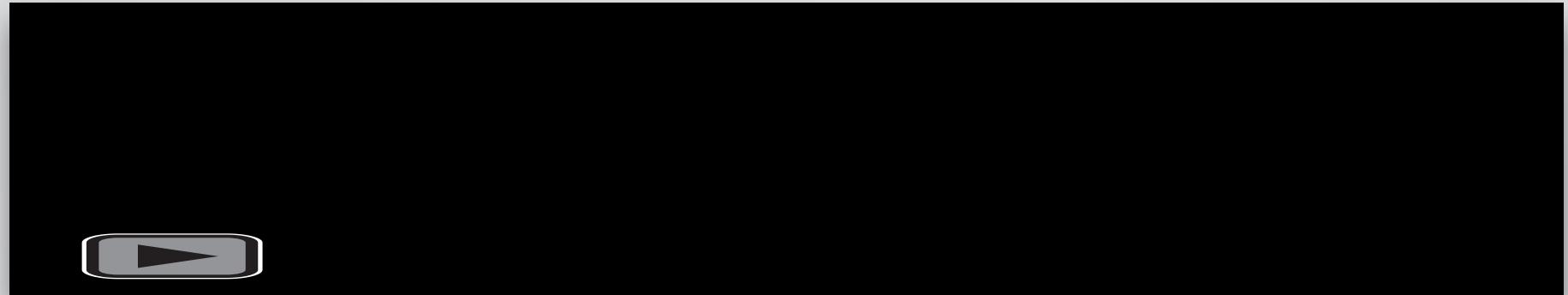
$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)} \longrightarrow \vec{E} = \vec{E}_o e^{-i\omega t}$$

where

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c \longrightarrow \infty$$

① index

② zero index



1 index

2 zero index



“Superluminal”?!

1 index

2 zero index

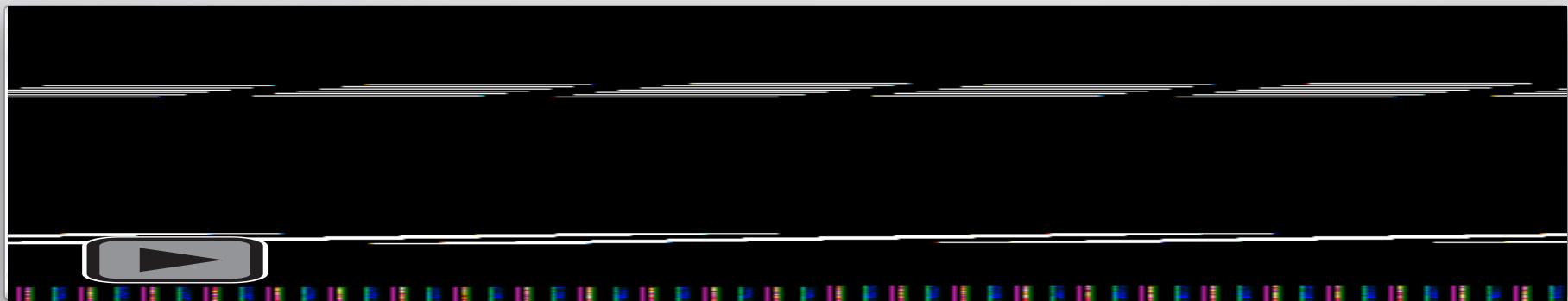
What about

WAVE PROPAGATION AND GROUP VELOCITY

LÉON BRILLOUIN

Member of the National Academy of Sciences

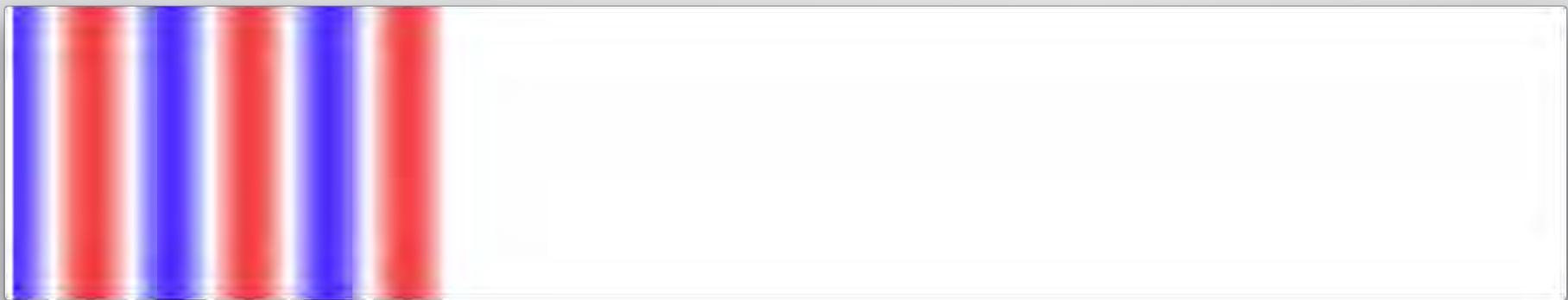
What about causality?



1 index

2 zero index

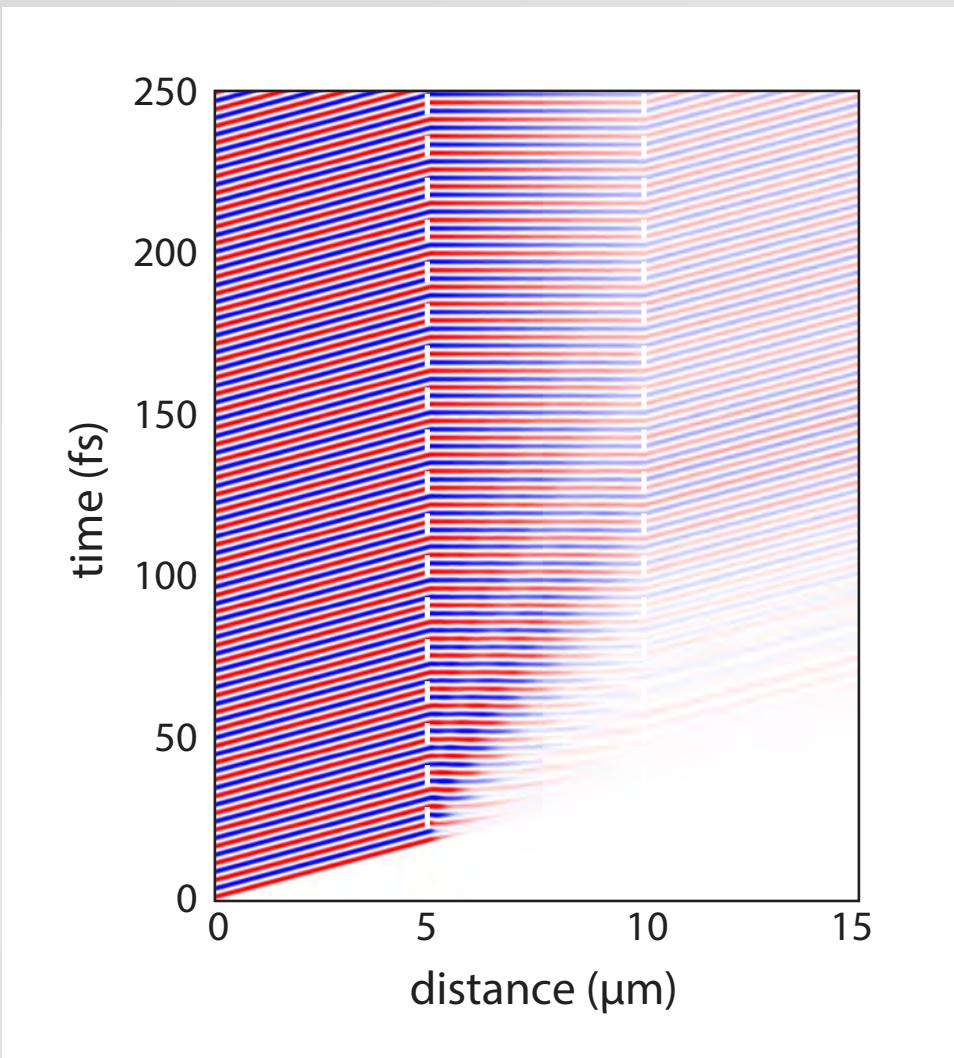
What about causality?



1 index

2 zero index

What about causality?

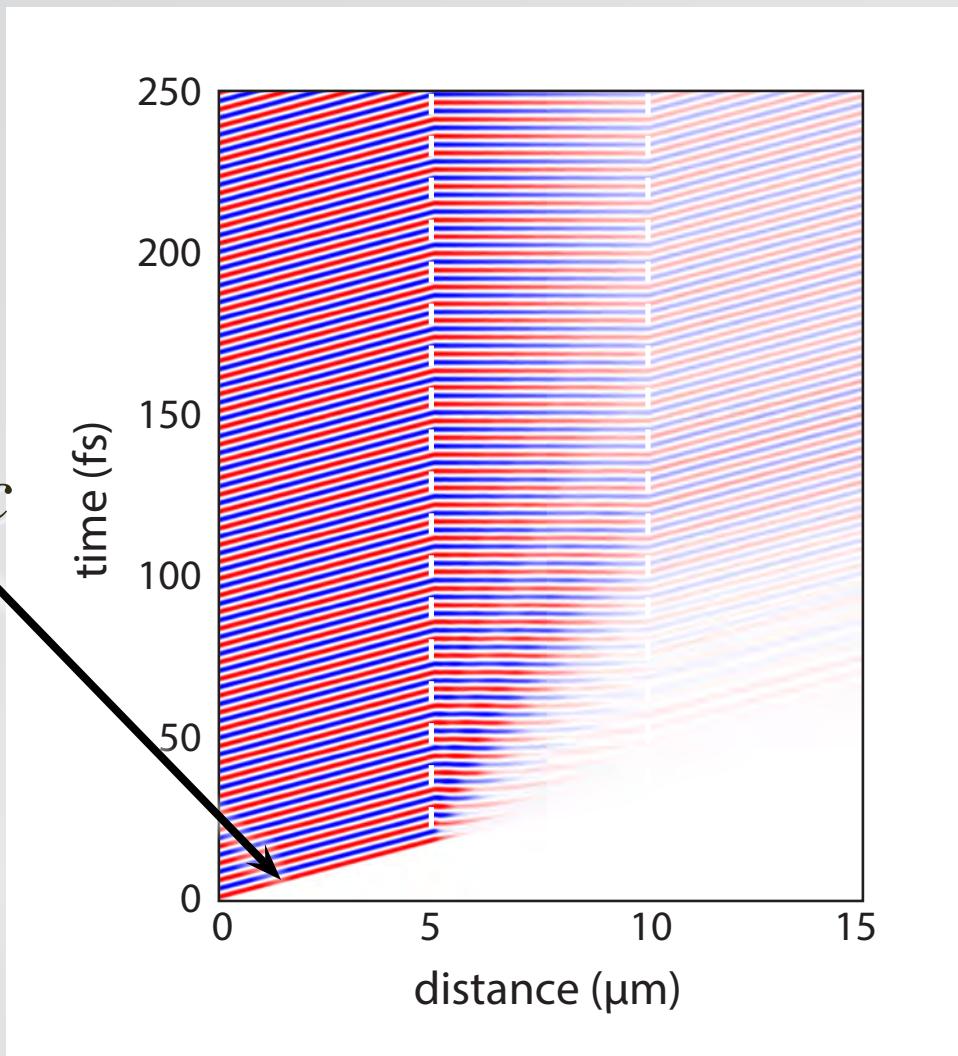


1 index

2 zero index

What about causality?

speed of light c

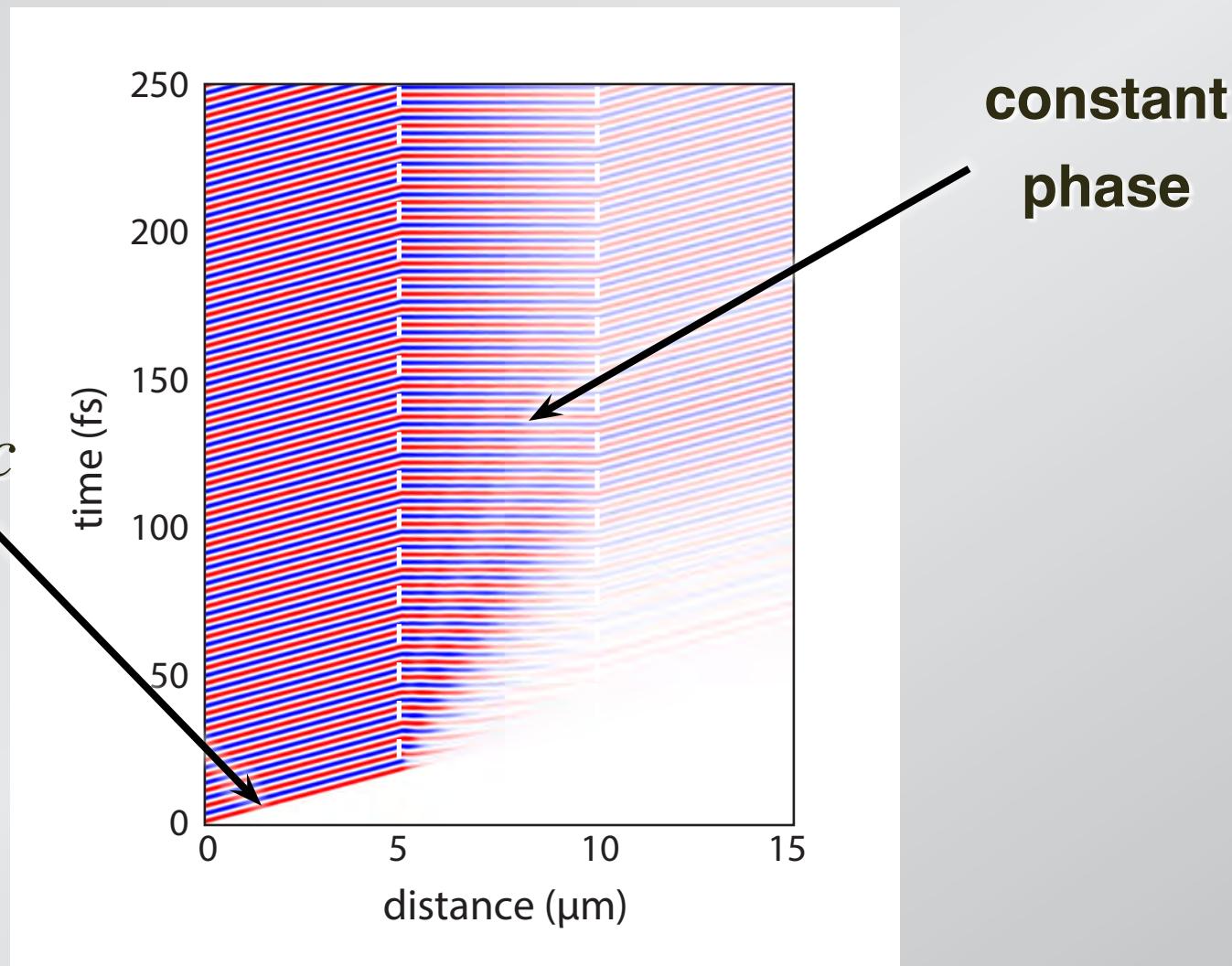


1 index

2 zero index

What about causality?

speed of light c



1 index

2 zero index

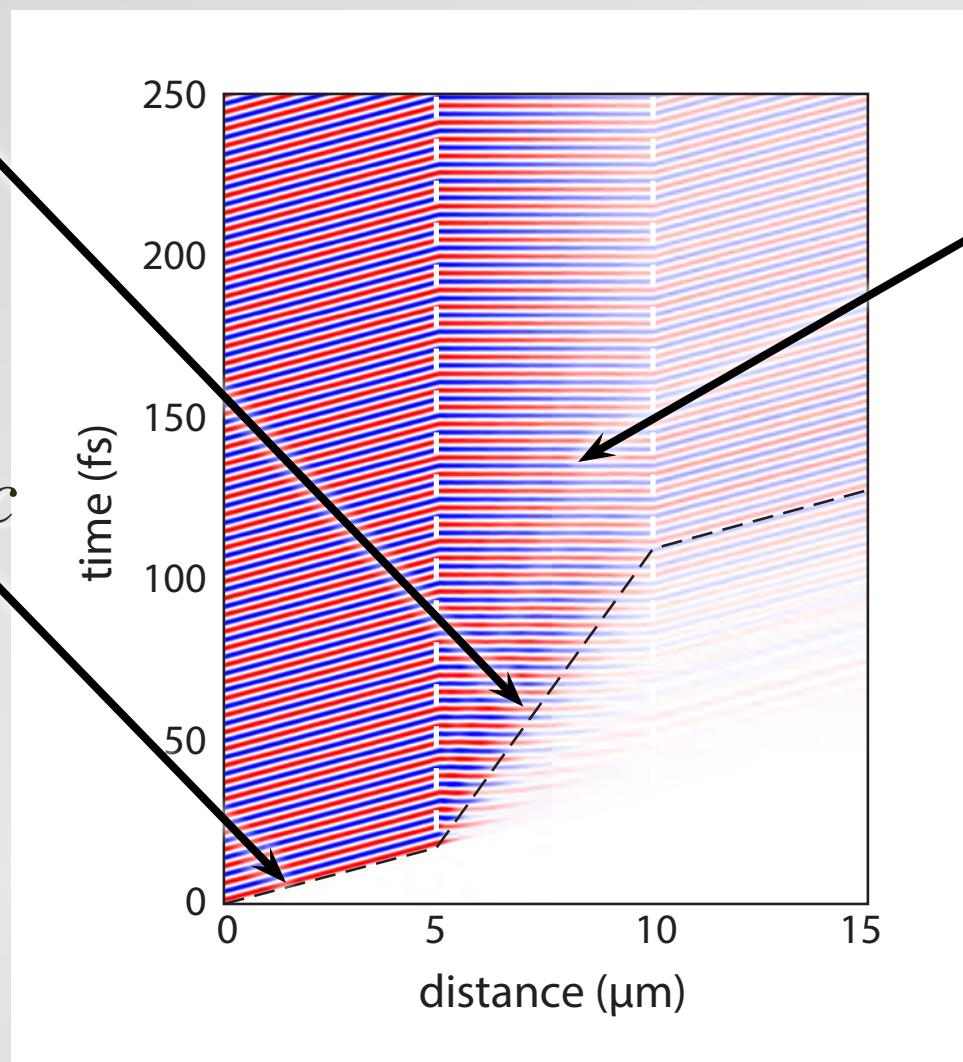
What about causality?

group velocity

$$v_g < c$$

speed of light c

constant phase



1 index

2 zero index

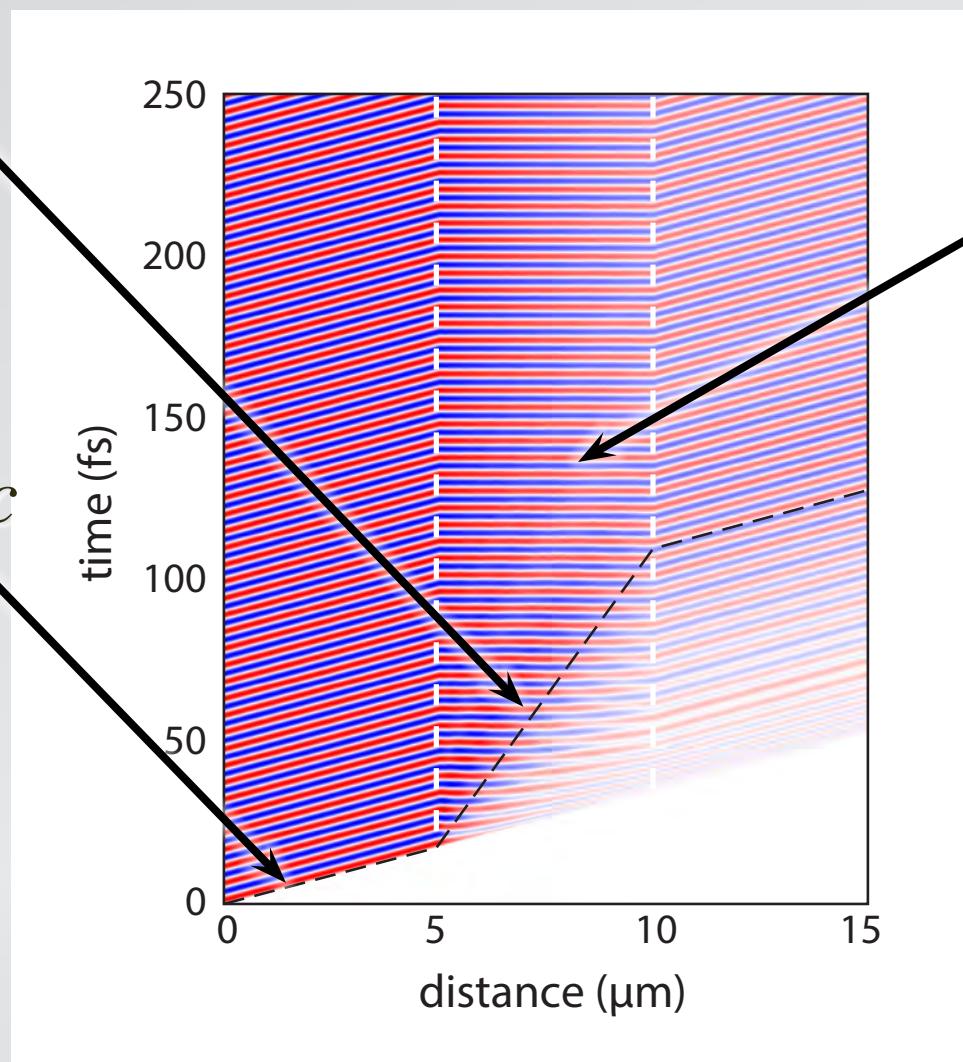
What about causality?

group velocity

$$v_g < c$$

speed of light c

constant phase



1 index

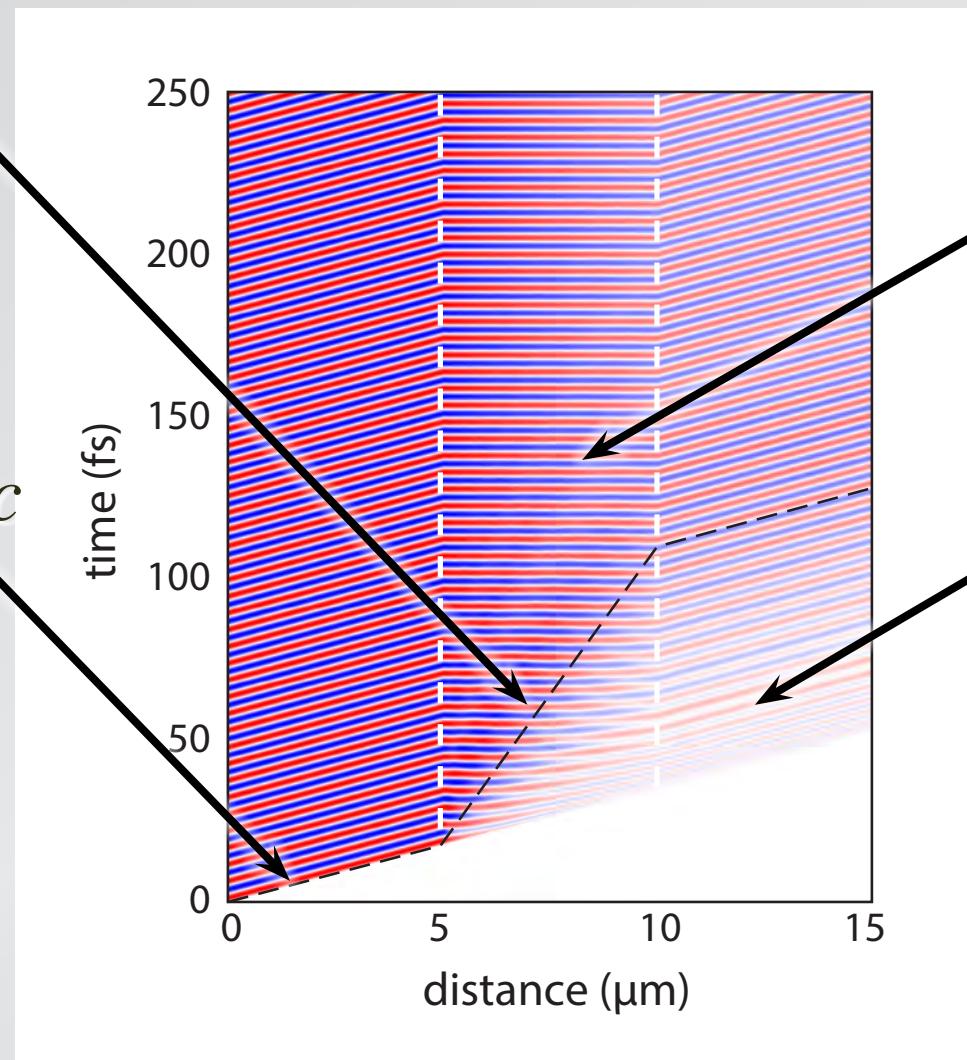
2 zero index

What about causality?

group velocity

$$v_g < c$$

speed of light c



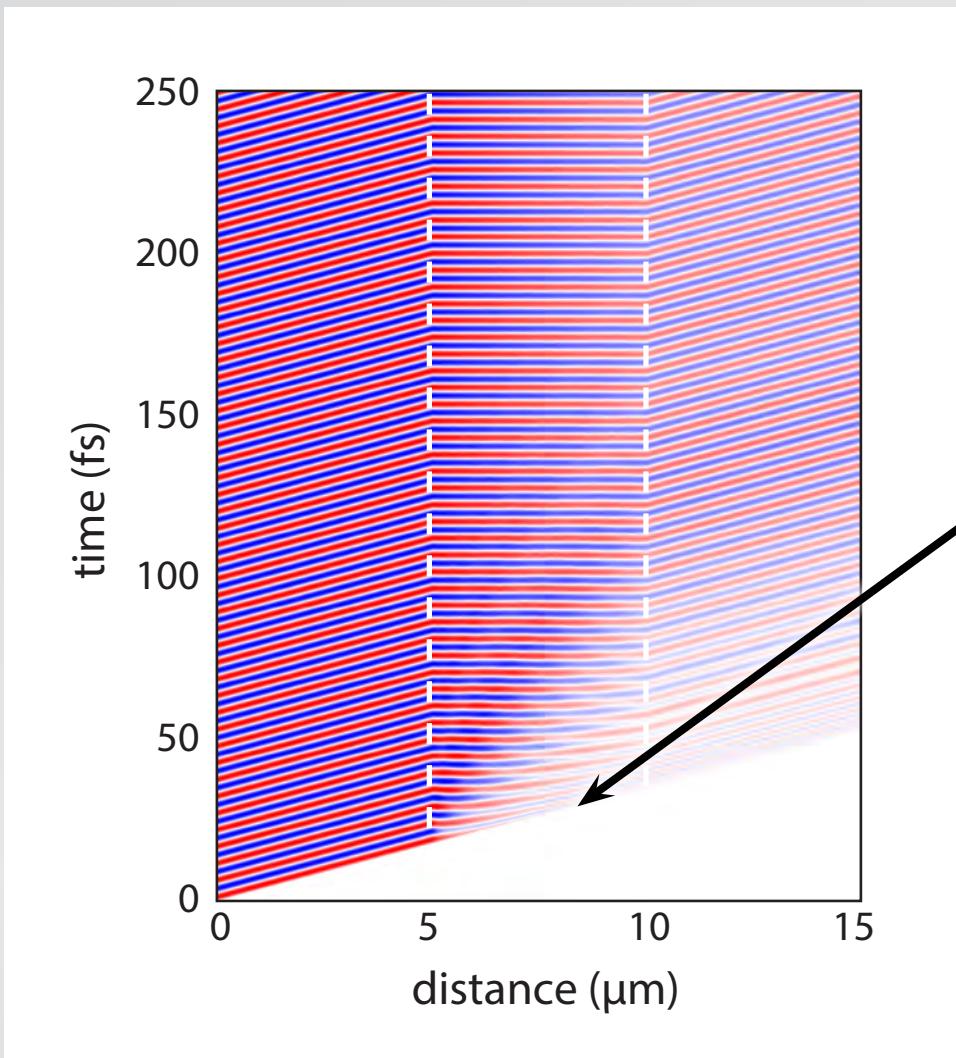
constant
phase

high-frequency
precursors

1 index

2 zero index

What about causality?

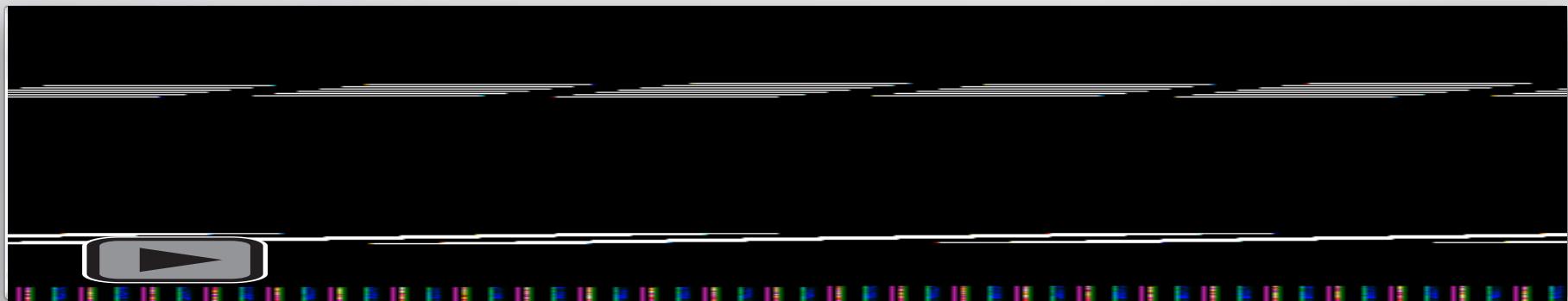


**signal *always*
travels at speed c !**

1 index

2 zero index

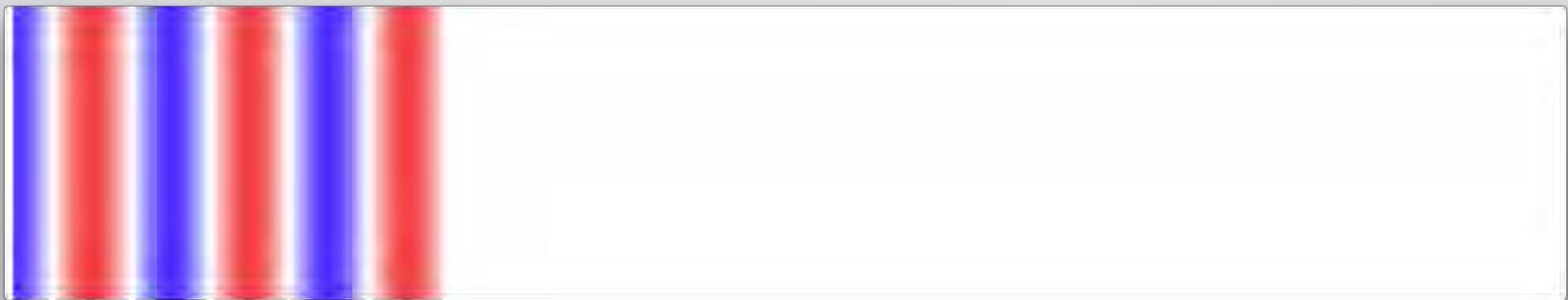
What about causality?



1 index

2 zero index

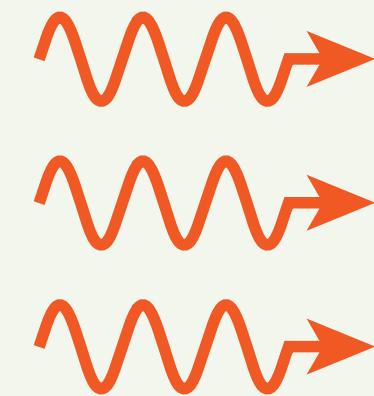
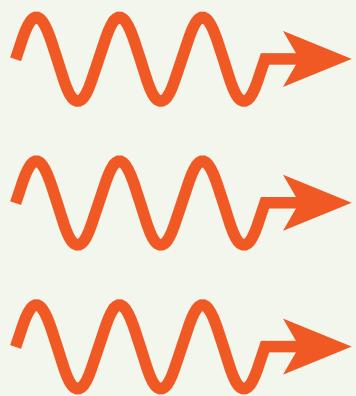
What about causality?



1 index

2 zero index

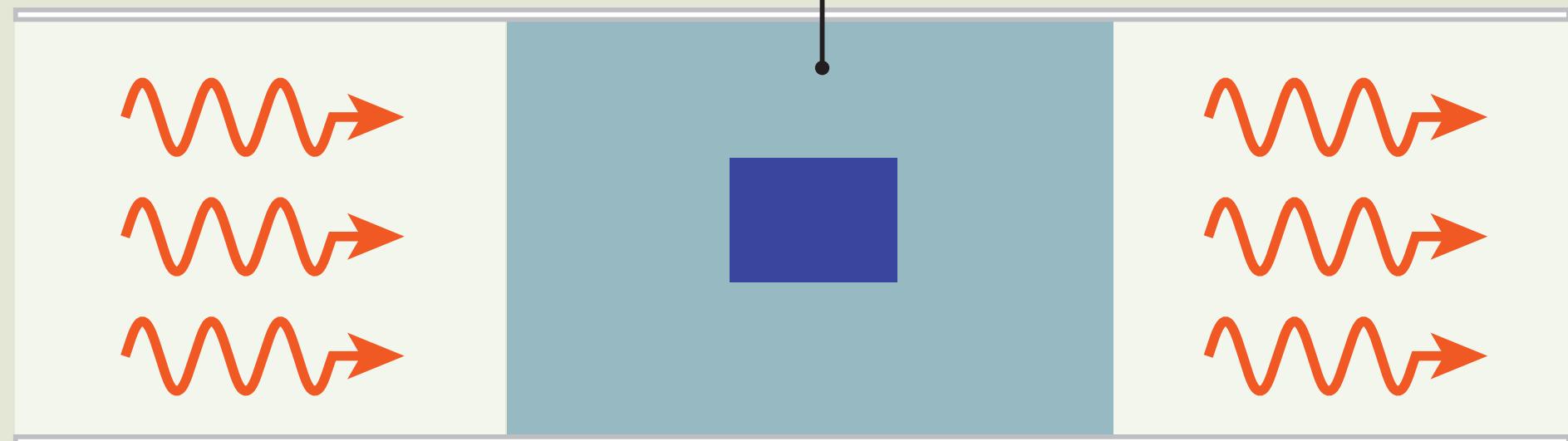
$n = 0$



1 index

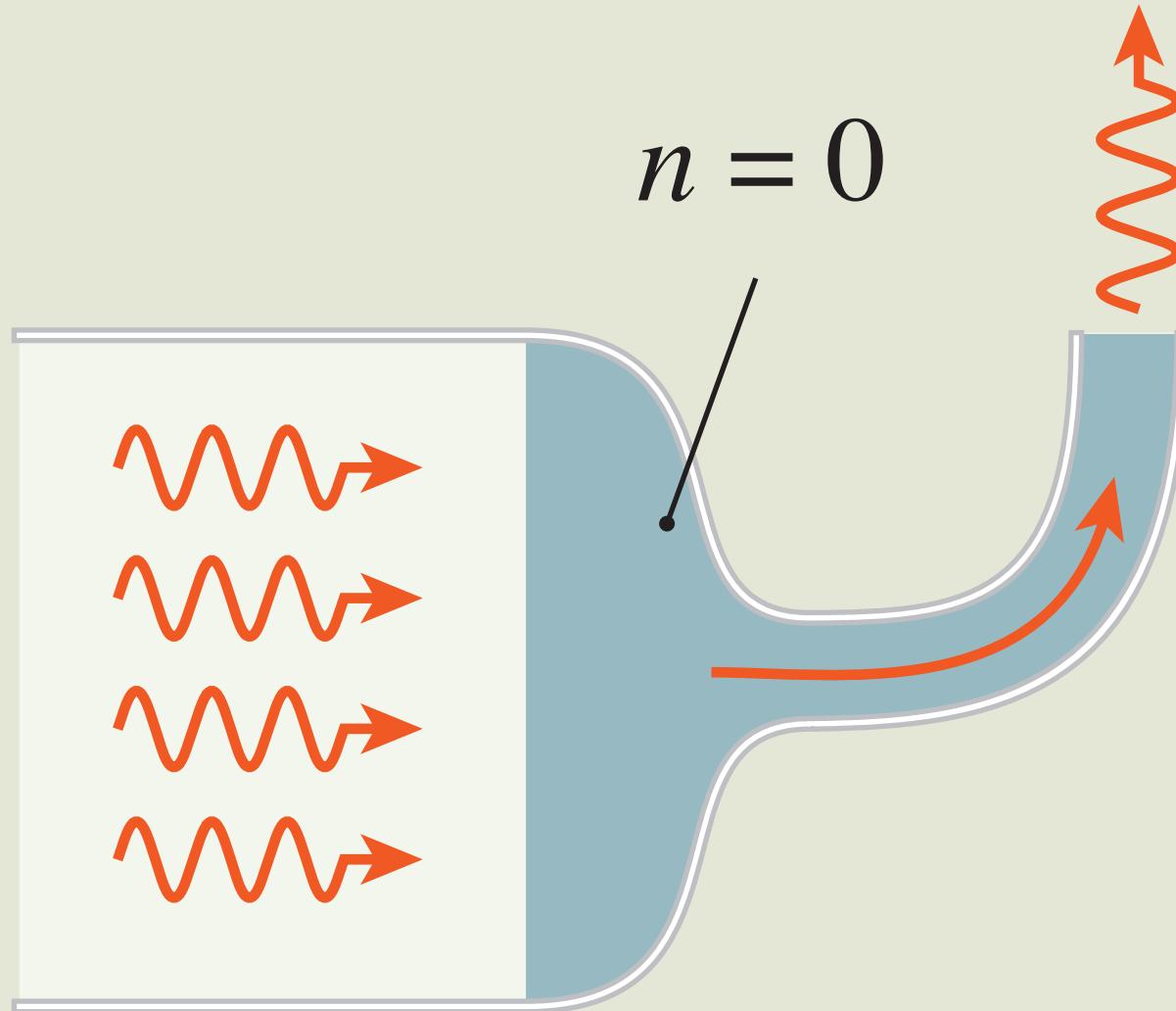
2 zero index

$n = 0$



1 index

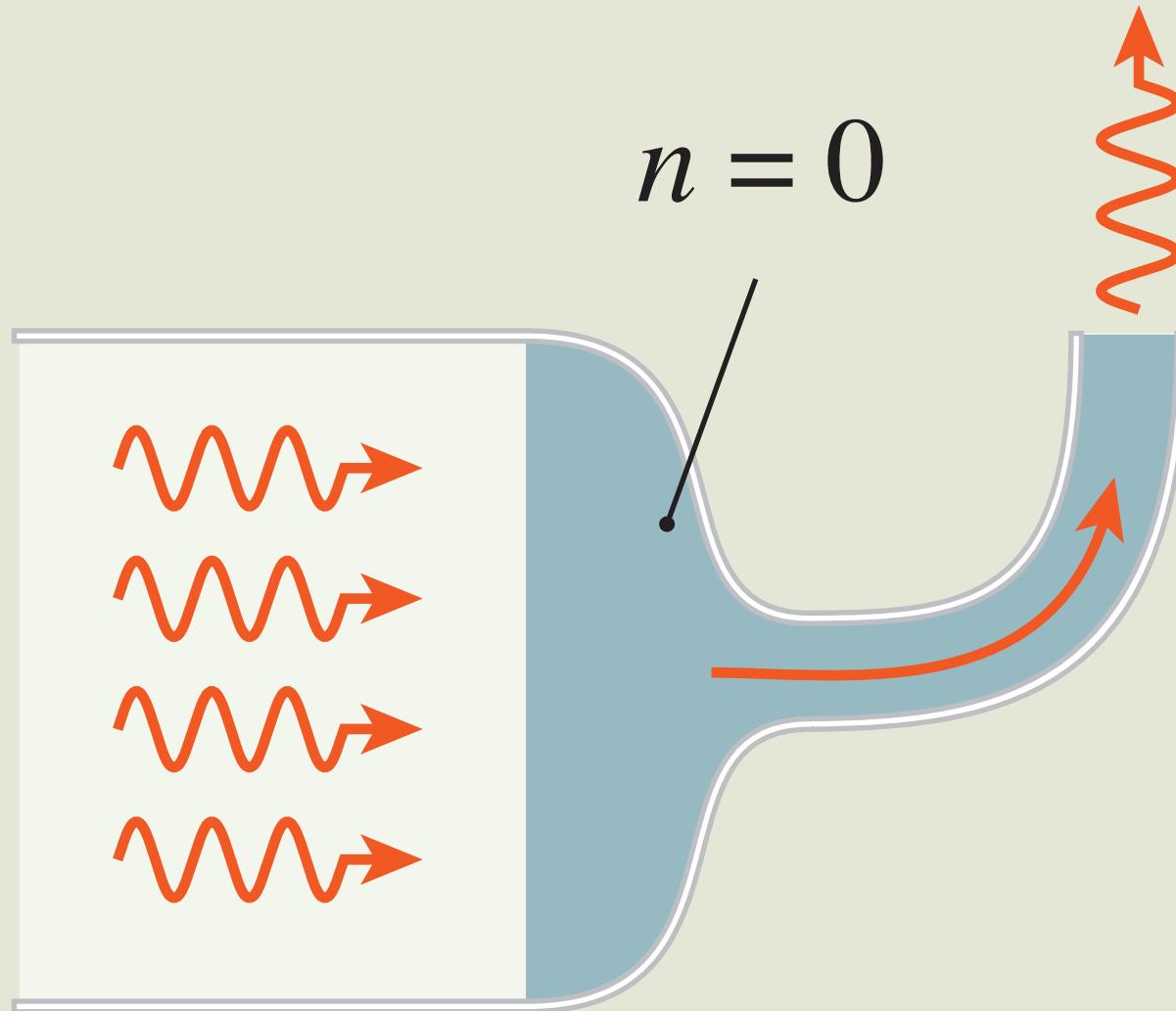
2 zero index



1 index

2 zero index

“tunneling with infinite decay length”



1 index

2 zero index

how?

$$n = \sqrt{\epsilon\mu}$$

but ϵ and μ also determine reflectivity

$$R = \frac{Z-1}{Z+1}$$

1 index

2 zero index

how?

$$n = \sqrt{\epsilon\mu}$$

but ϵ and μ also determine reflectivity

$$R = \frac{Z - 1}{Z + 1}$$

where

$$Z = \sqrt{\frac{\mu}{\epsilon}}$$

1 index

2 zero index

how?

$$\varepsilon, \mu \rightarrow 0$$

$$n = \sqrt{\varepsilon\mu} \rightarrow 0$$

but ε and μ also determine reflectivity

$$R = \frac{Z - 1}{Z + 1}$$

where

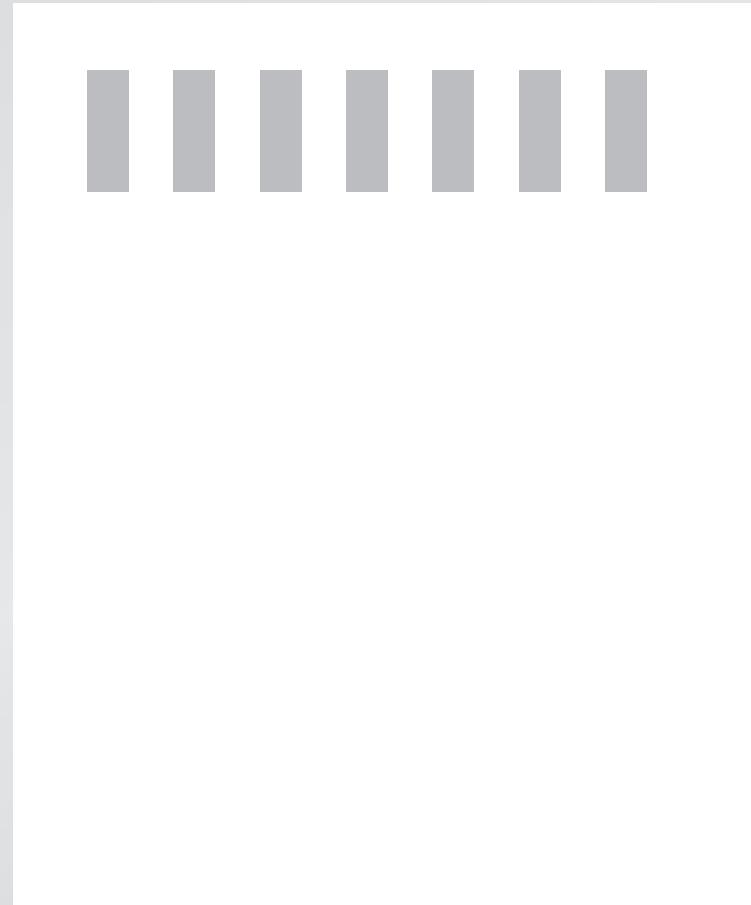
$$Z = \sqrt{\frac{\mu}{\varepsilon}} \quad \text{finite!}$$

1 index

2 zero index

Engineering a magnetic response

use array of dielectric rods

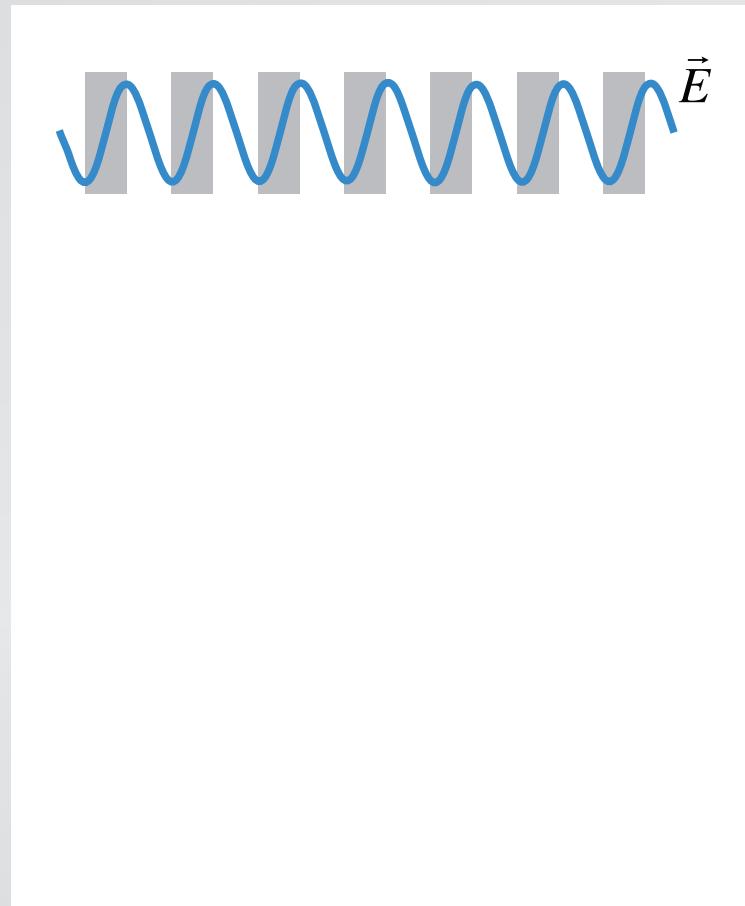


1 index

2 zero index

Engineering a magnetic response

incident electromagnetic wave ($\lambda_{\text{eff}} \approx a$)

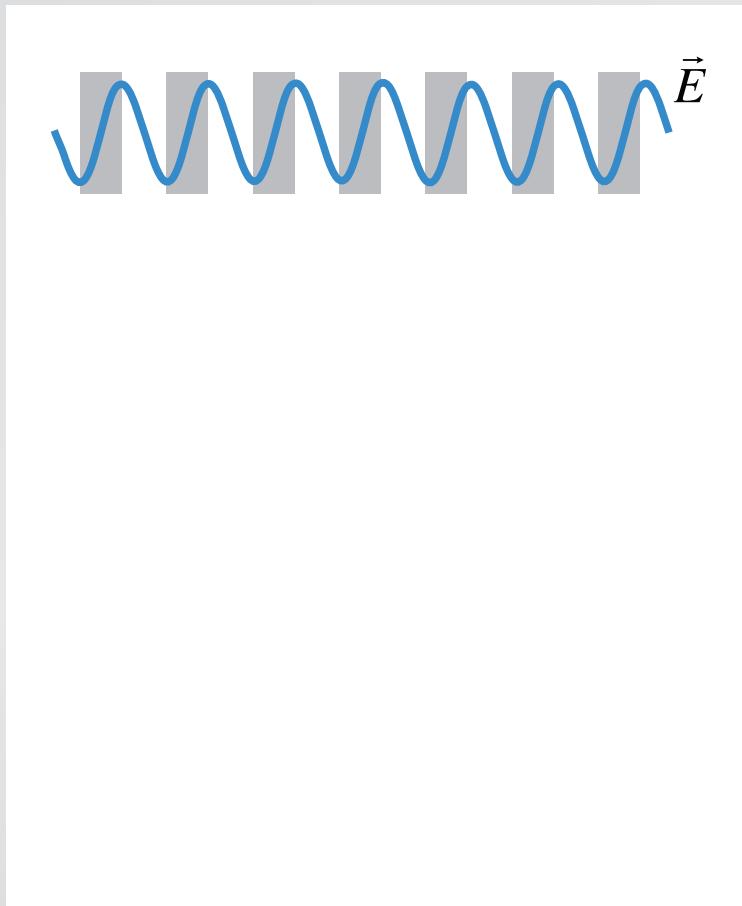


1 index

2 zero index

Engineering a magnetic response

produces an electric response...

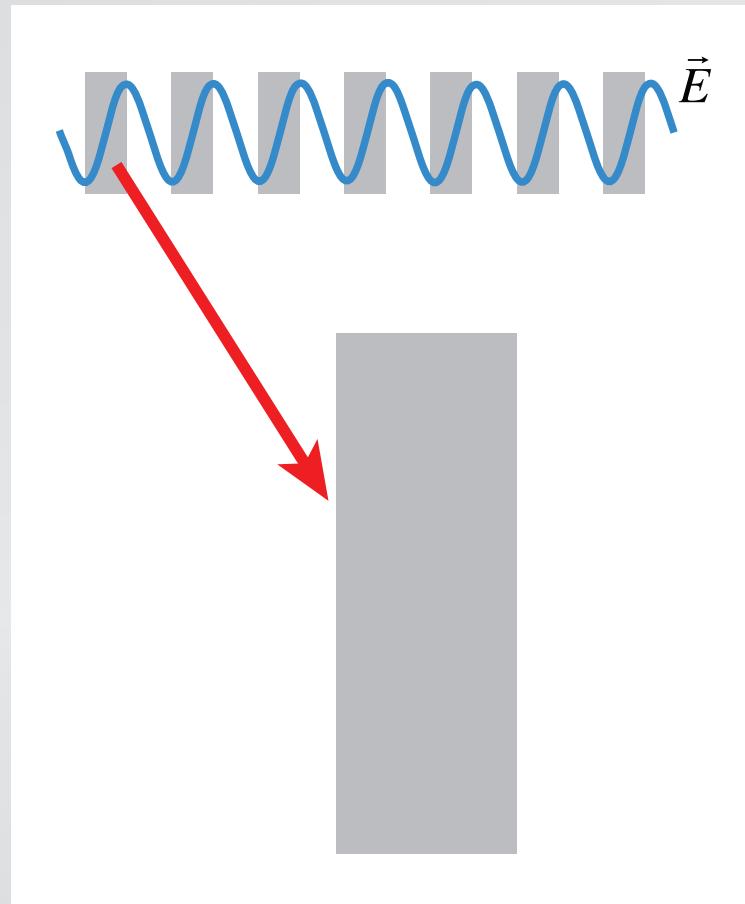


1 index

2 zero index

Engineering a magnetic response

... but different electric fields front and back...

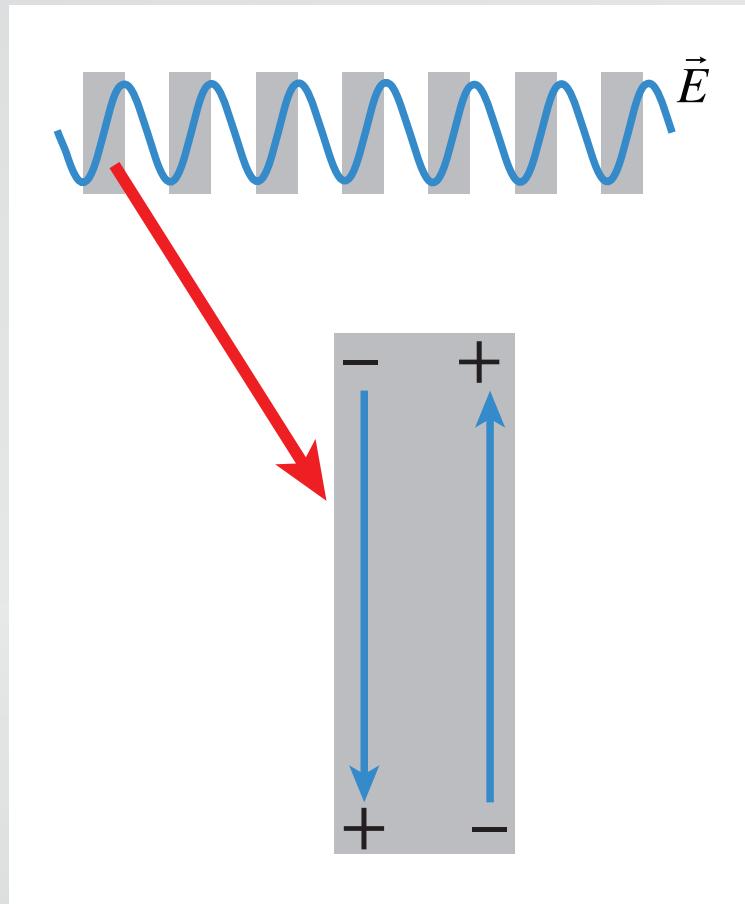


1 index

2 zero index

Engineering a magnetic response

...induce different polarizations on opposite sides...

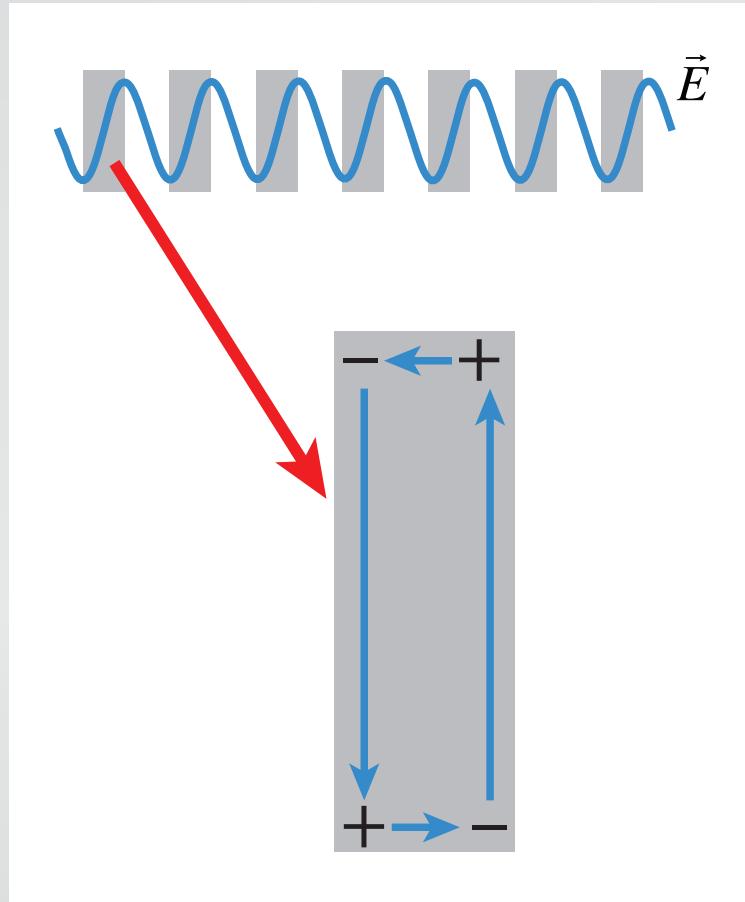


1 index

2 zero index

Engineering a magnetic response

...causing a current loop...

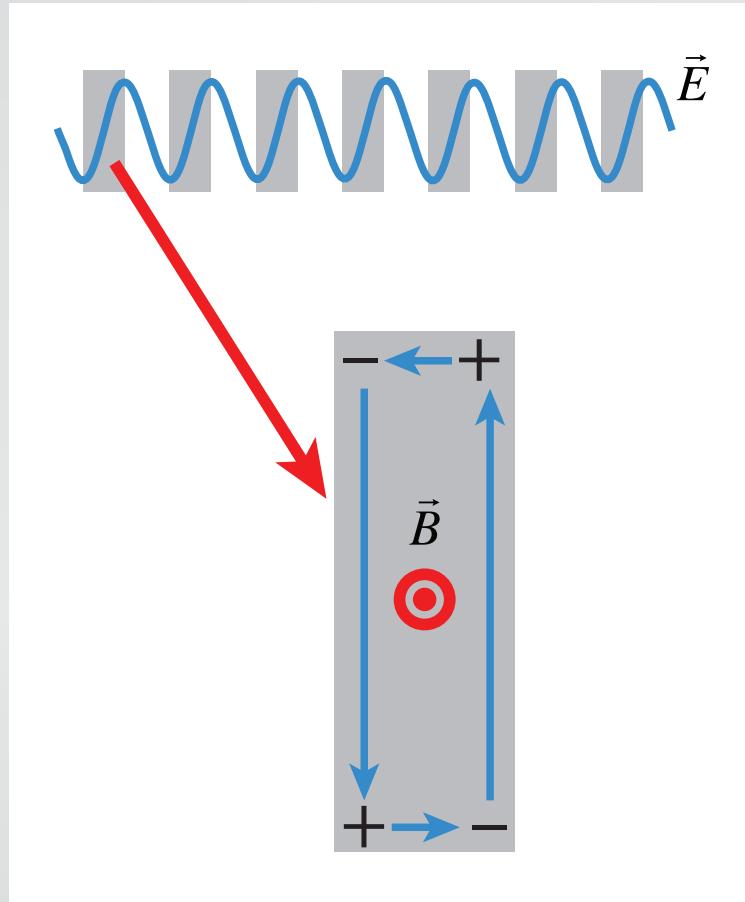


1 index

2 zero index

Engineering a magnetic response

...which, in turn, produces an induced magnetic field

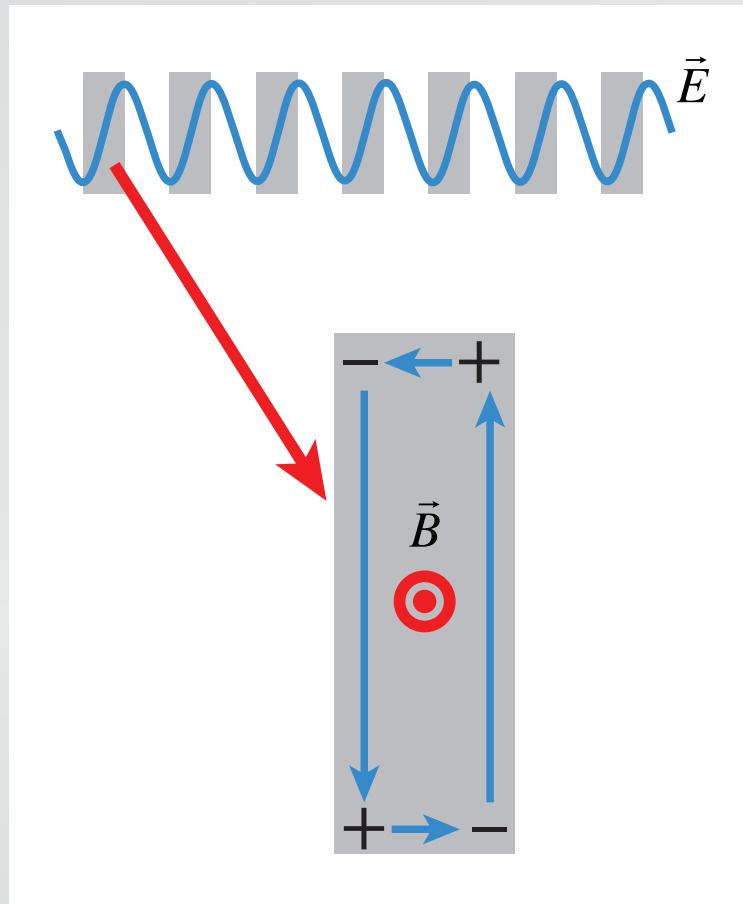


1 index

2 zero index

Engineering a magnetic response

adjust design so electrical and magnetic resonances coincide



1 index

2 zero index

Engineering a magnetic response

adjustable parameters



1 index

2 zero index

Engineering a magnetic response

adjustable parameters

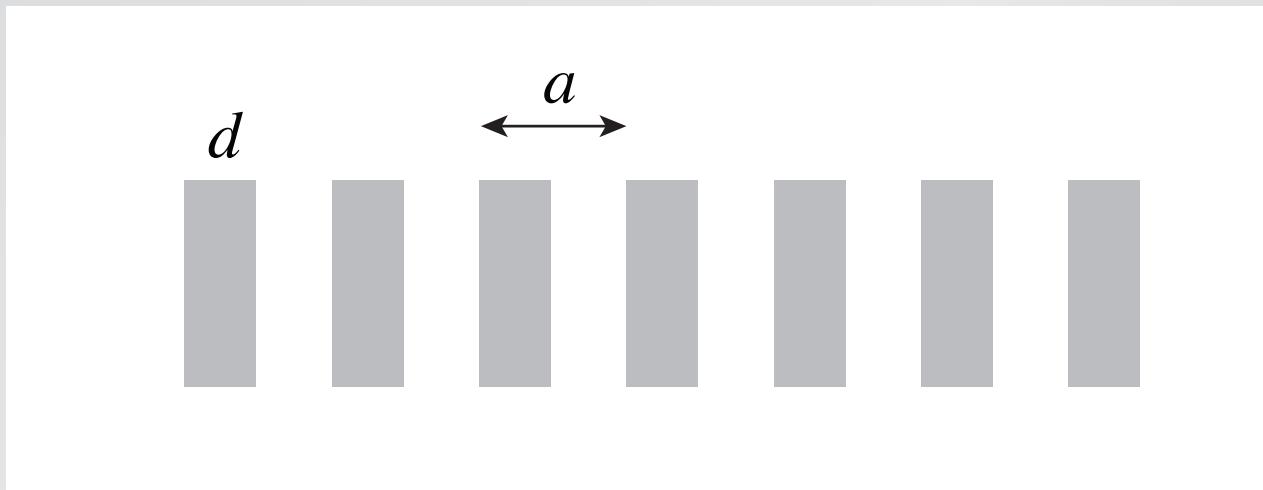


1 index

2 zero index

Engineering a magnetic response

adjustable parameters

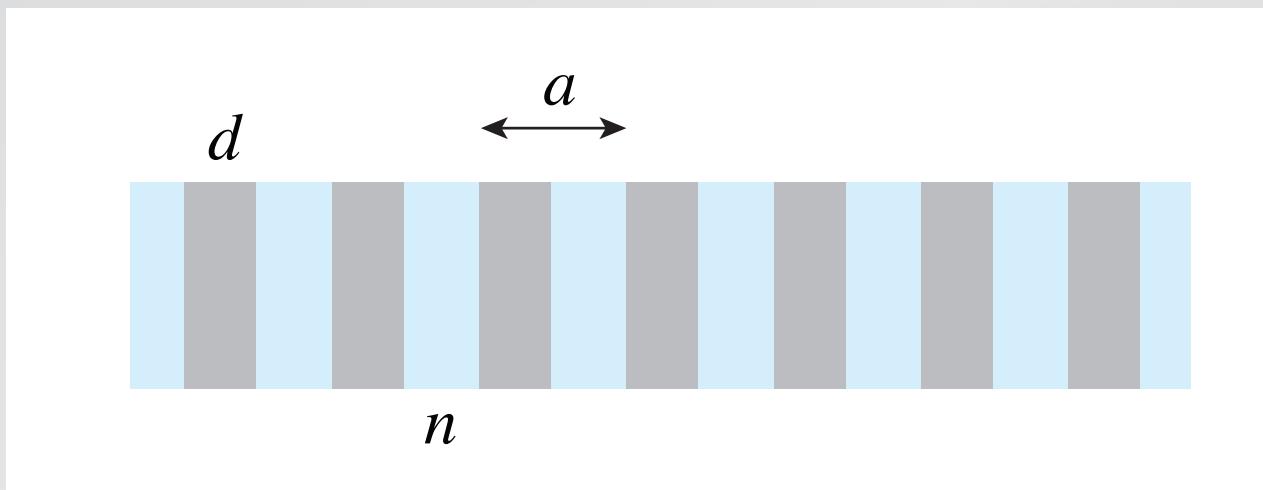


1 index

2 zero index

Engineering a magnetic response

adjustable parameters

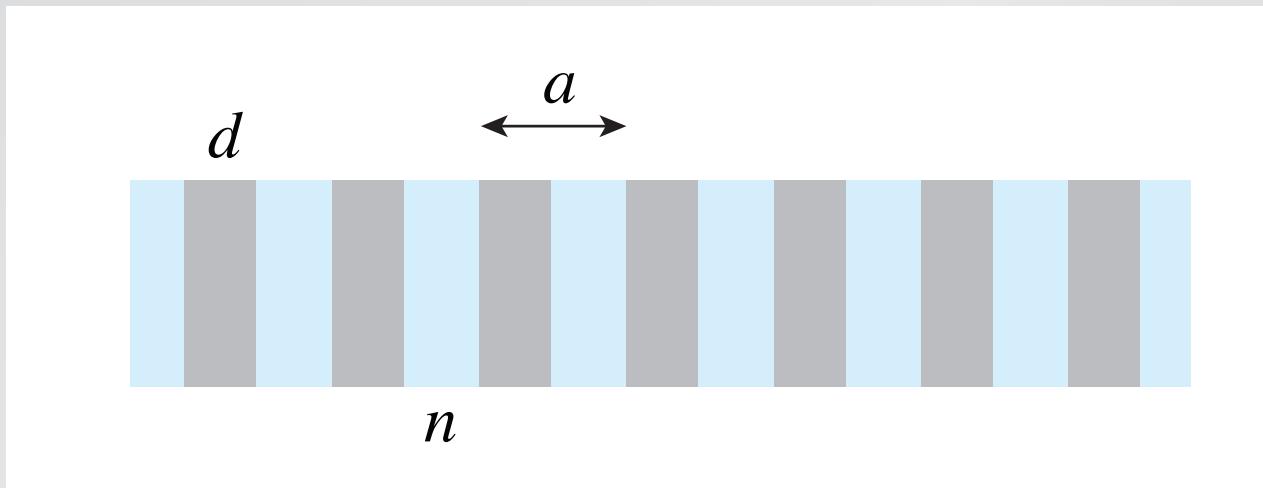


1 index

2 zero index

Engineering a magnetic response

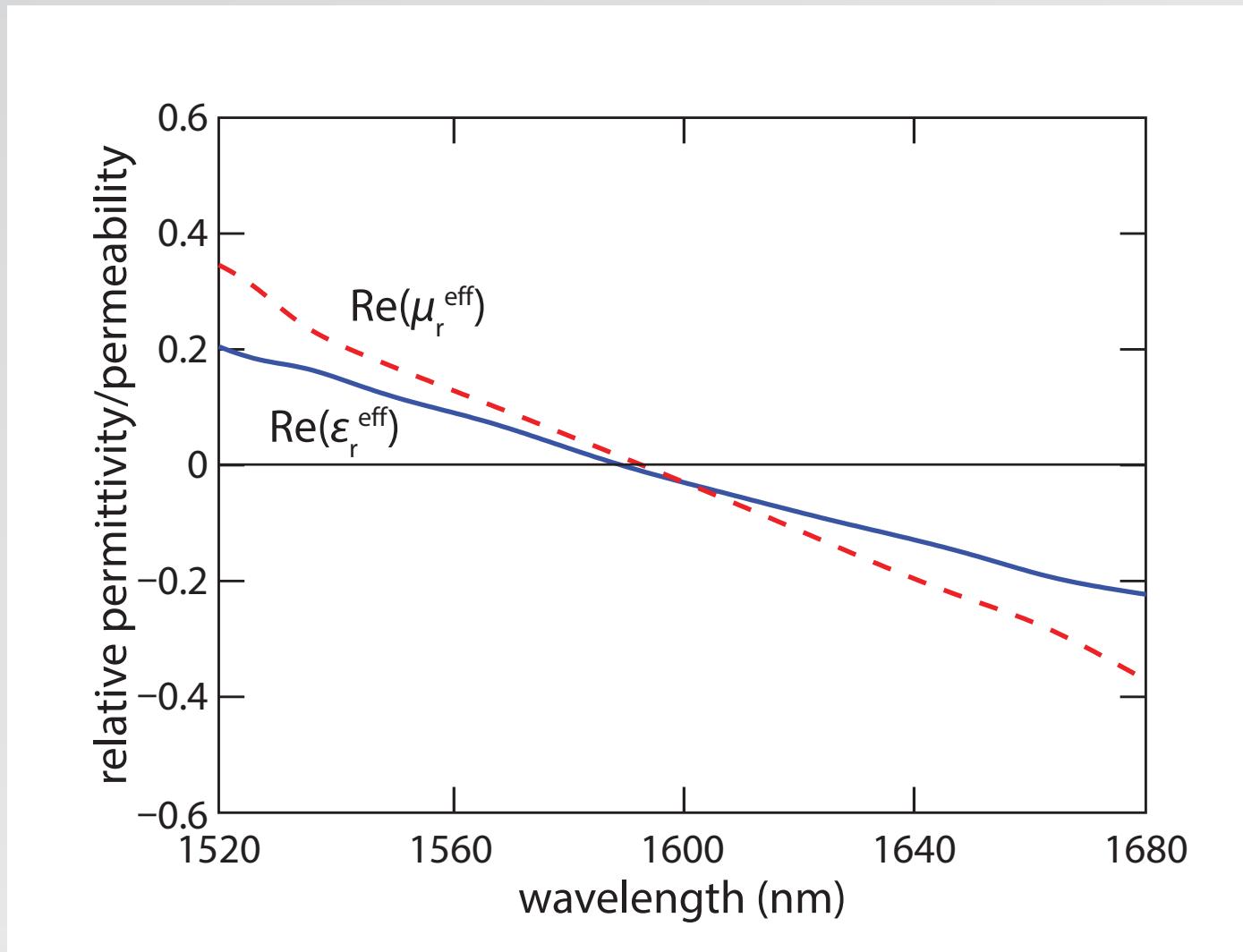
adjustable parameters



$$d = 422 \text{ nm}, \quad a = 690 \text{ nm}, \quad n = 1.57 \text{ (SU8)}$$

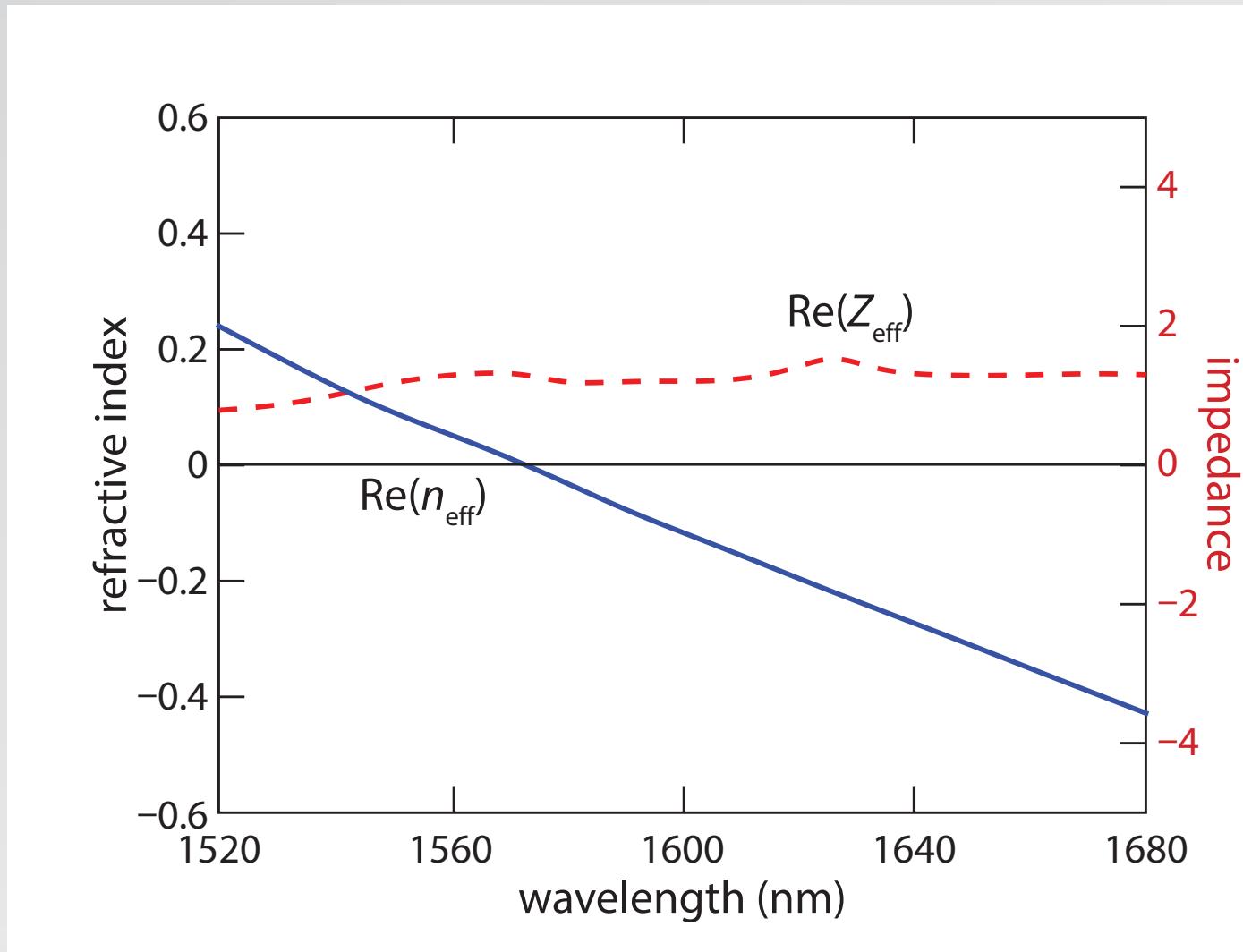
1 index

2 zero index



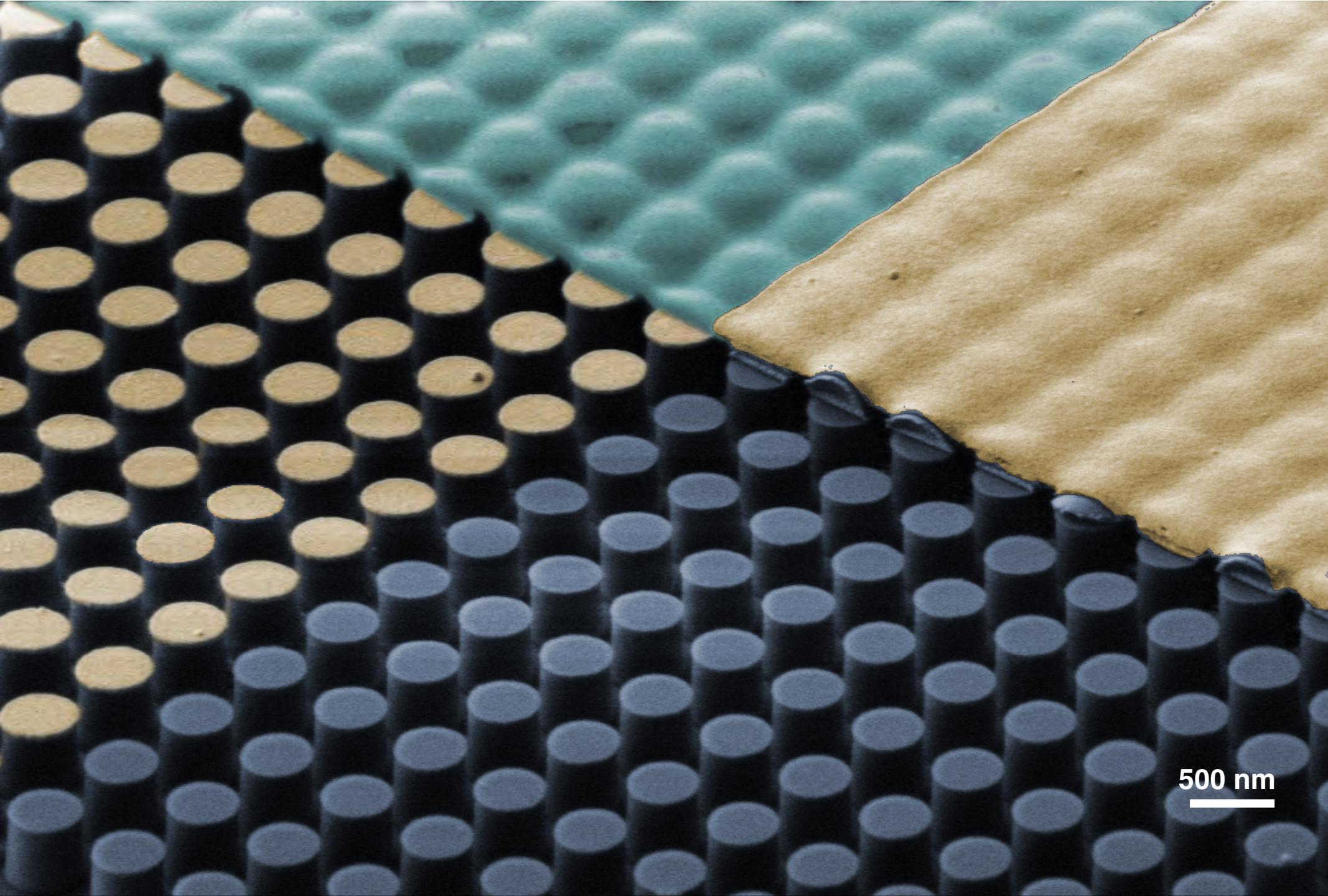
1 index

2 zero index



1 index

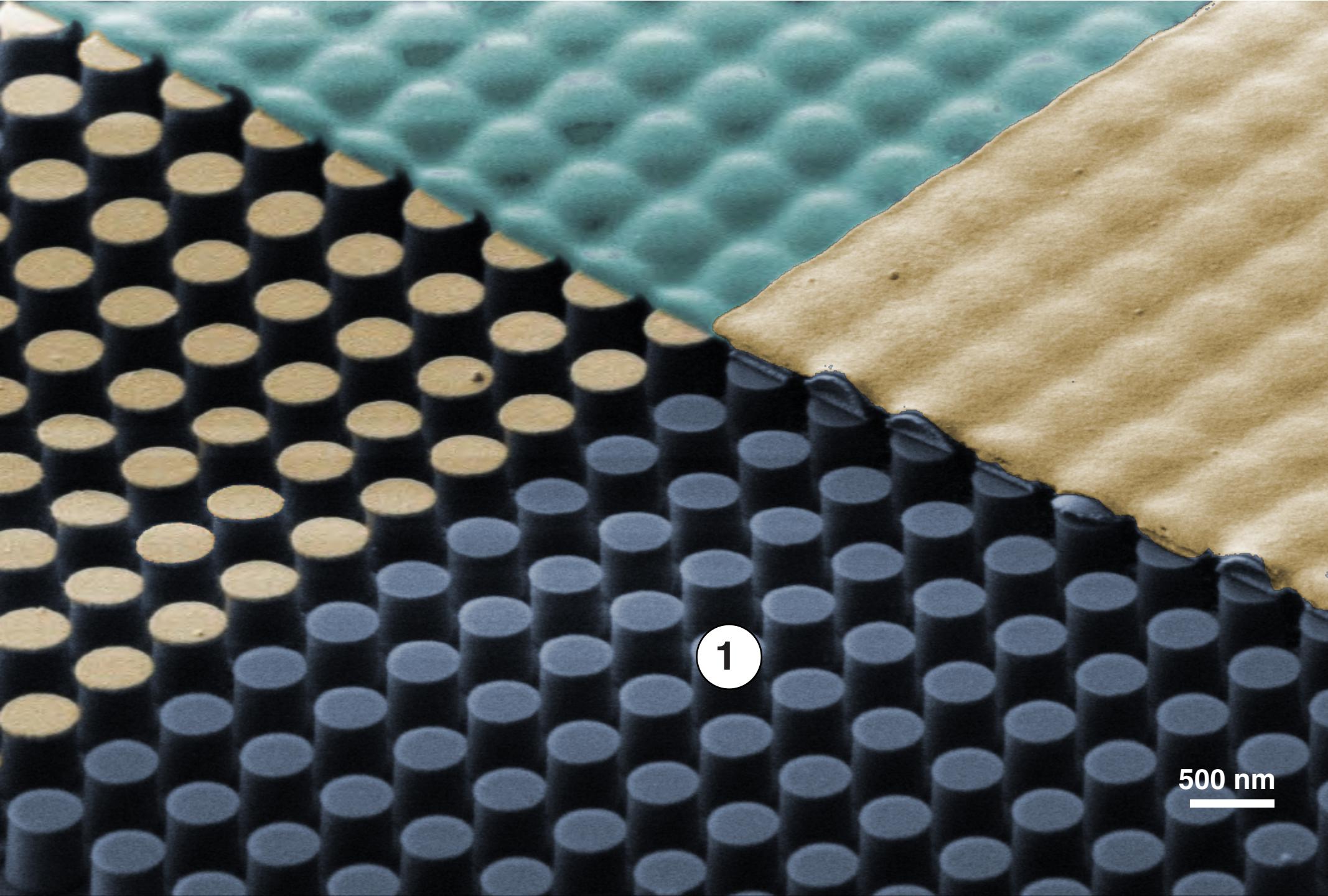
2 zero index



1 index

2 zero index

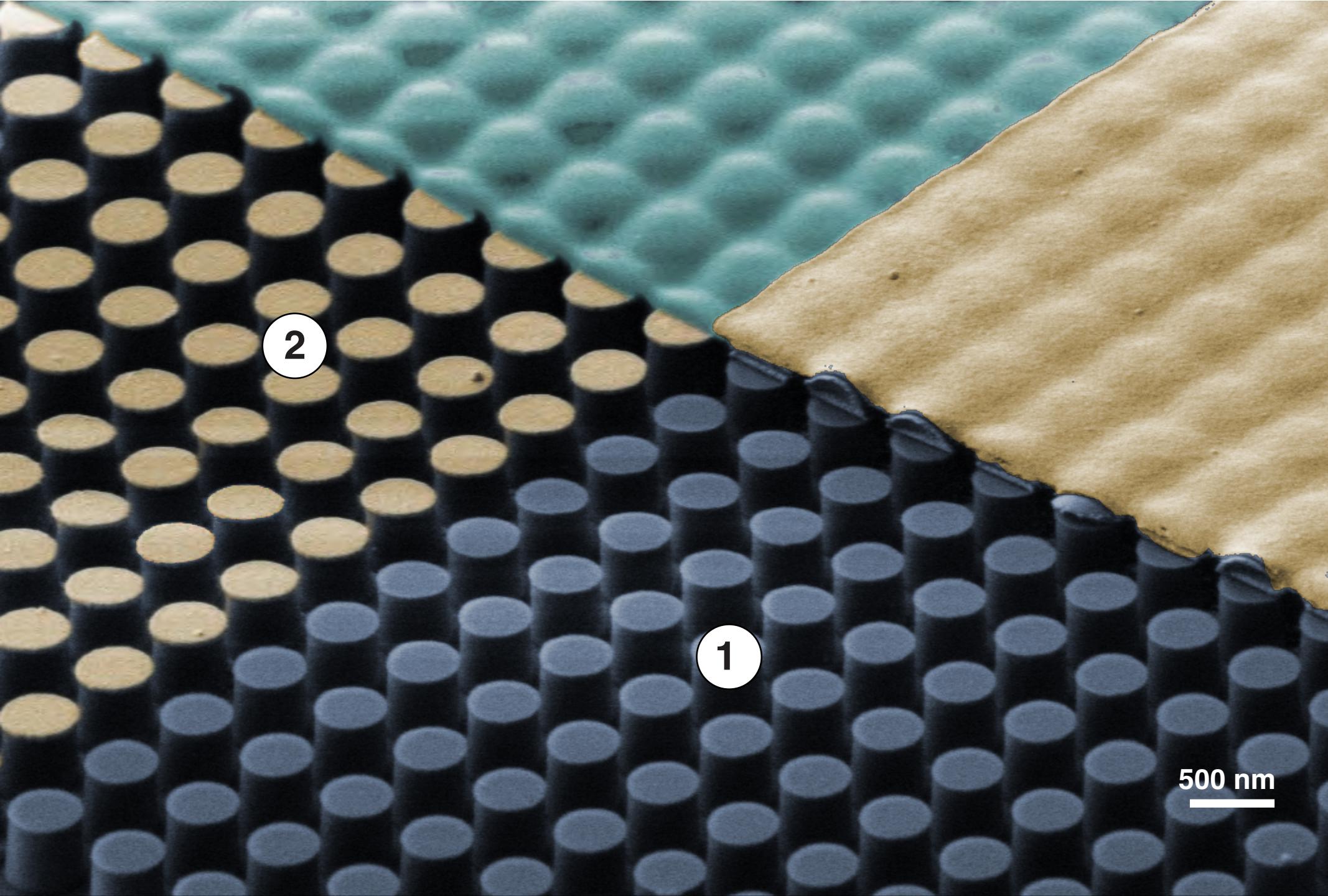
500 nm



1 index

2 zero index

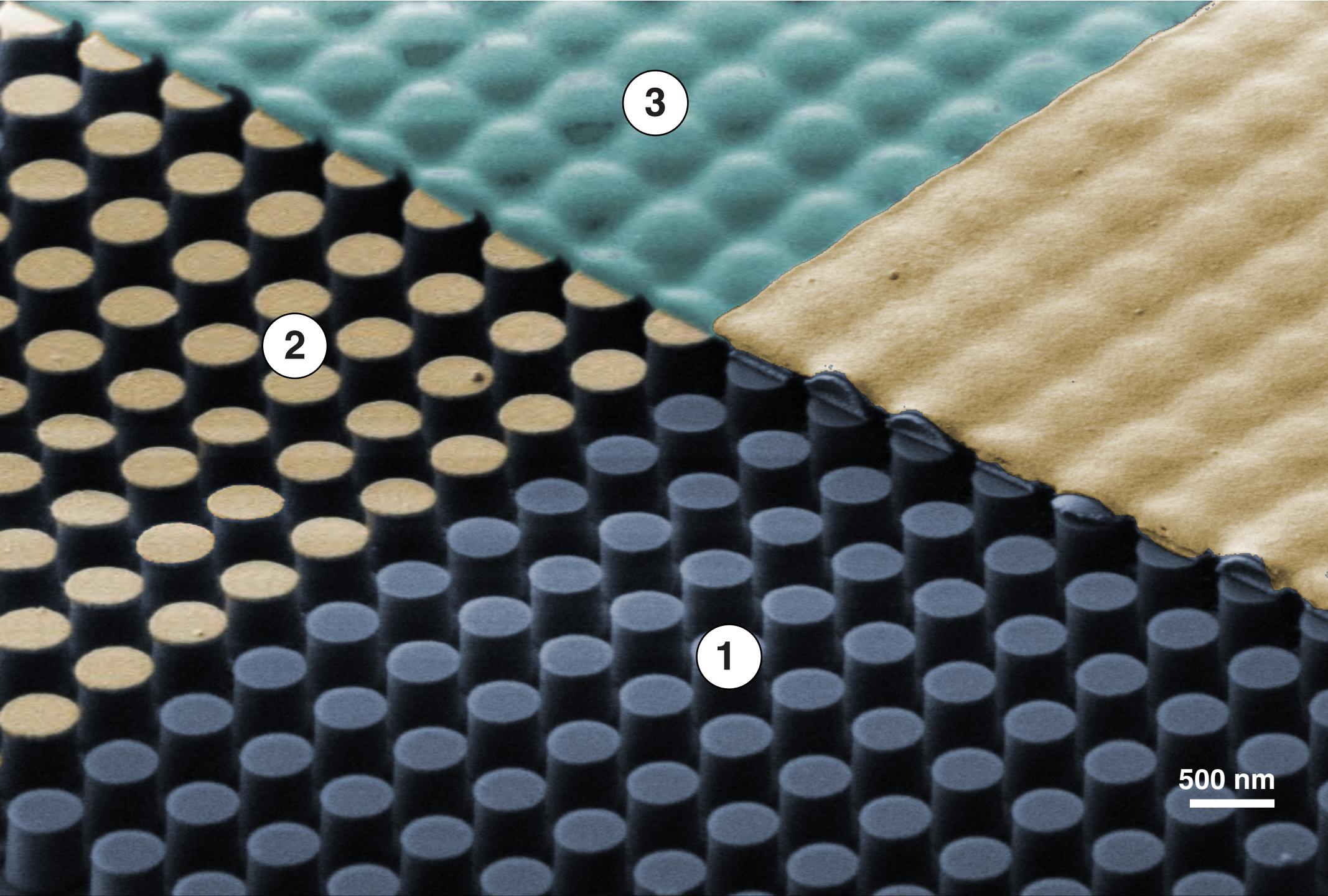
500 nm



1 index

2 zero index

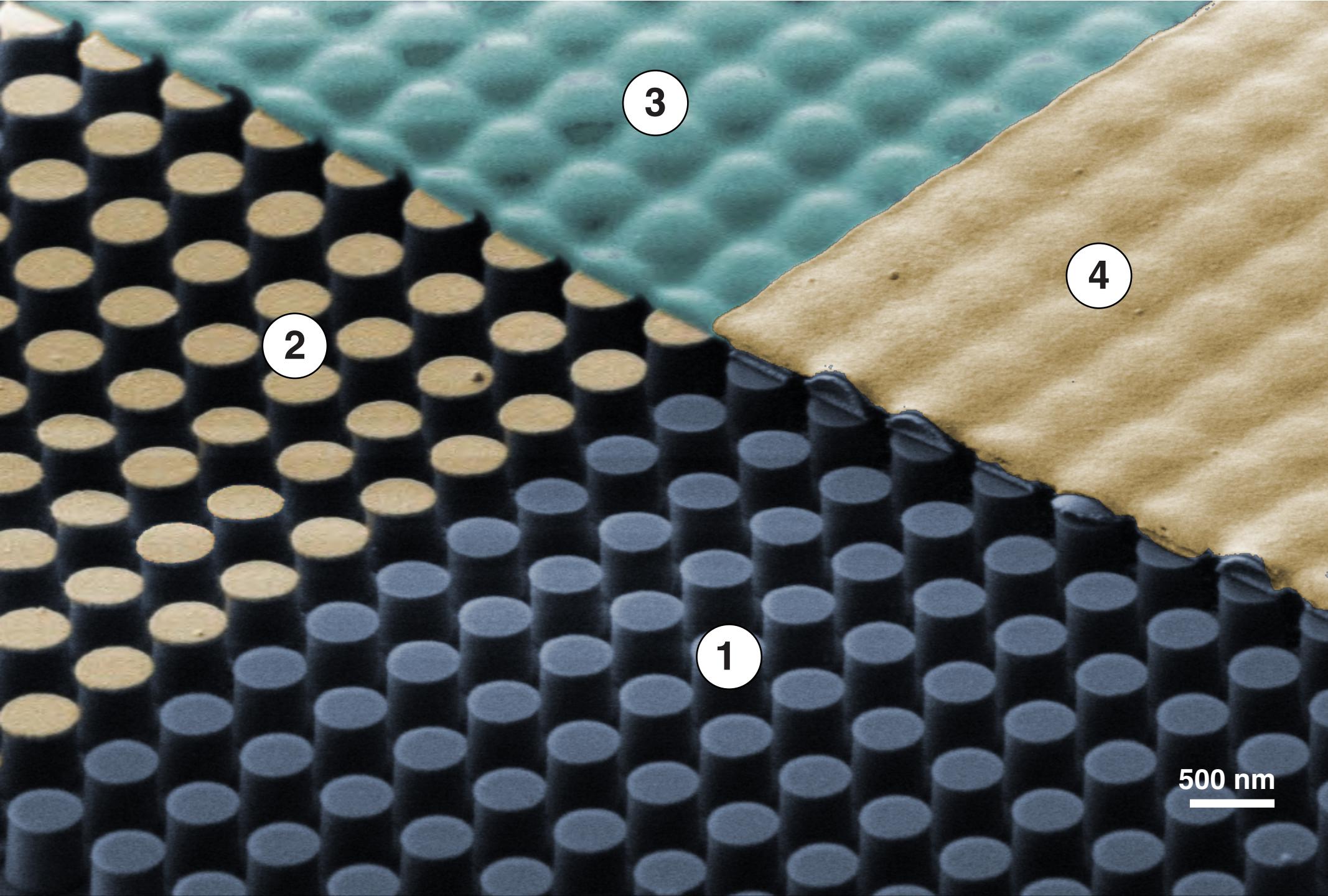
500 nm



1 index

2 zero index

500 nm



1 index

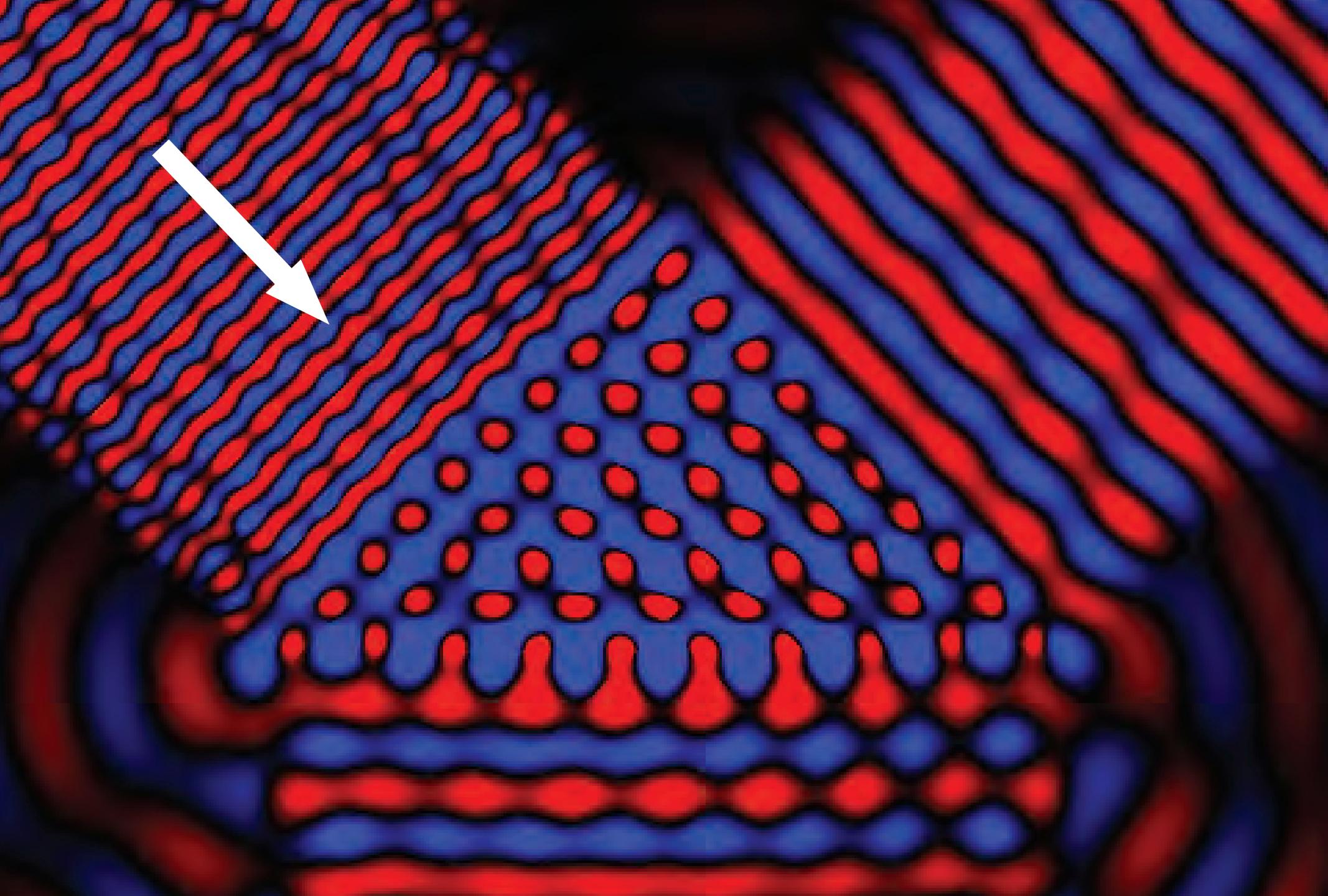
2 zero index

500 nm

Can make this in any shape!

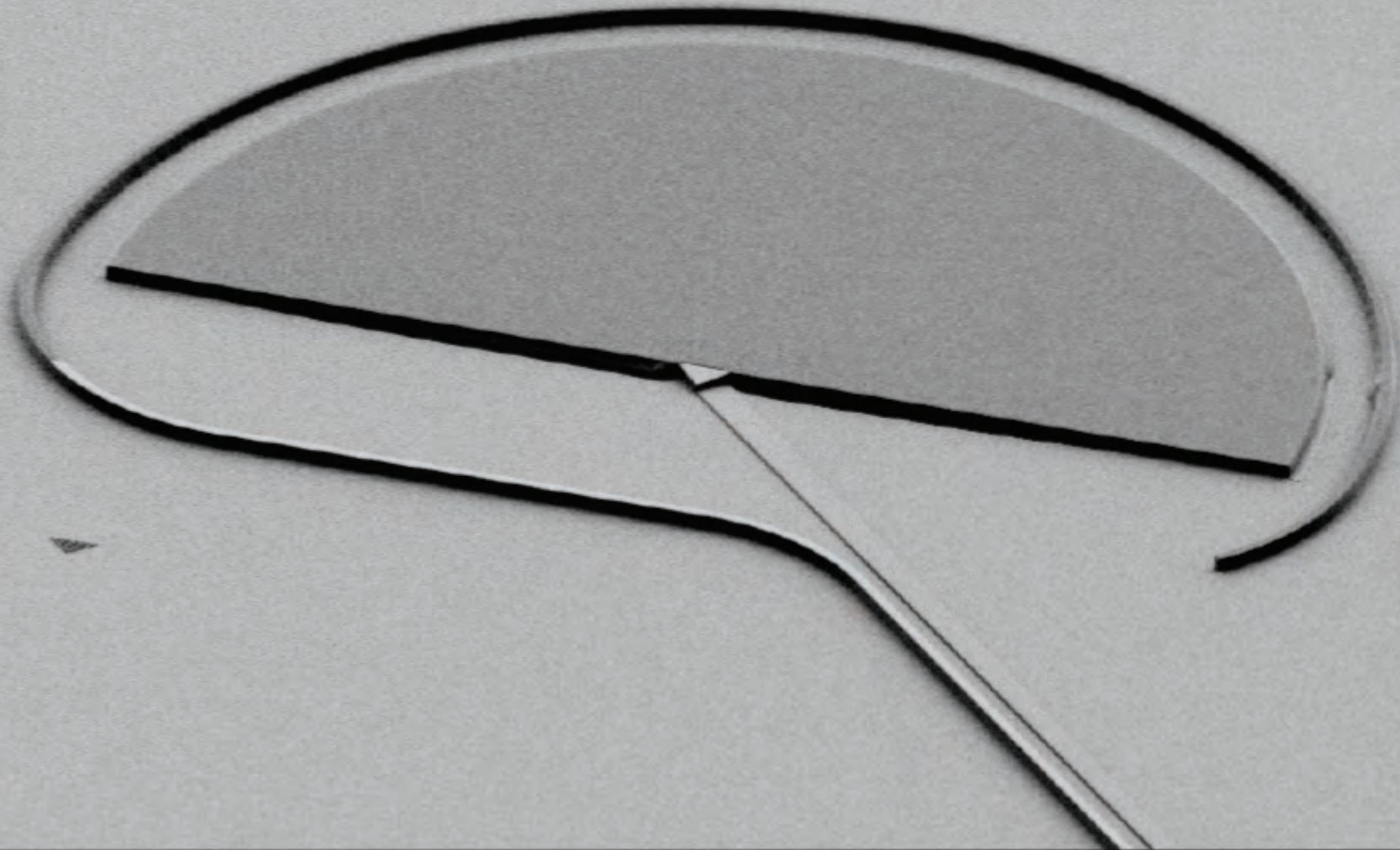
1 index

2 zero index



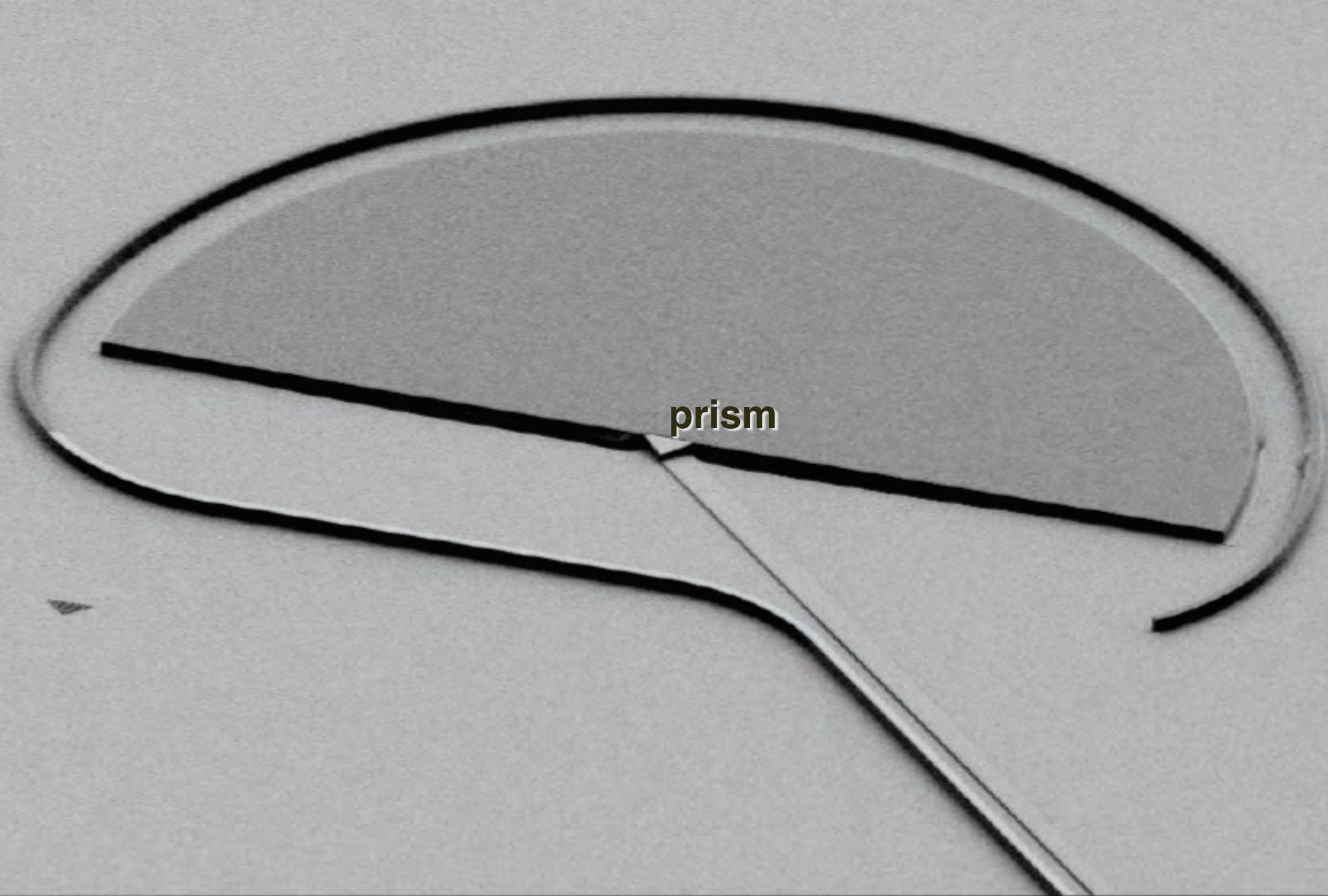
1 index

2 zero index



1 index

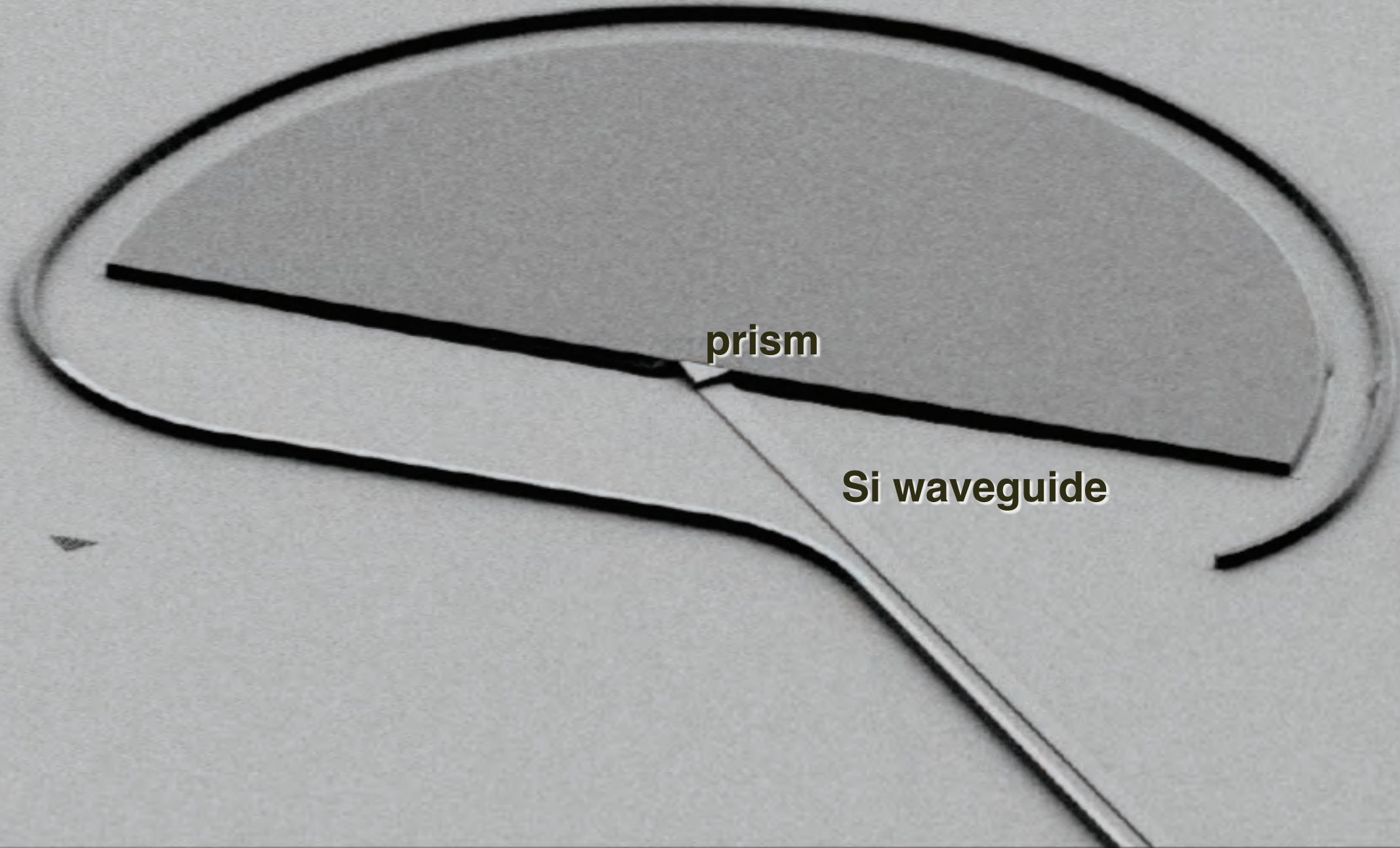
2 zero index



prism

1 index

2 zero index

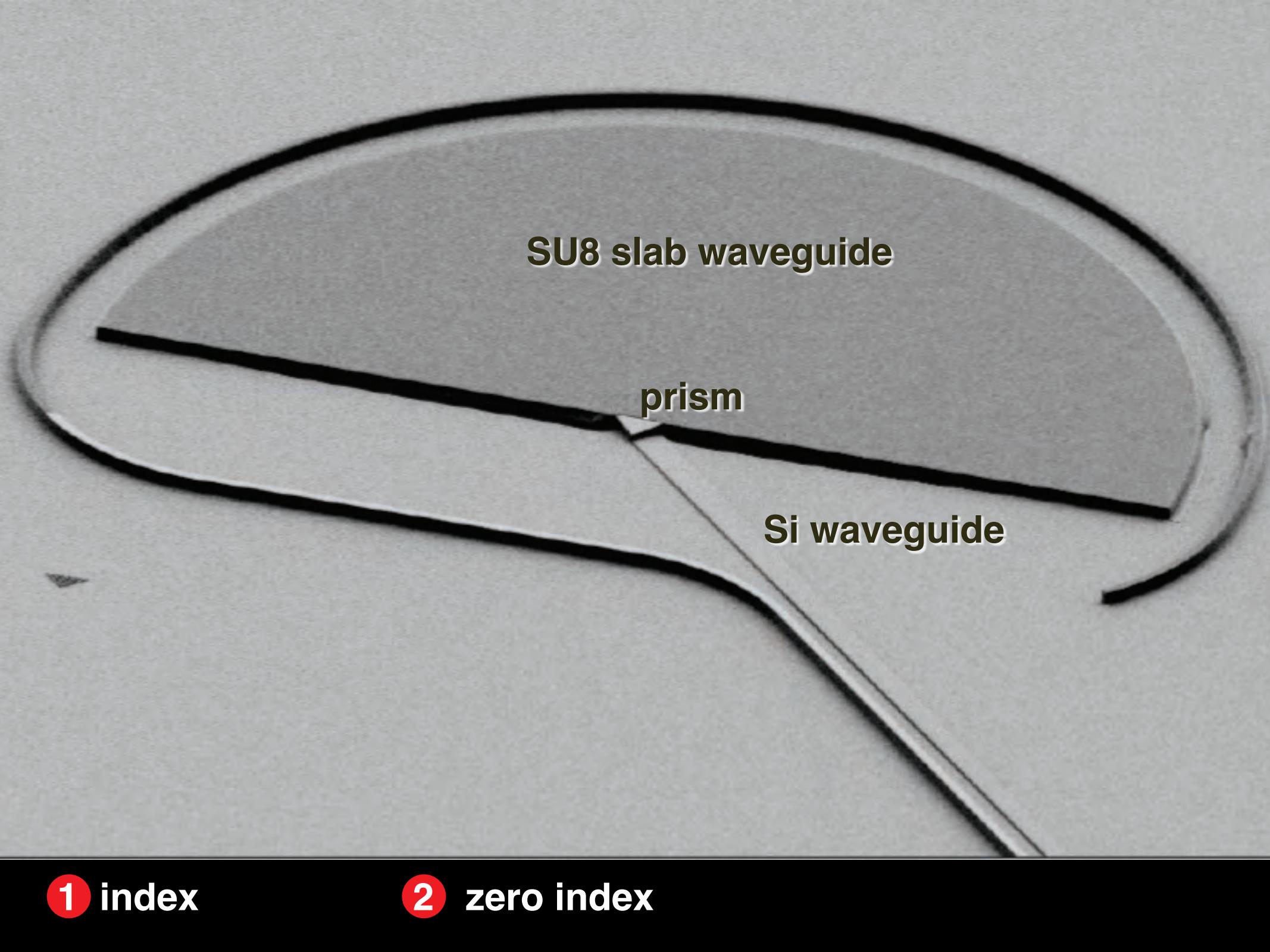


prism

Si waveguide

1 index

2 zero index

A scanning electron micrograph (SEM) showing a waveguide structure. A large, thick, curved waveguide is labeled "SU8 slab waveguide". A smaller, thinner waveguide that splits off from the main one is labeled "Si waveguide". The two waveguides meet at a junction point. The word "prism" is written vertically along the main waveguide curve.

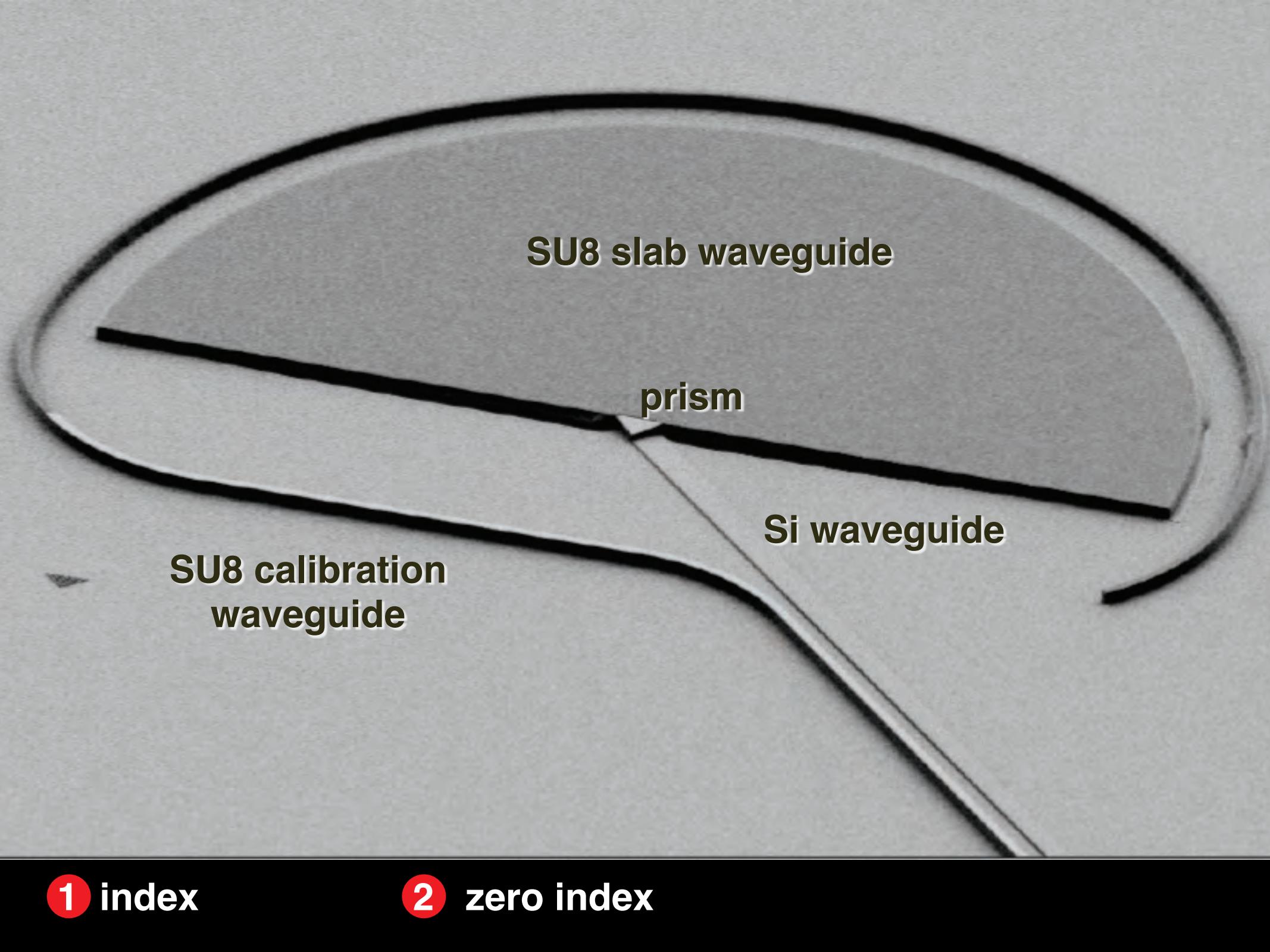
SU8 slab waveguide

prism

Si waveguide

1 index

2 zero index

A scanning electron micrograph (SEM) showing a cross-section of a waveguide structure. The structure consists of a thick, curved SU8 slab waveguide at the top, which tapers down to a thinner Si waveguide. A small triangular prism is positioned between the two waveguides. A horizontal calibration bar is visible on the left side of the image.

SU8 slab waveguide

prism

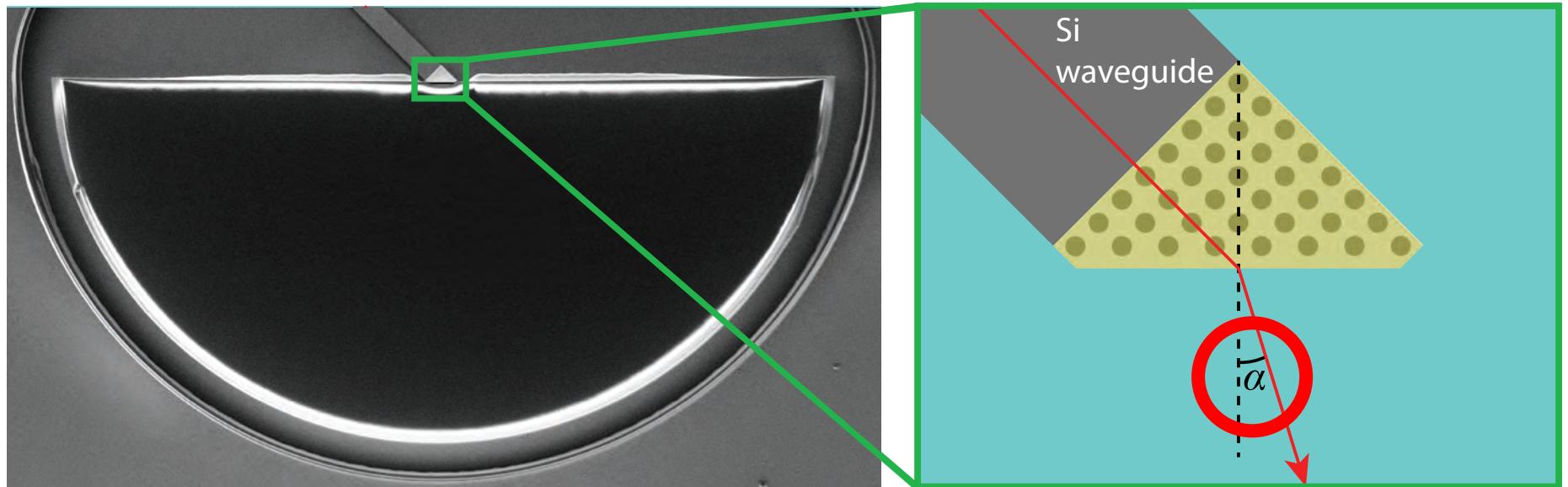
**SU8 calibration
waveguide**

Si waveguide

1 index

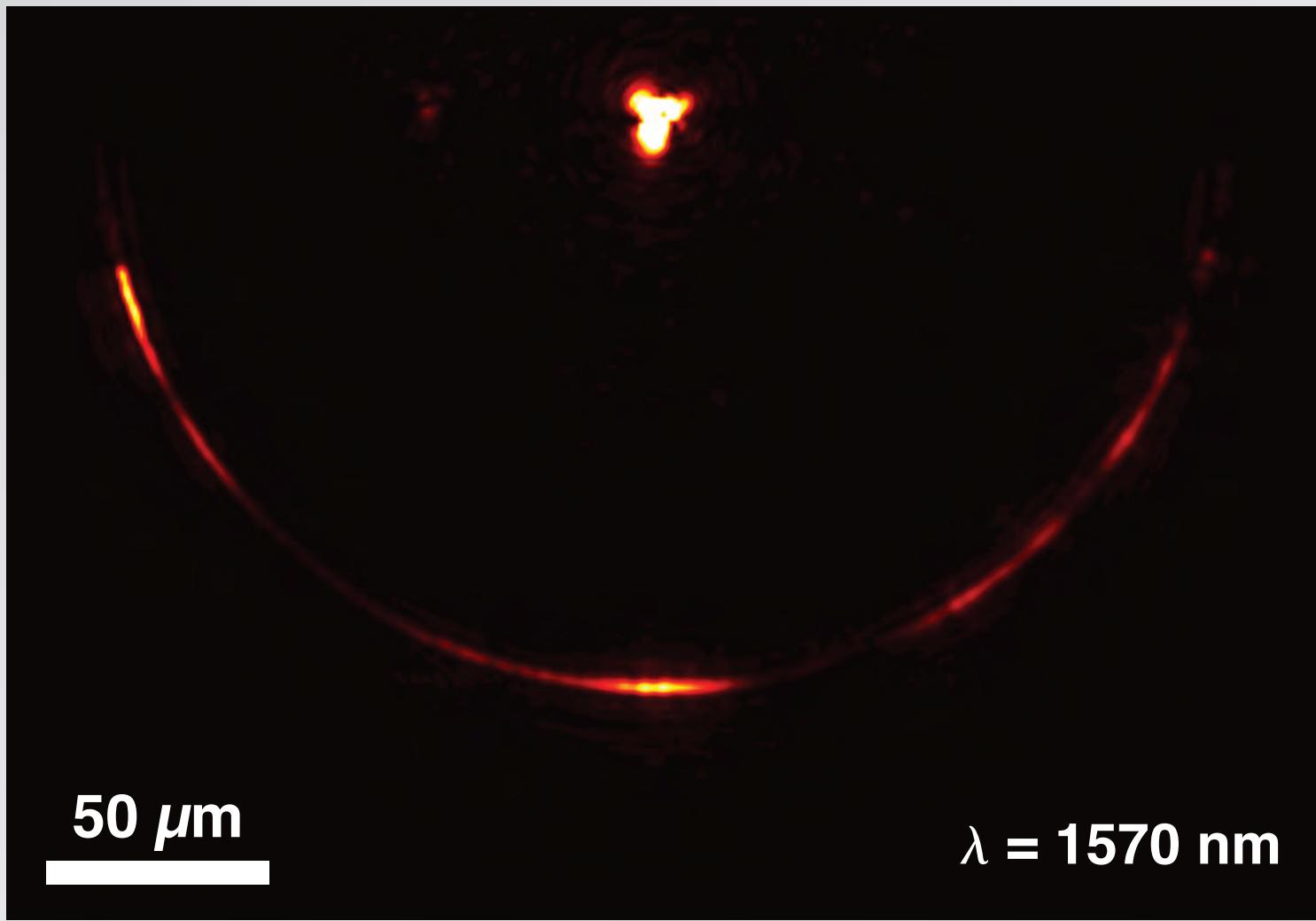
2 zero index

On-chip zero-index prism



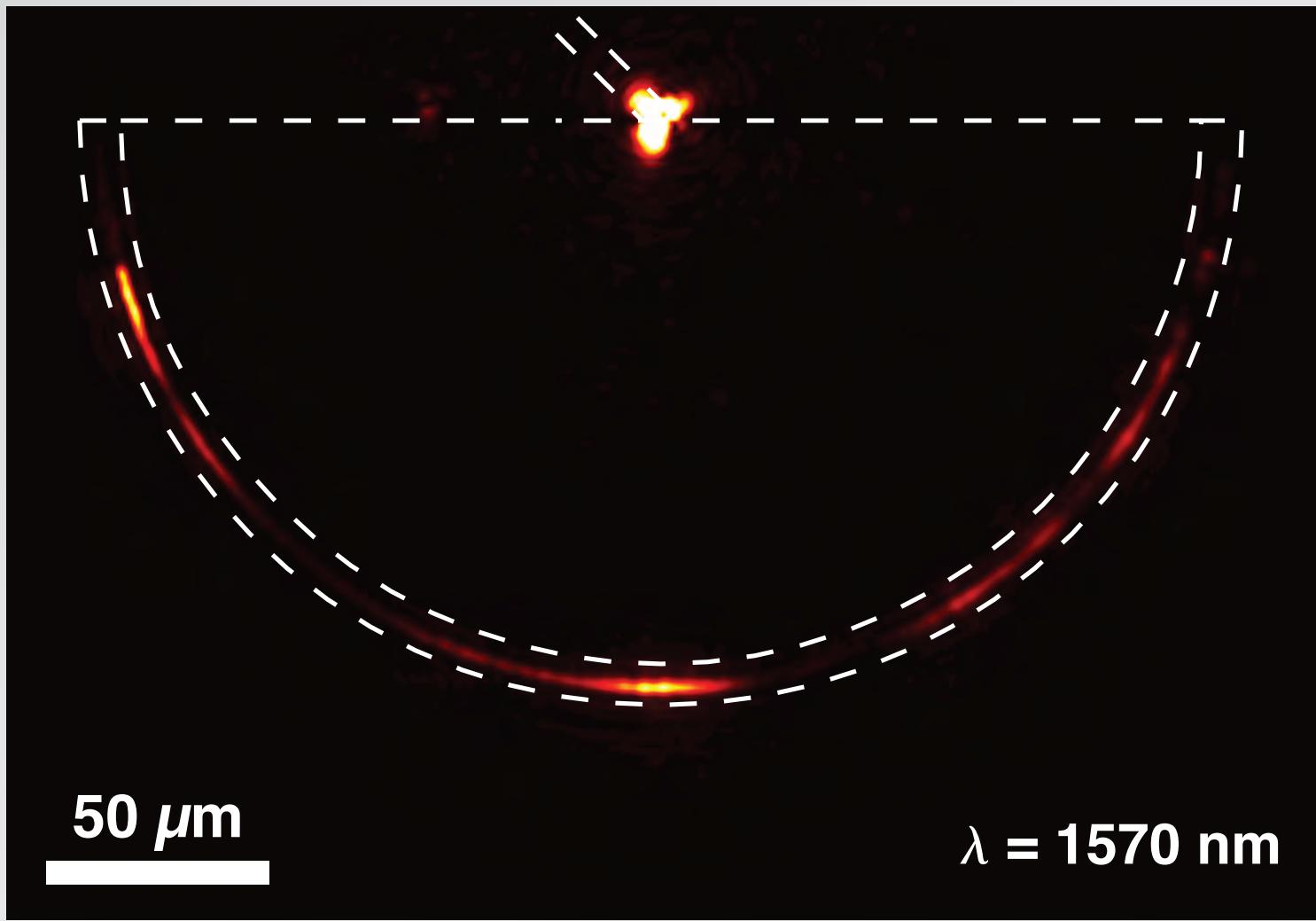
1 index

2 zero index



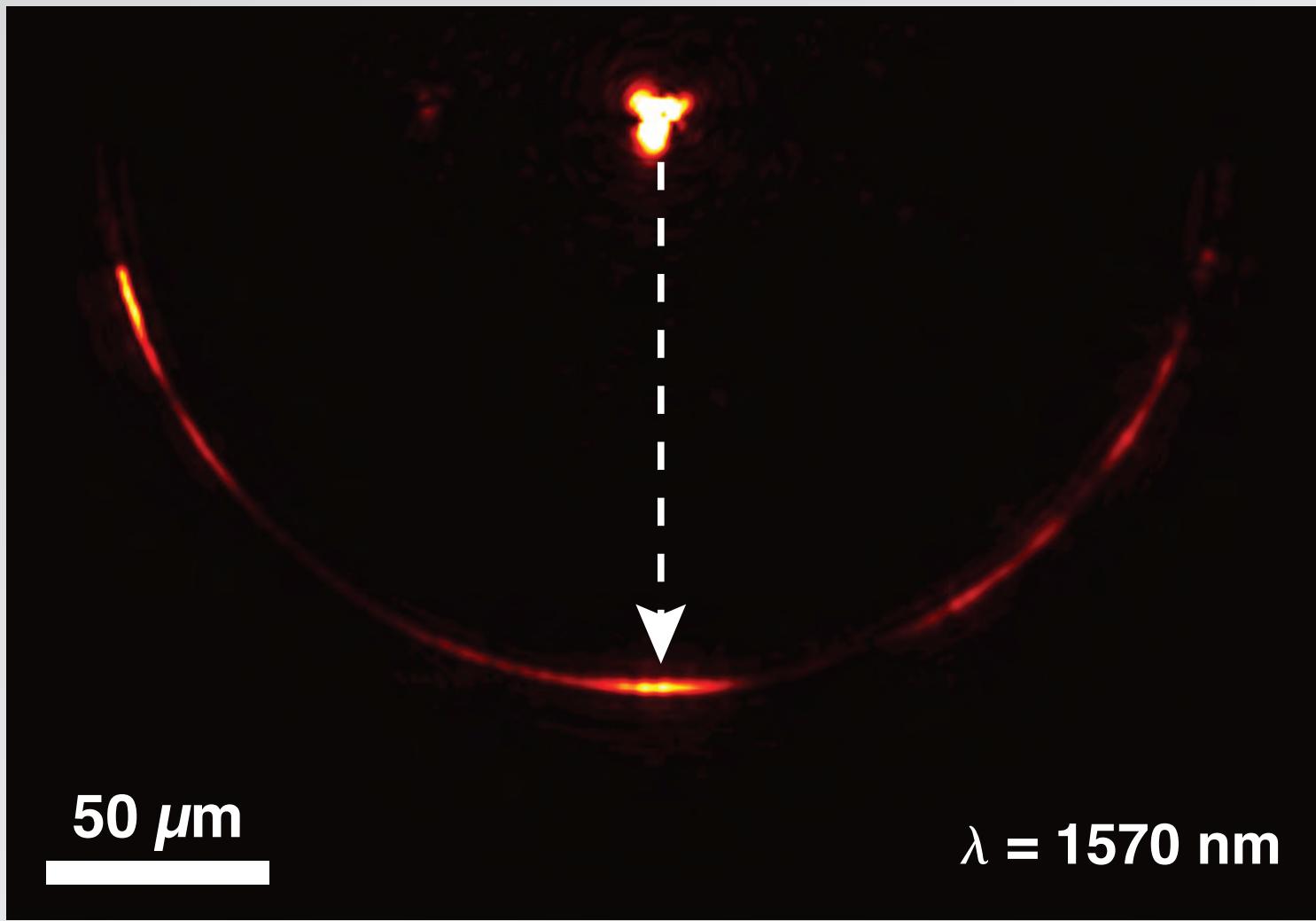
1 index

2 zero index



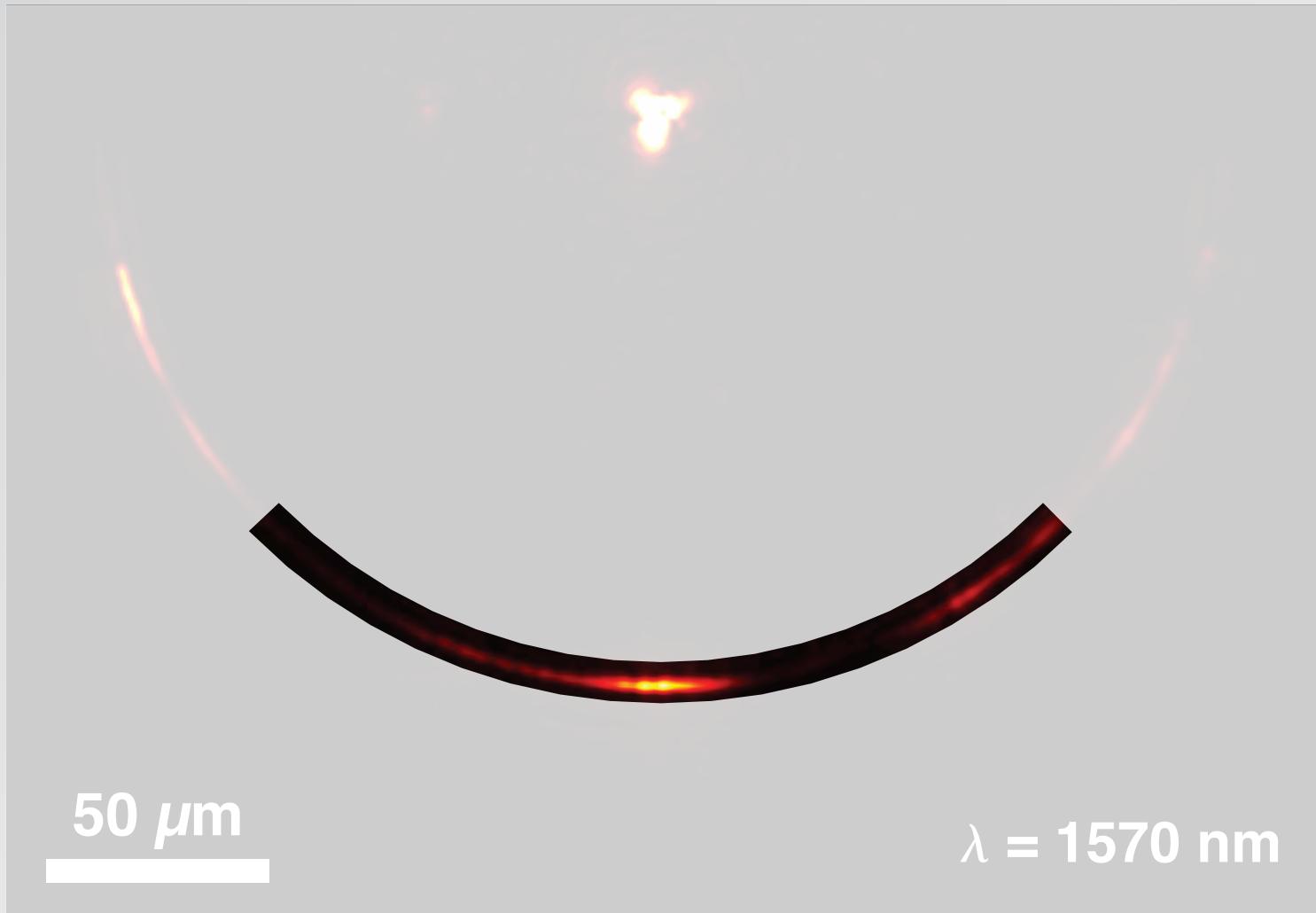
1 index

2 zero index



1 index

2 zero index



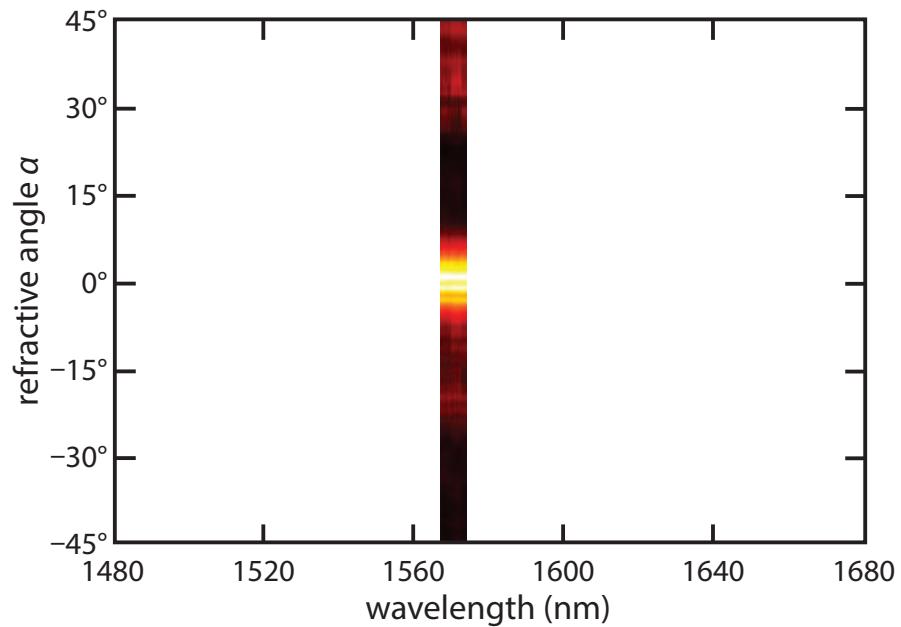
50 μm

$\lambda = 1570 \text{ nm}$

1 index

2 zero index

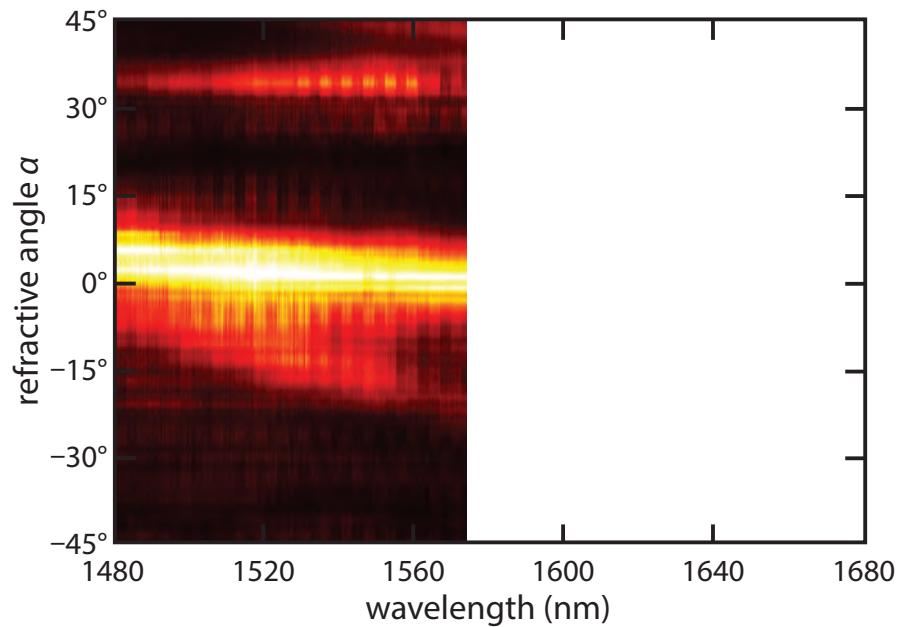
Wavelength dependence of refraction angle



1 index

2 zero index

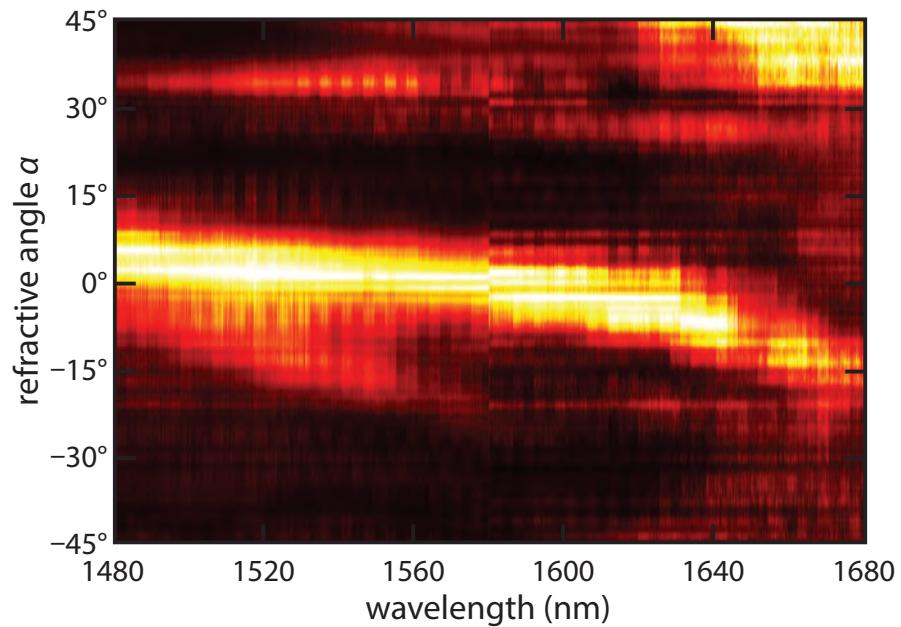
Wavelength dependence of refraction angle



1 index

2 zero index

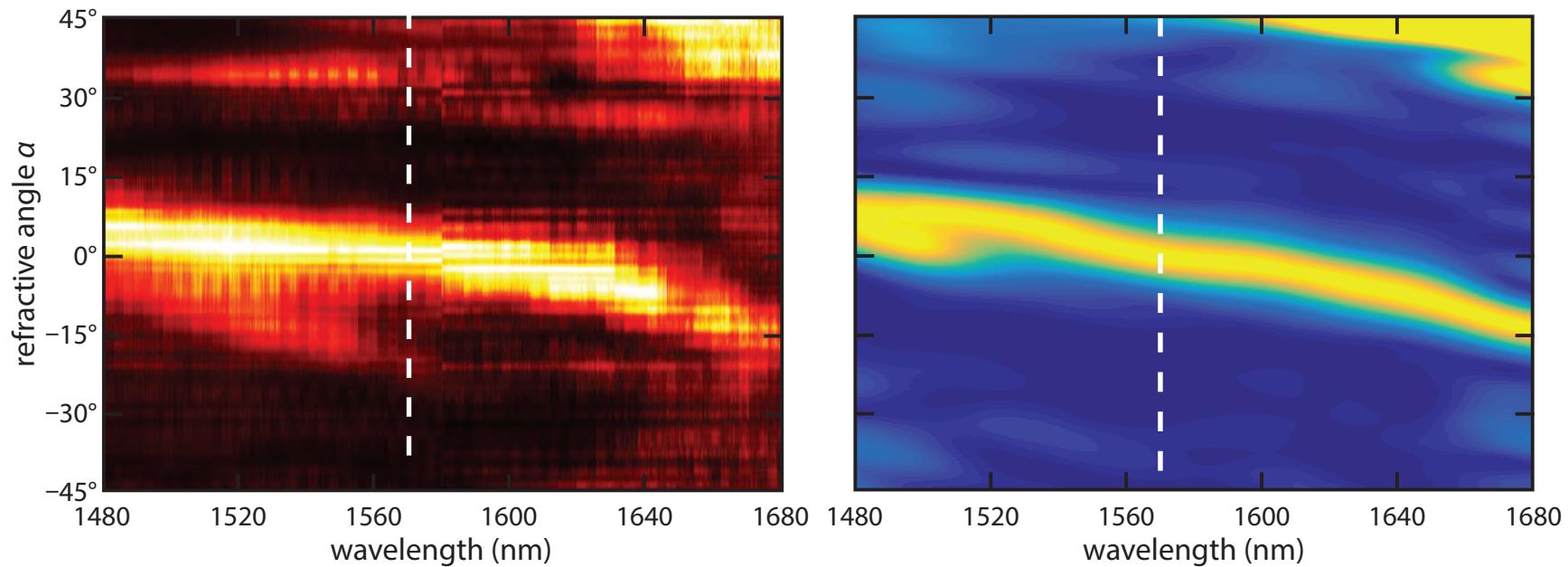
Wavelength dependence of refraction angle



1 index

2 zero index

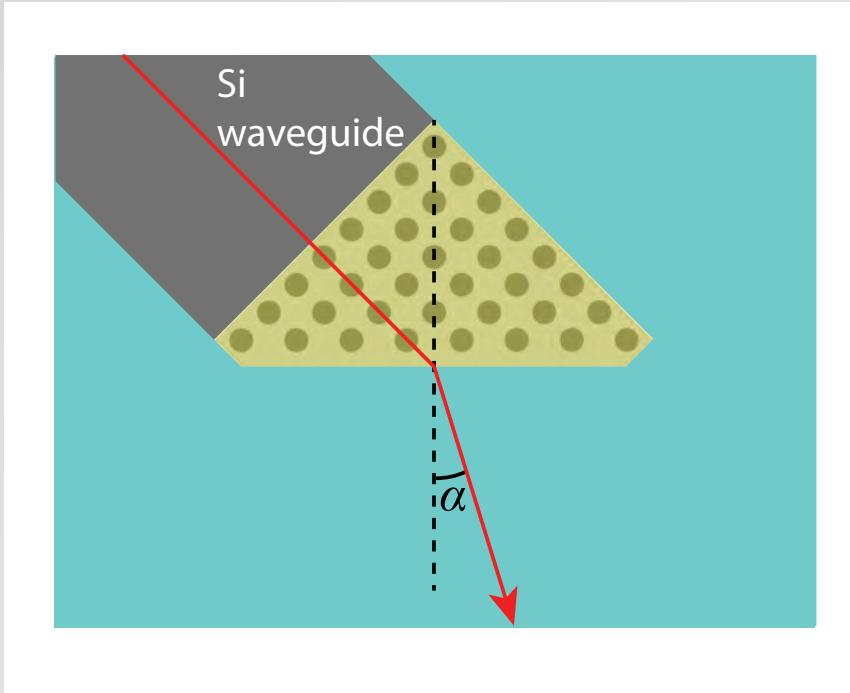
Wavelength dependence of refraction angle



1 index

2 zero index

Wavelength dependence of index

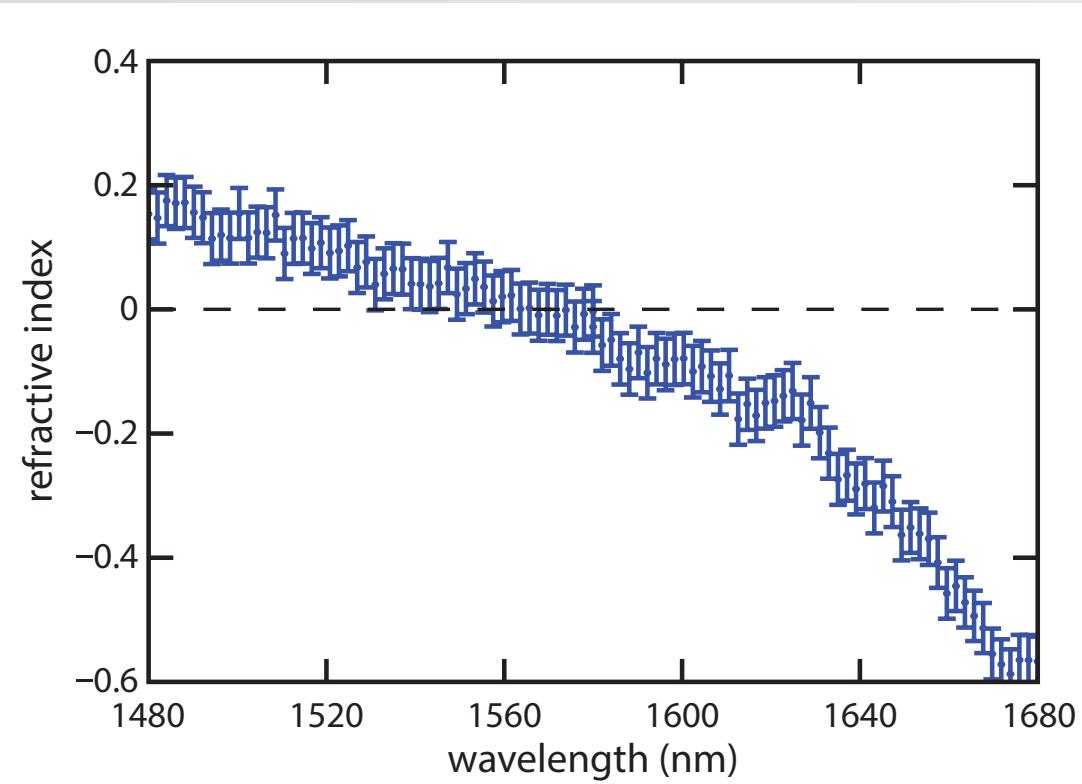


$$n_{\text{prism}} = n_{\text{slab}} \frac{\sin \alpha}{\sin 45^\circ}$$

1 index

2 zero index

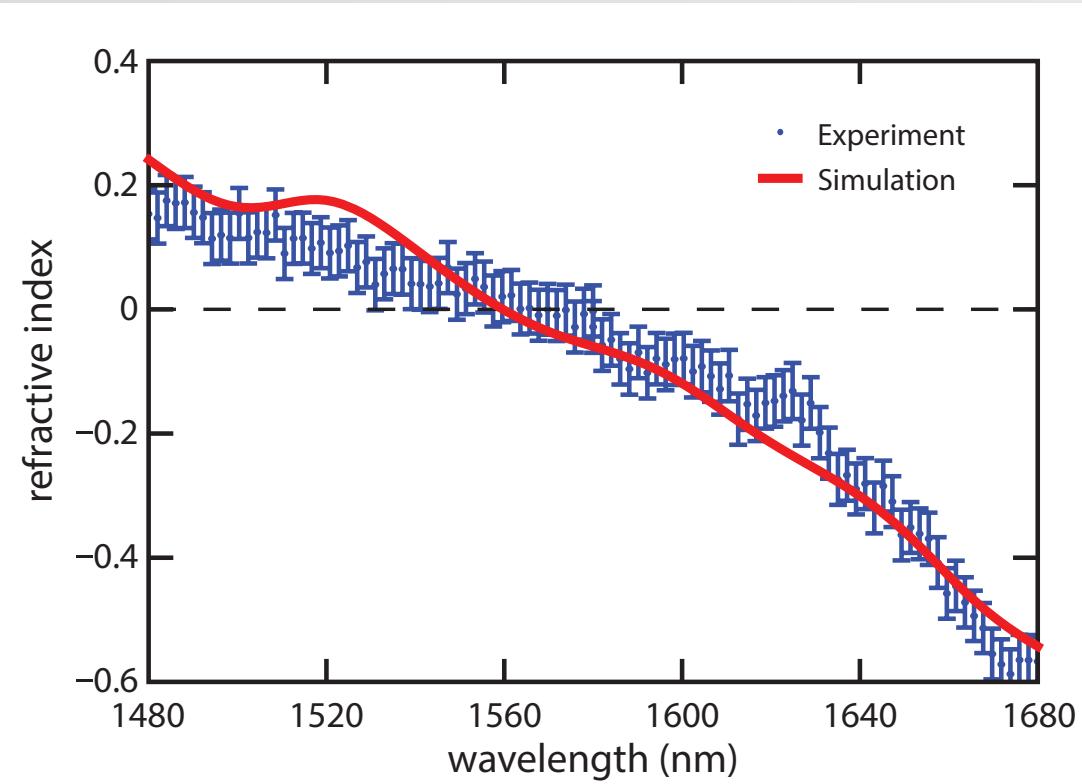
Wavelength dependence of index



1 index

2 zero index

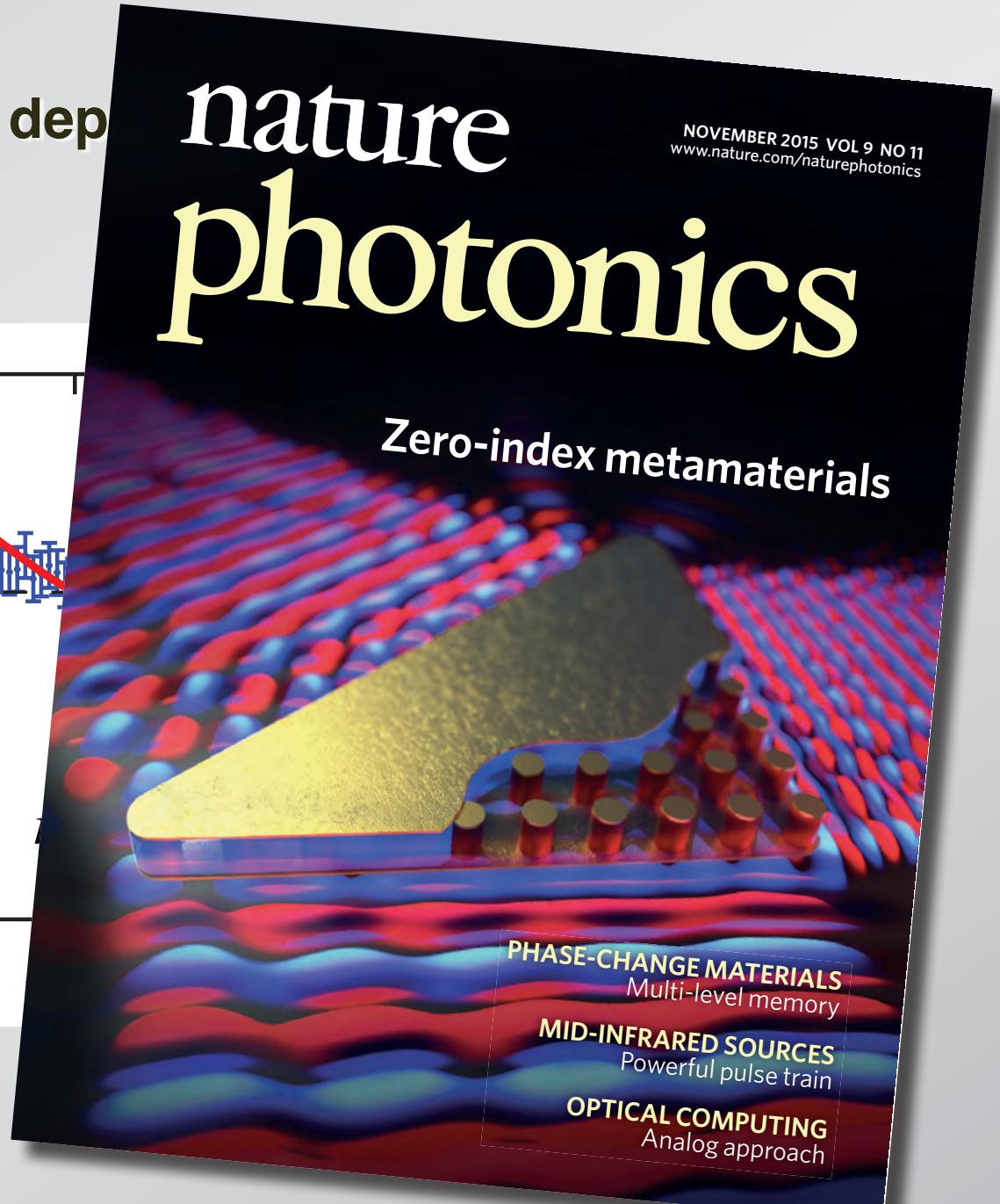
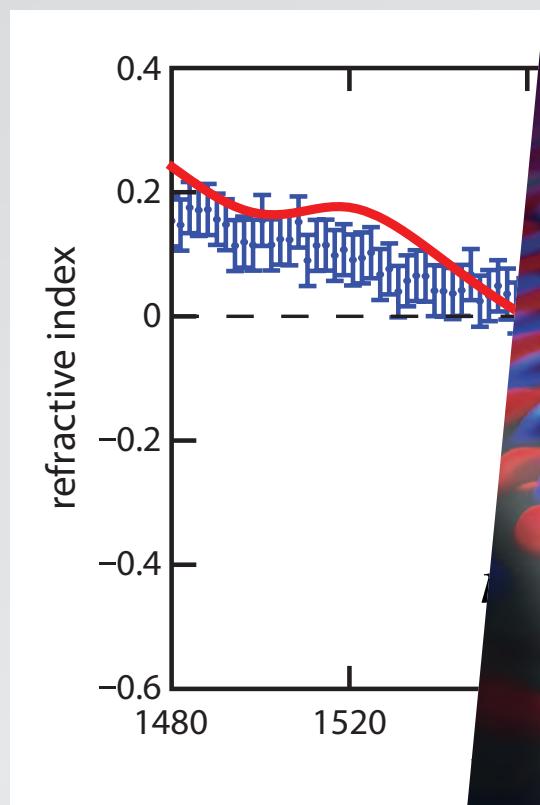
Wavelength dependence of index



1 index

2 zero index

Wavelength dep

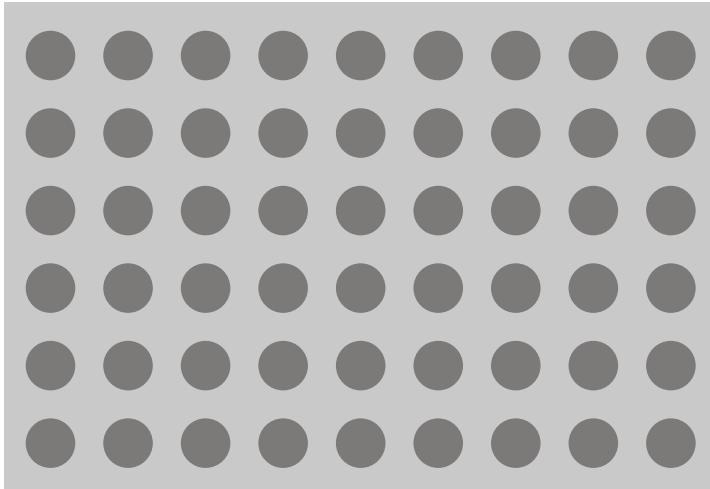


1 index

2 zero index

simplify fabrication

pillar array

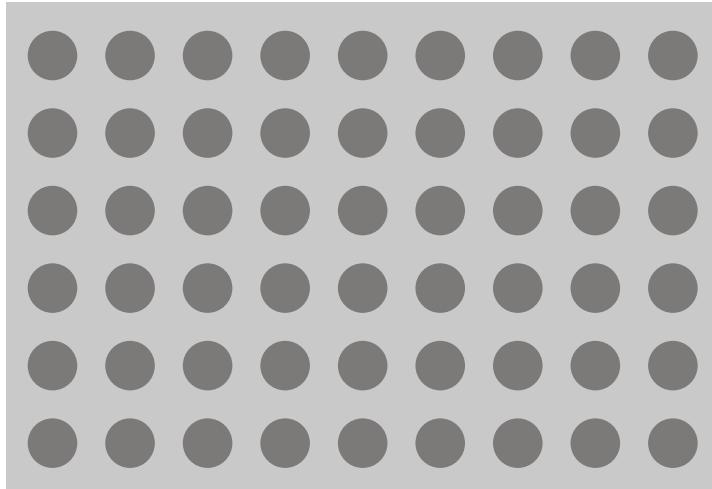


1 index

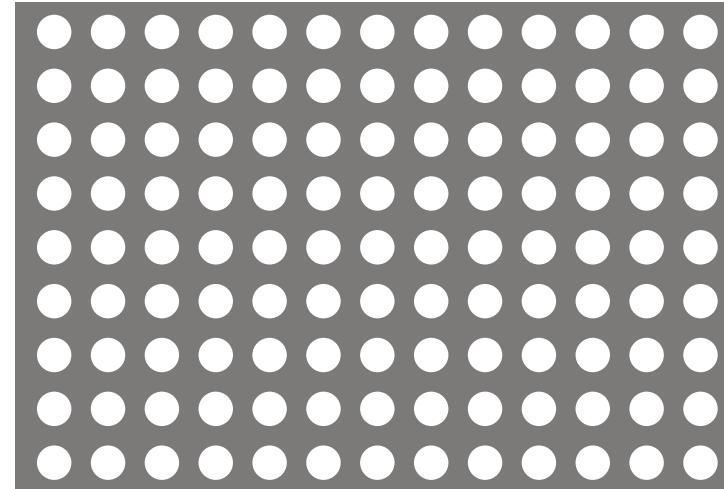
2 zero index

simplify fabrication

pillar array



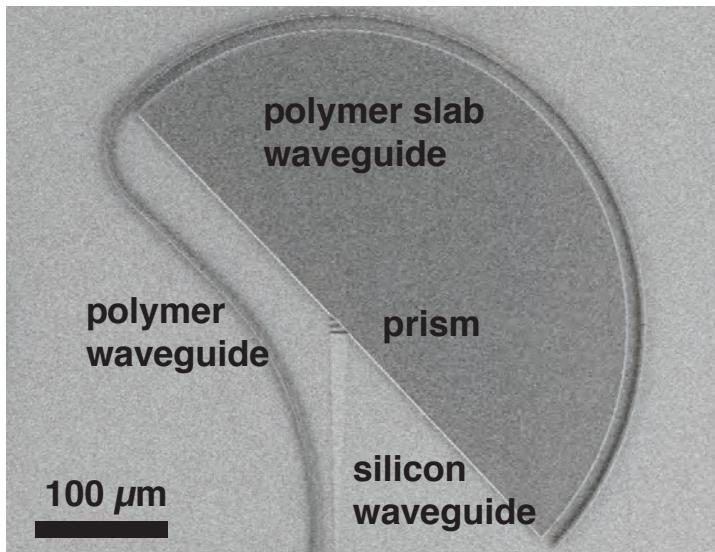
airhole array



1 index

2 zero index

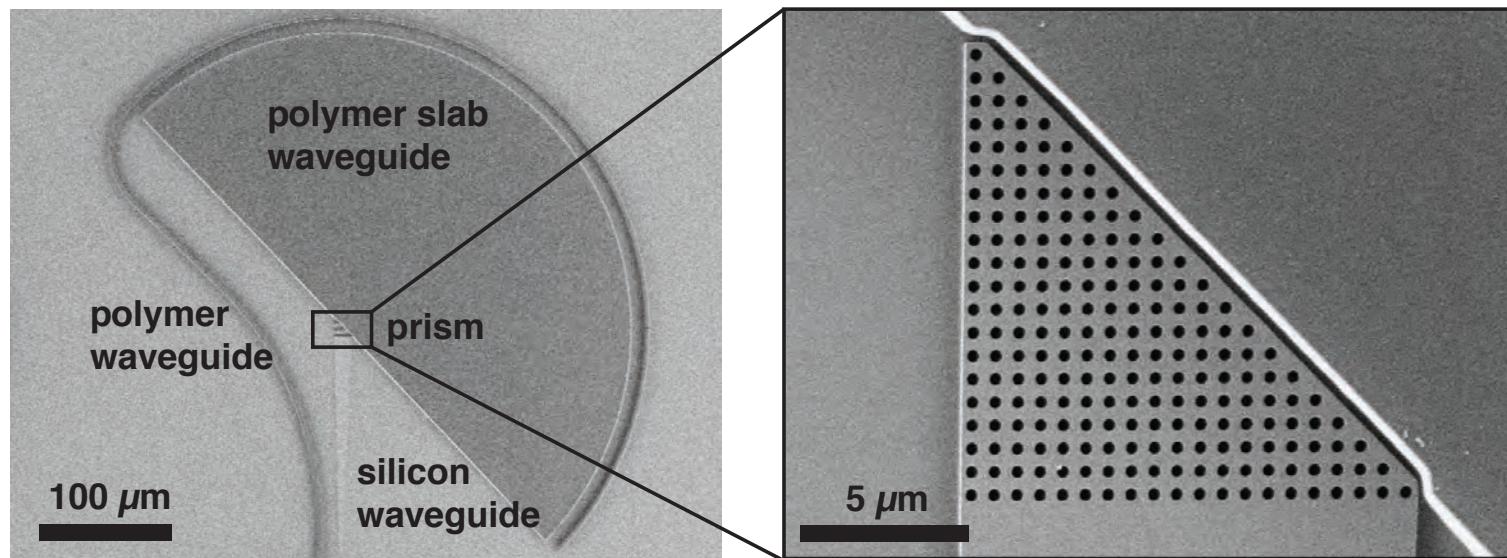
simplify fabrication



1 index

2 zero index

simplify fabrication

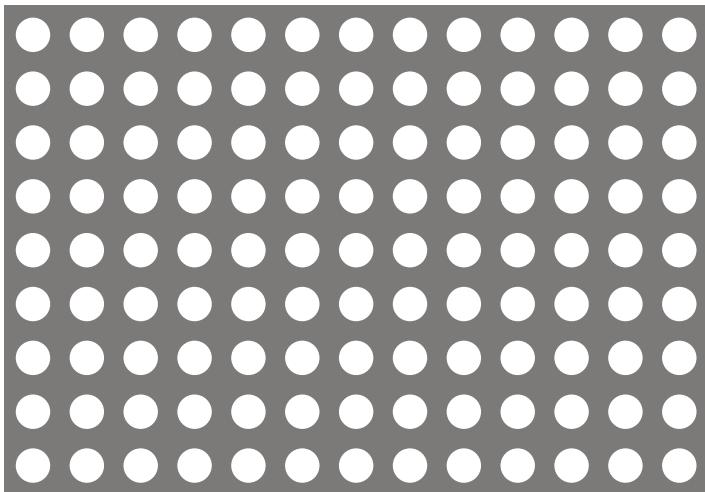


1 index

2 zero index

simplify further!

airhole array

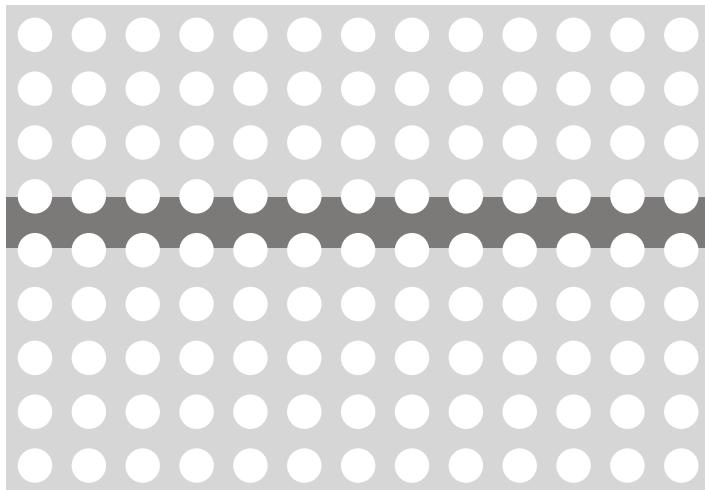


1 index

2 zero index

simplify further!

airhole array

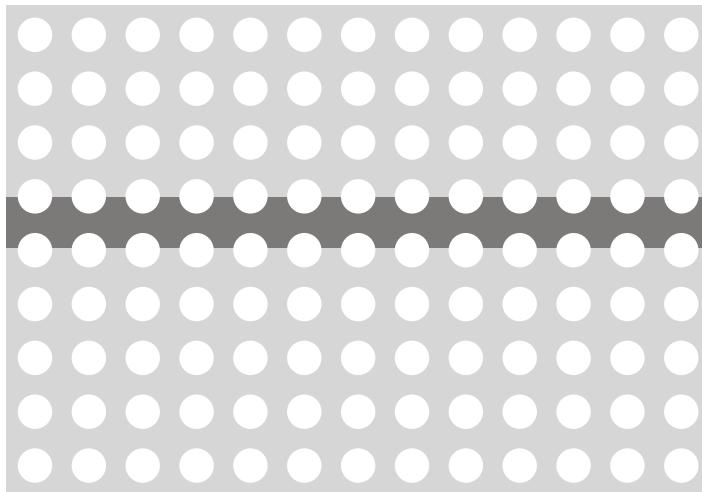


1 index

2 zero index

simplify further!

airhole array



1D ZIM waveguide

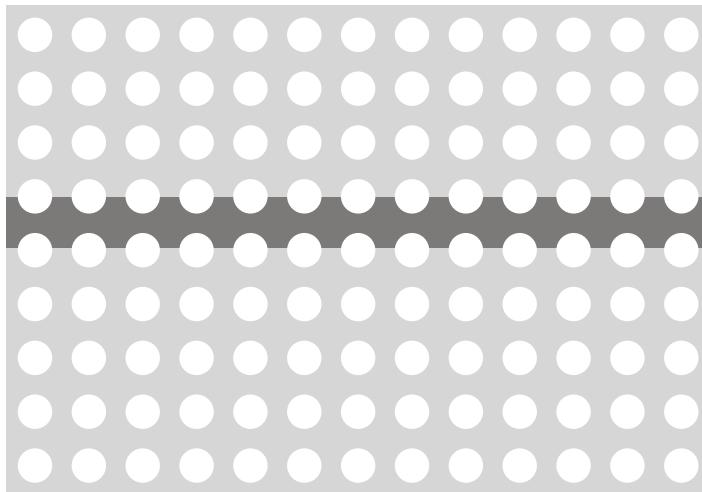


1 index

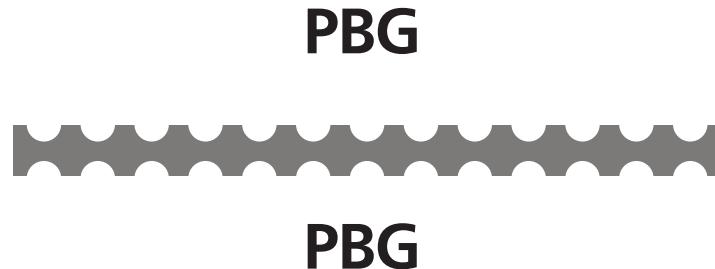
2 zero index

simplify further!

airhole array



1D ZIM waveguide



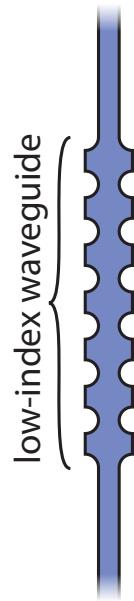
PBG

PBG

1 index

2 zero index

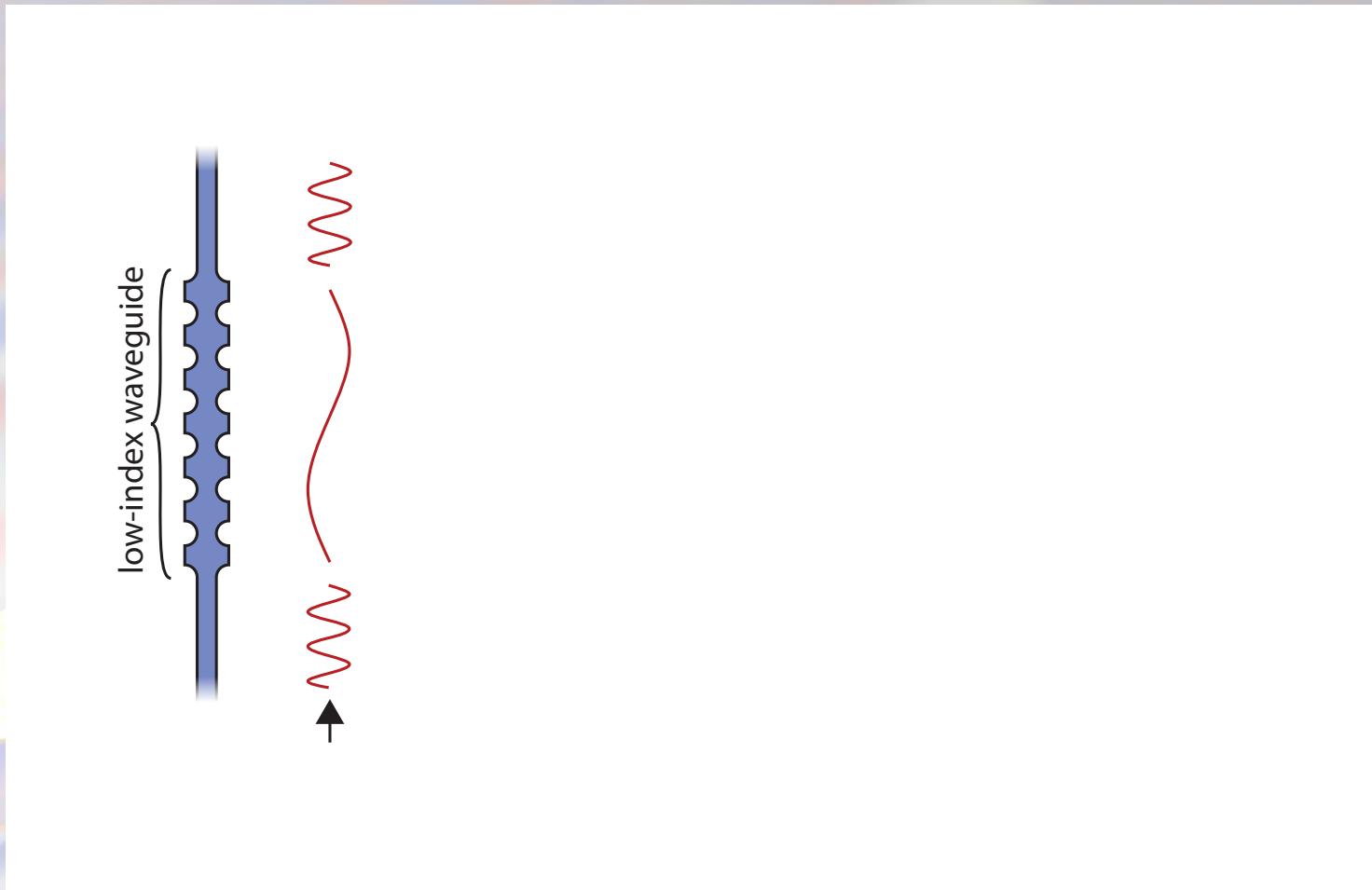
waveguiding



1 index

2 zero index

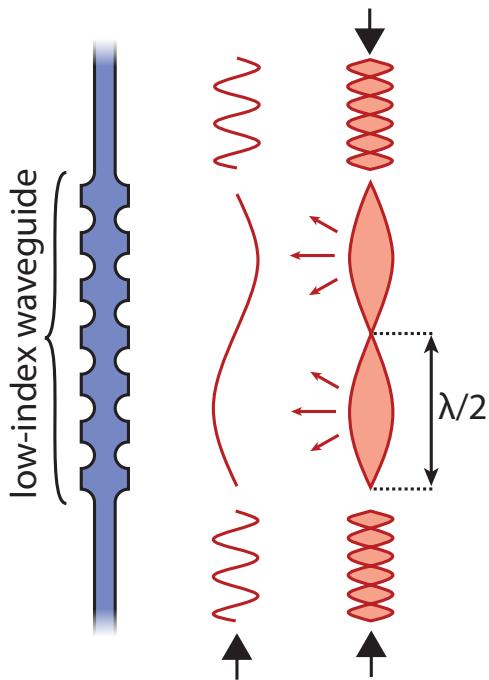
waveguiding



1 index

2 zero index

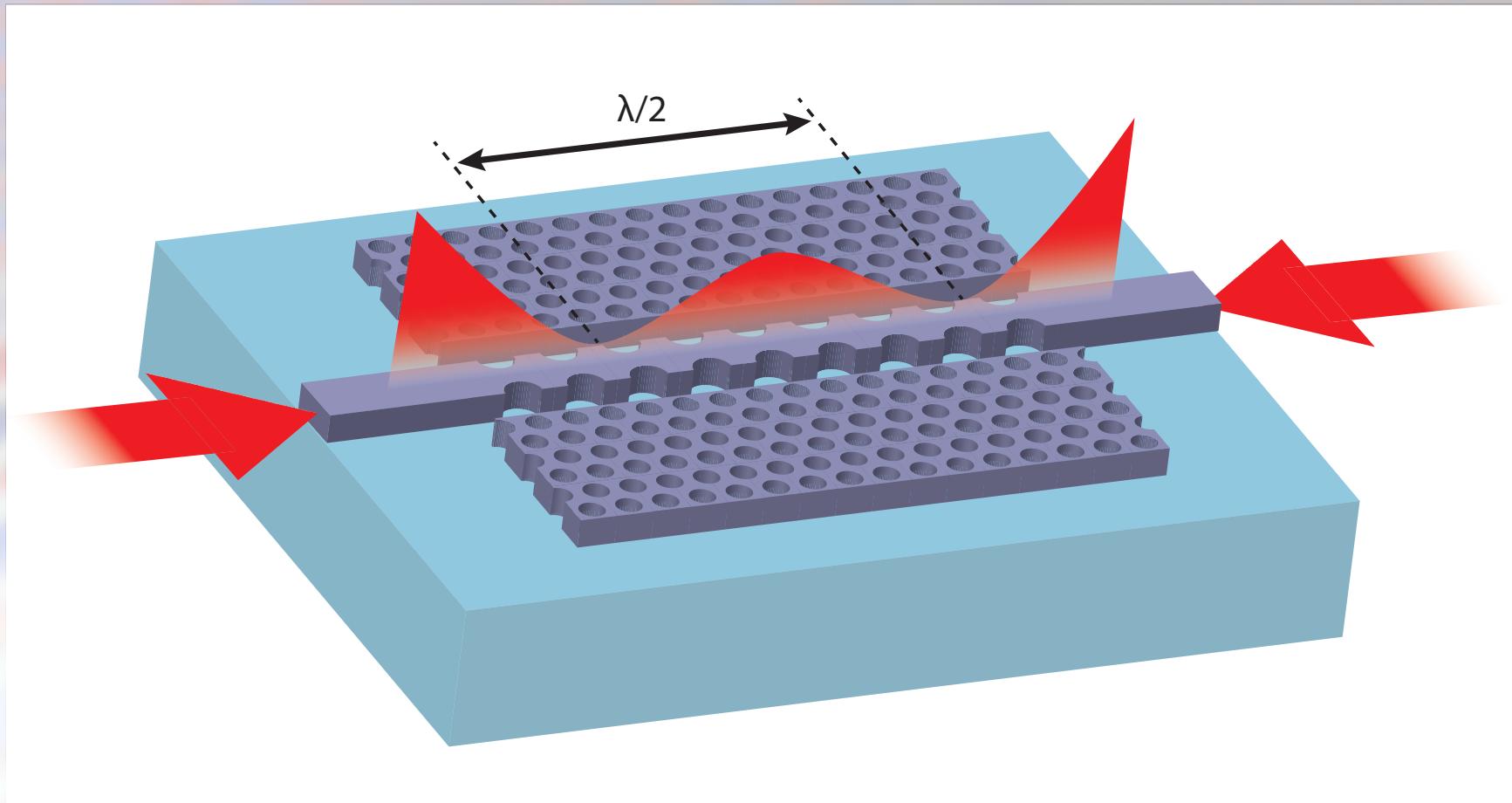
waveguiding



1 index

2 zero index

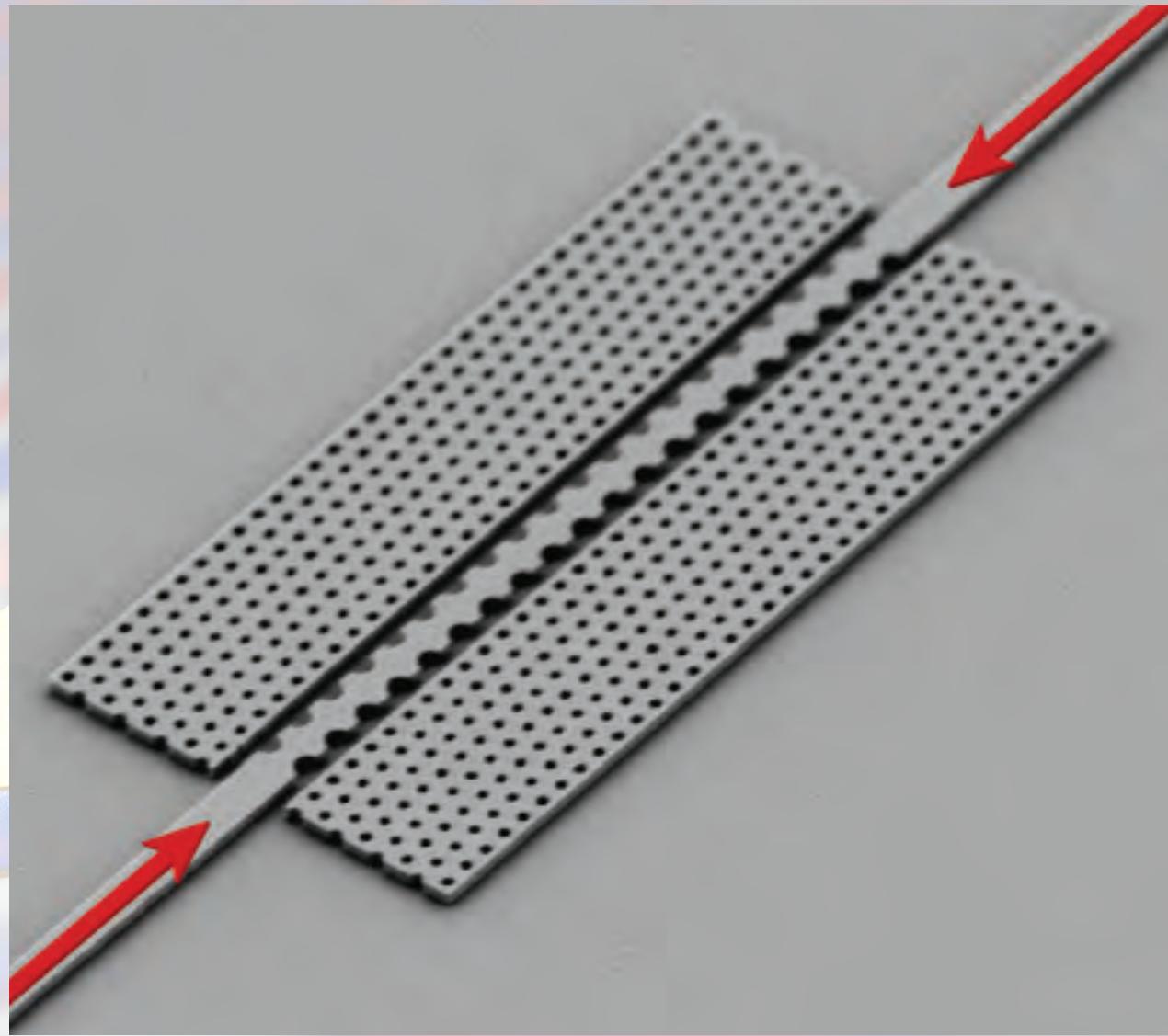
look at standing waves



1 index

2 zero index

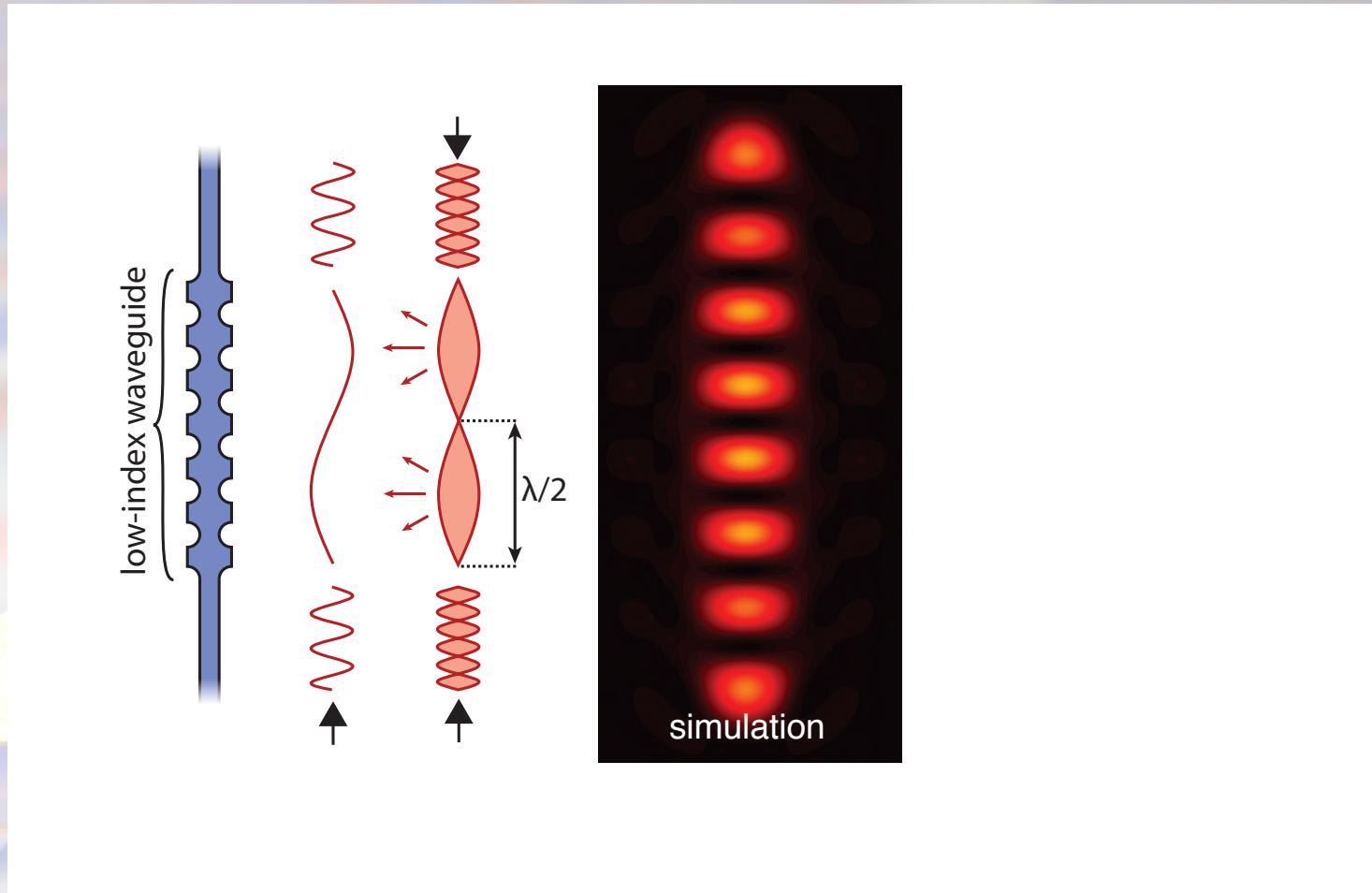
look at standing waves



1 index

2 zero index

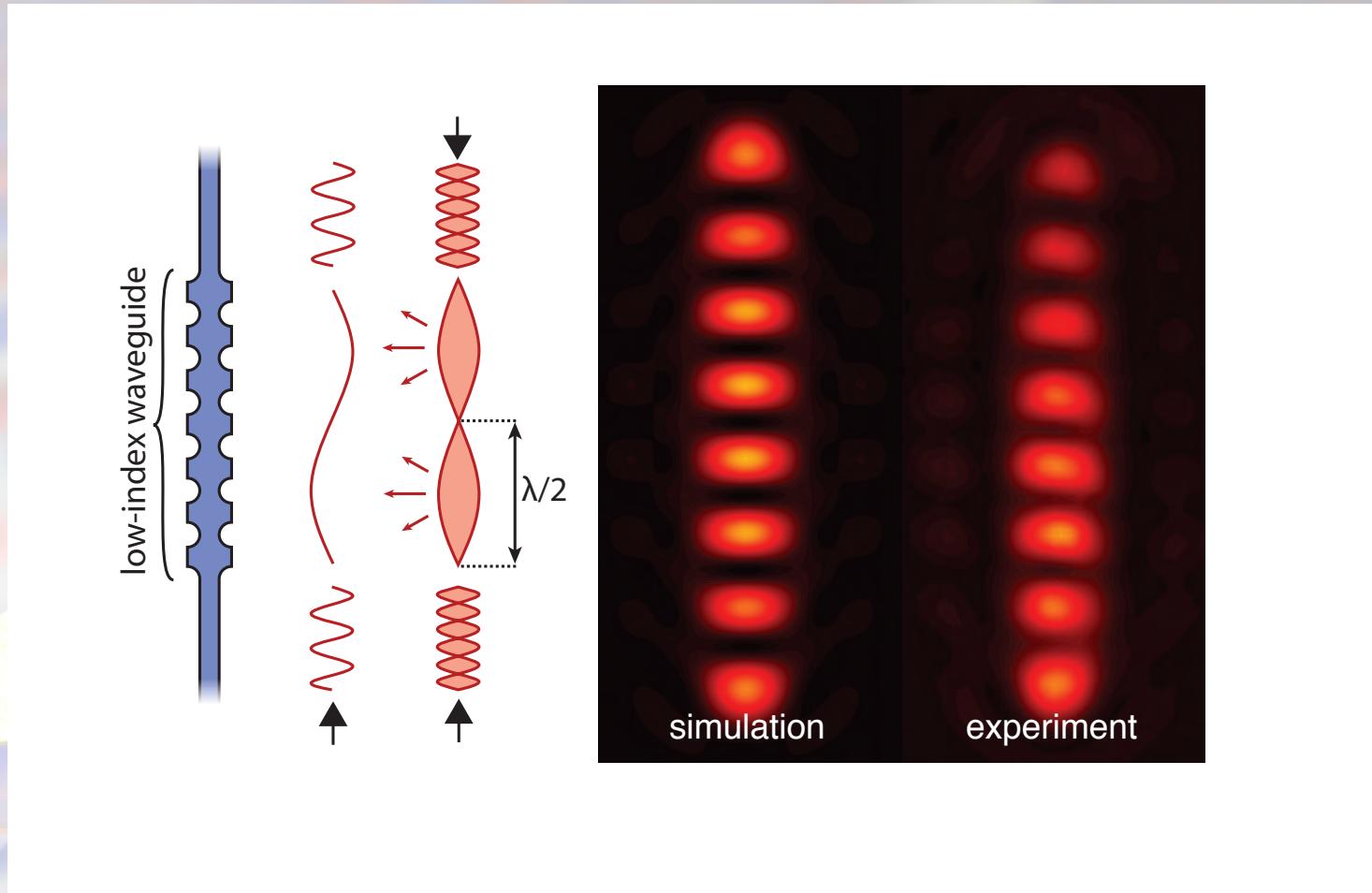
look at standing waves



1 index

2 zero index

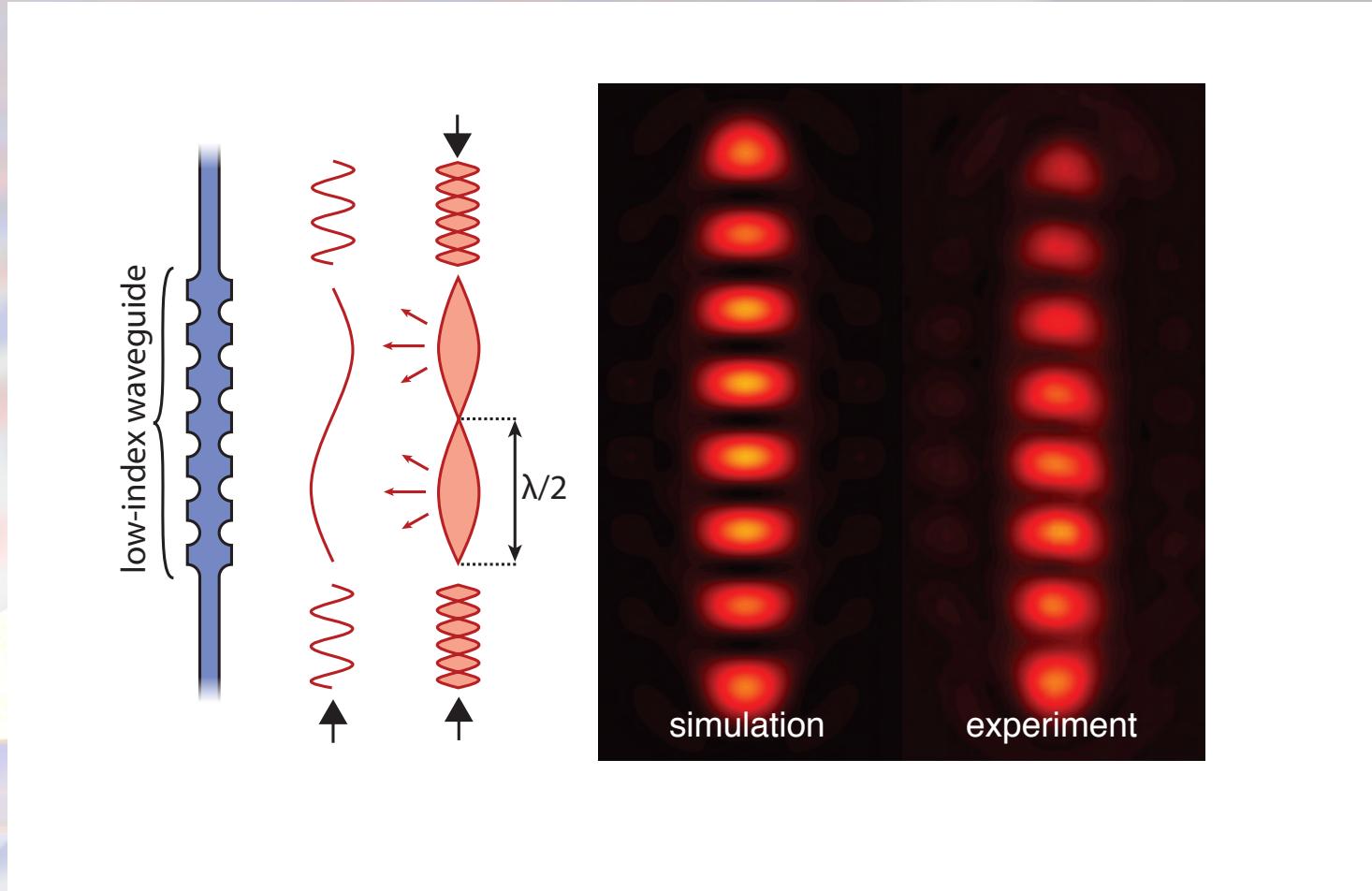
look at standing waves



1 index

2 zero index

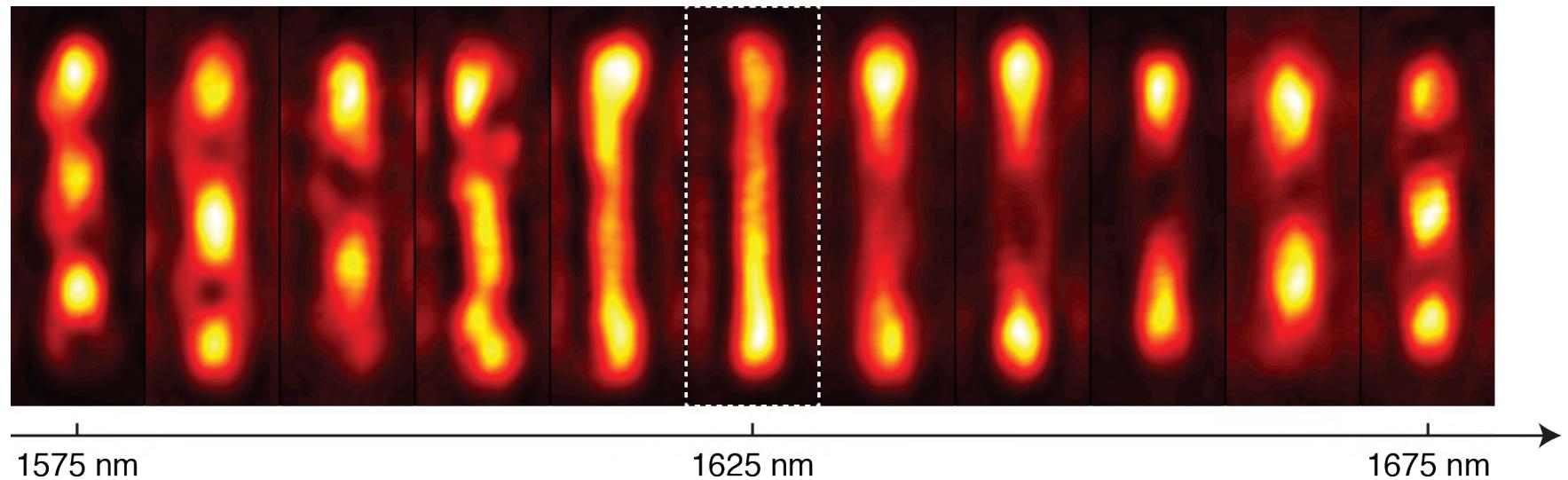
direct observation of effective wavelength!!



1 index

2 zero index

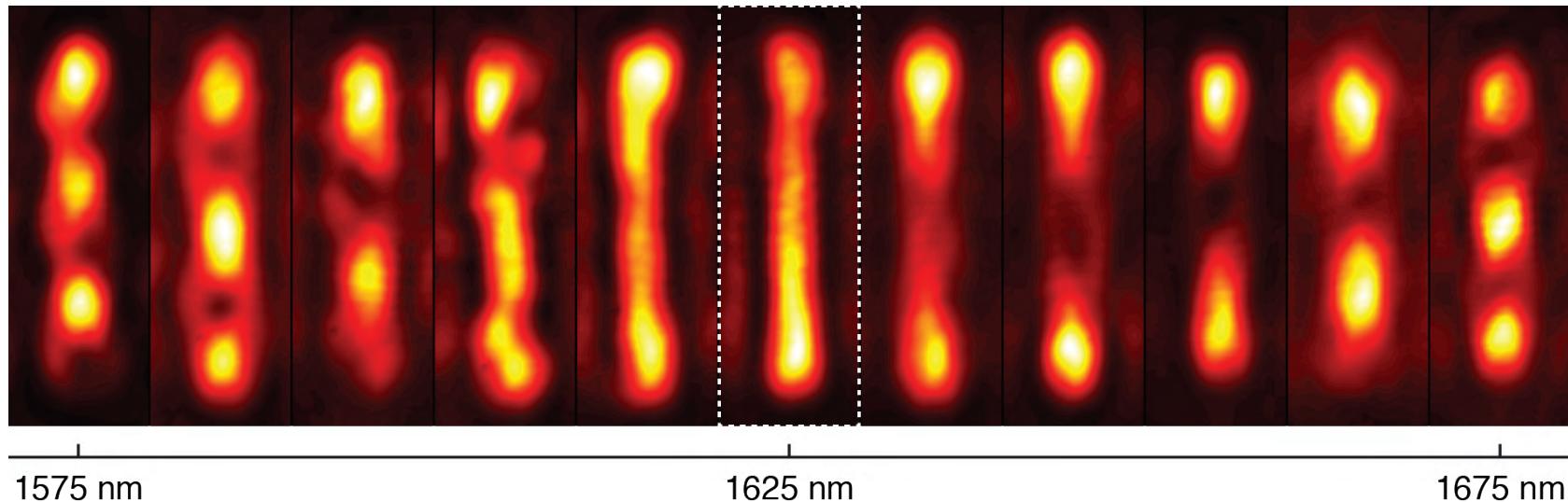
look at standing waves



1 index

2 zero index

look at standing waves

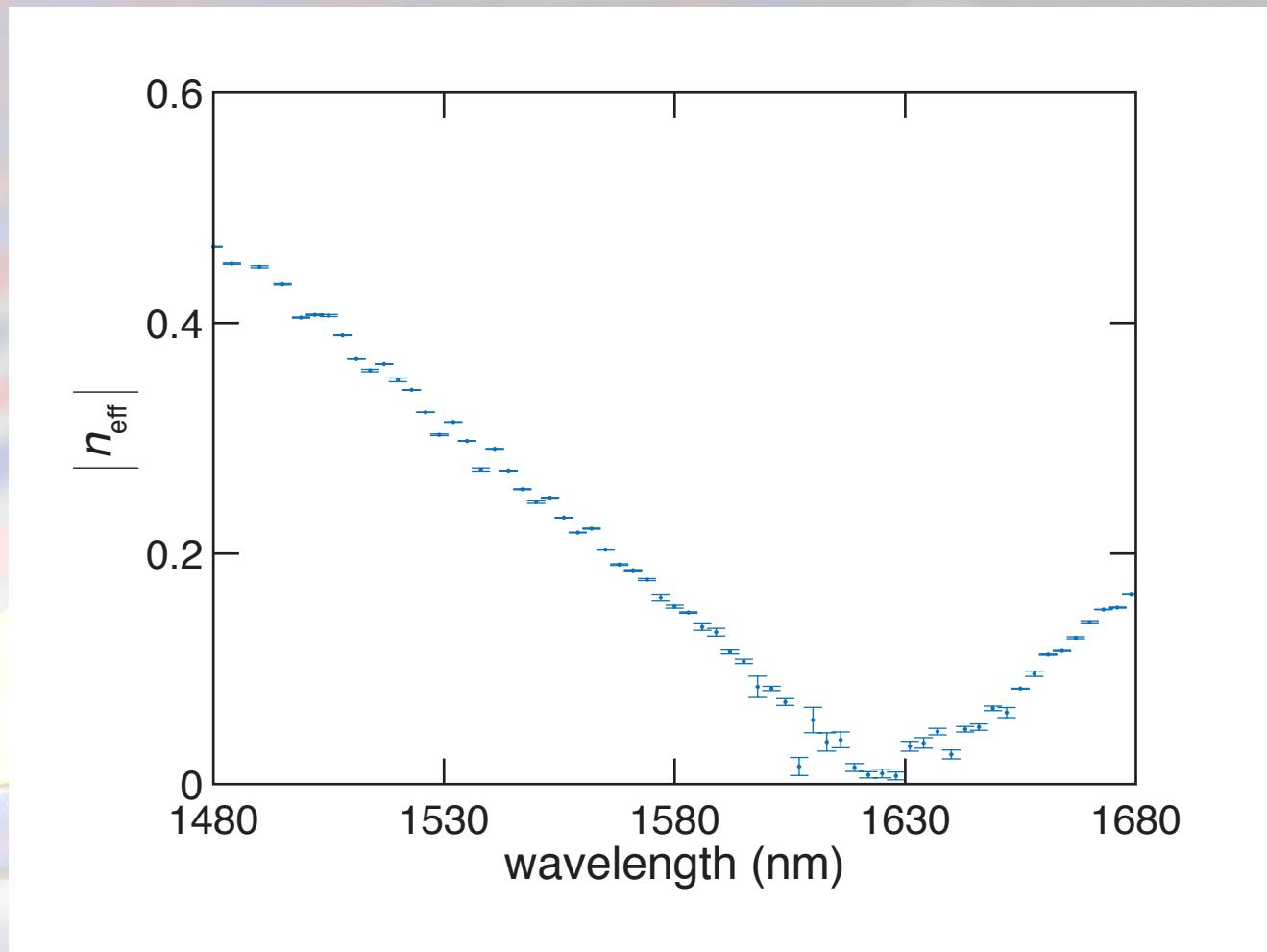


$$| n_{\text{eff}} | = \frac{\lambda_0}{\lambda_{\text{eff}}}$$

1 index

2 zero index

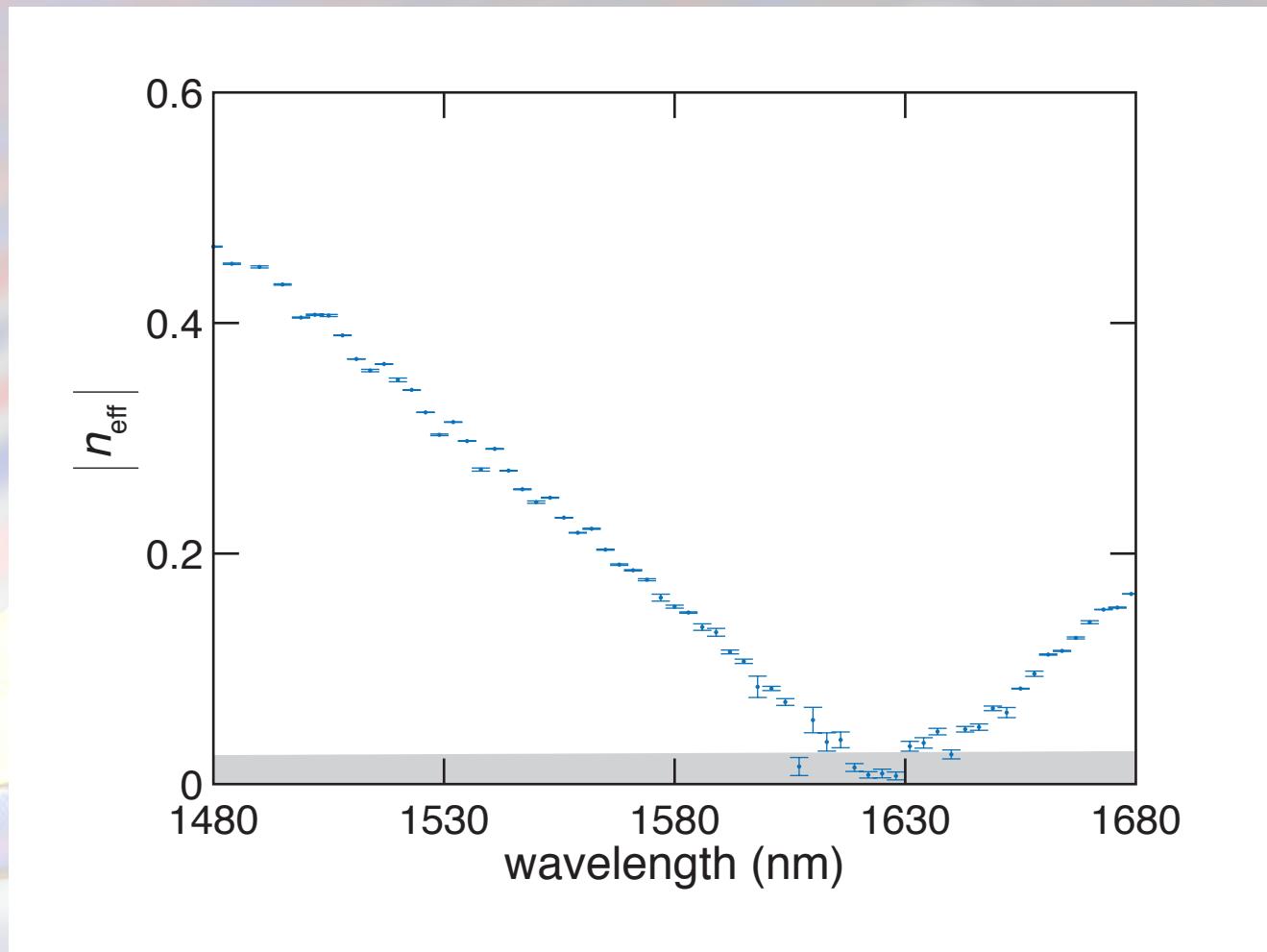
comparison of experiment and simulation



1 index

2 zero index

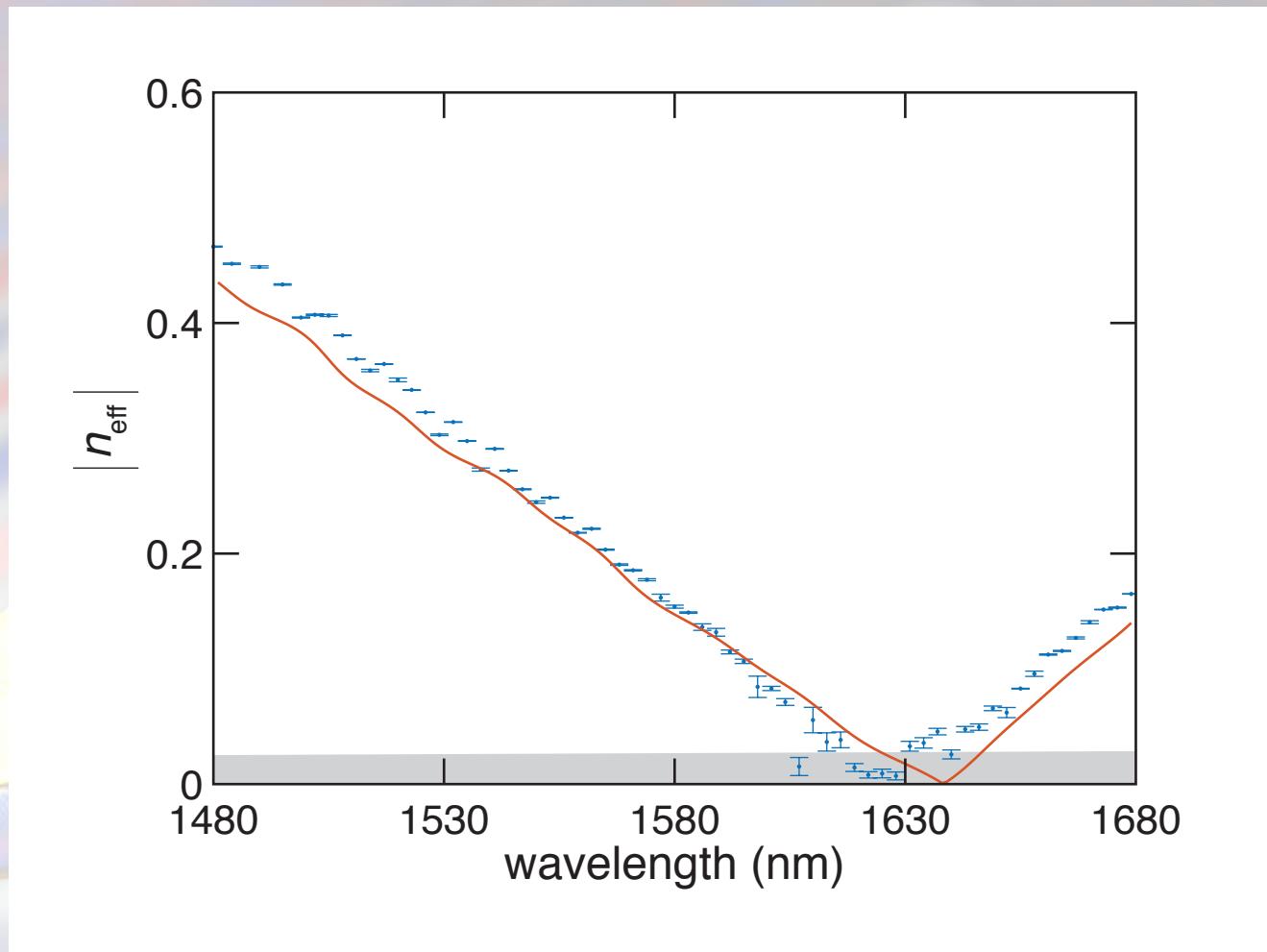
comparison of experiment and simulation



1 index

2 zero index

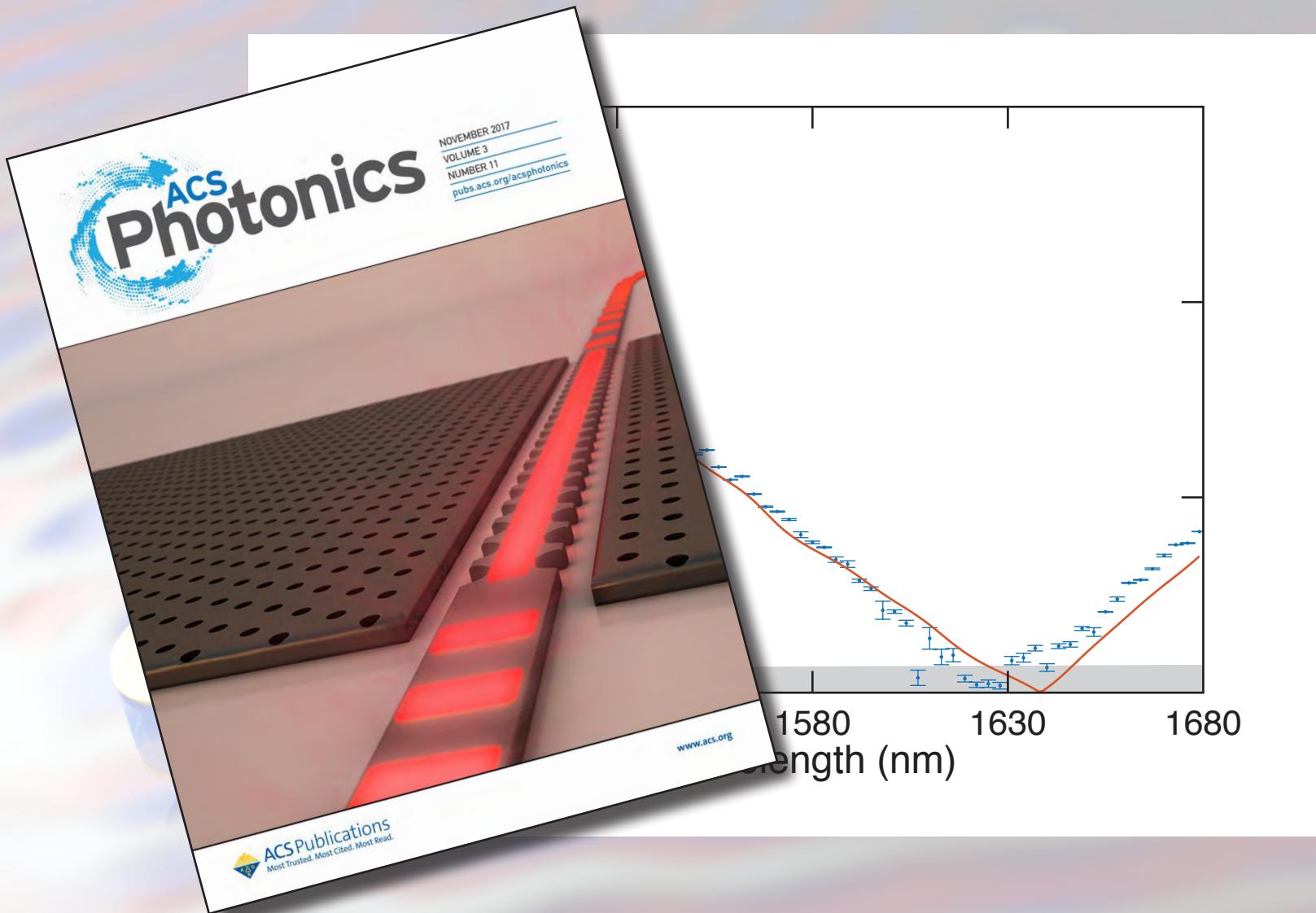
comparison of experiment and simulation



1 index

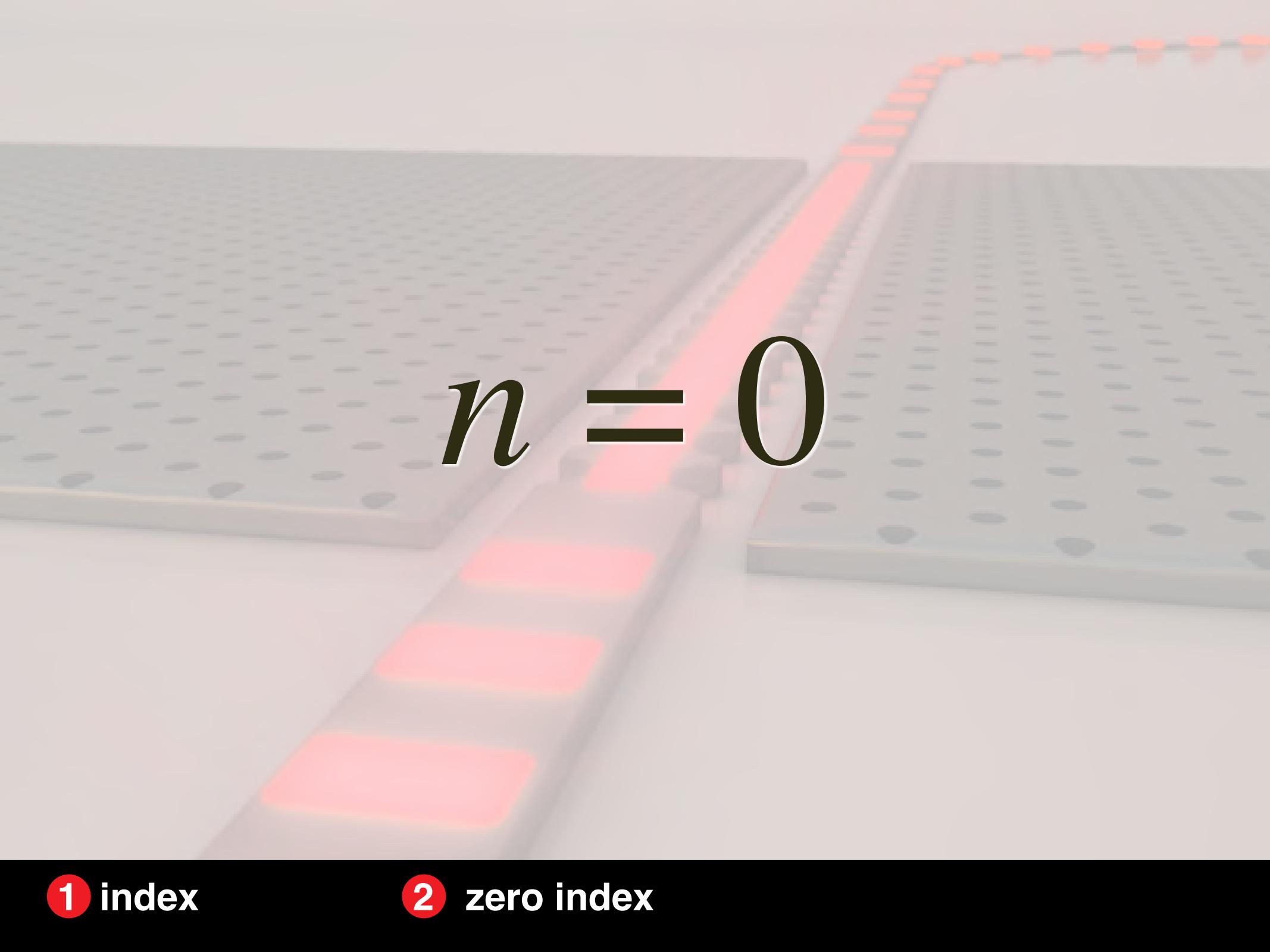
2 zero index

comparison of experiment and simulation



1 index

2 zero index



$n = 0$

1 index

2 zero index

(extreme) opportunities

relaxed phase matching constraints as $k \rightarrow 0$

1 index

2 zero index

(extreme) opportunities

PHYSICAL REVIEW LETTERS 128, 203902 (2022)

Editors' Suggestion

Relaxed Phase-Matching Constraints in Zero-Index Waveguides

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(Received 1 July 2021; accepted 4 April 2022; published 17 May 2022)

The utility of all parametric nonlinear optical processes is hampered by phase-matching requirements. Quasi-phase-matching, birefringent phase matching, and higher-order-mode phase matching have all been developed to address this constraint, but the methods demonstrated to date suffer from the inconvenience of only being phase matched for a single, specific arrangement of beams, typically copropagating, resulting in cumbersome experimental configurations and large footprints for integrated devices. Here, we experimentally demonstrate that these phase-matching requirements may be satisfied in a parametric nonlinear waveguide using two counterpropagating input and output beams when using low-index

$$v_g = \frac{d\omega}{dk}$$

1 index

2 zero index

3 zero group velocity

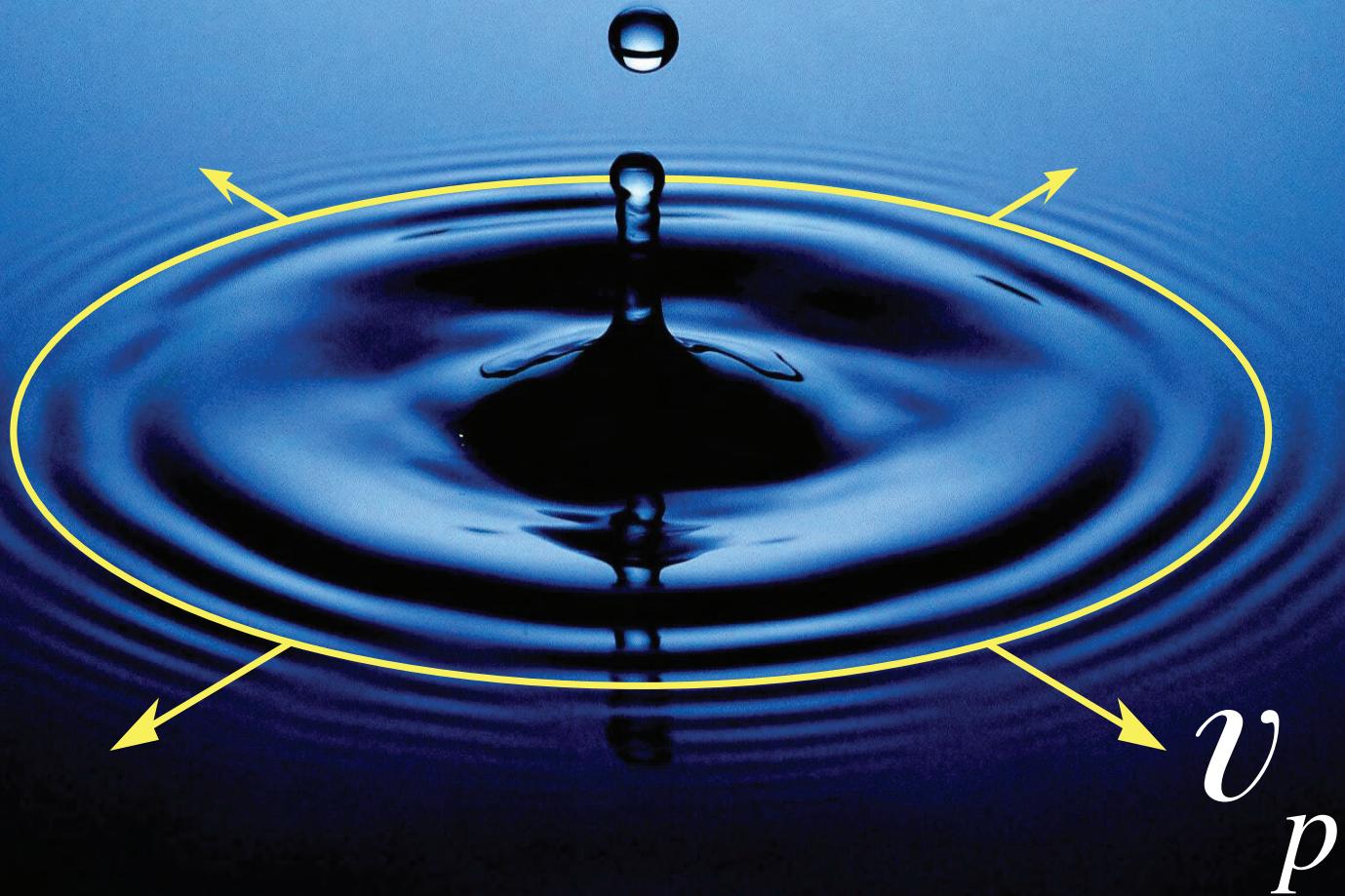


1 index

2 zero index

3 zero group velocity

phase velocity



1 index

2 zero index

3 zero group velocity

group velocity



1 index

2 zero index

3 zero group velocity

$$v_g = \frac{1}{2} v_p$$



1 index

2 zero index

3 zero group velocity

$$v_g = 0$$



1 index

2 zero index

3 zero group velocity

$$v_g = 0$$



localization

1 index

2 zero index

3 zero group velocity

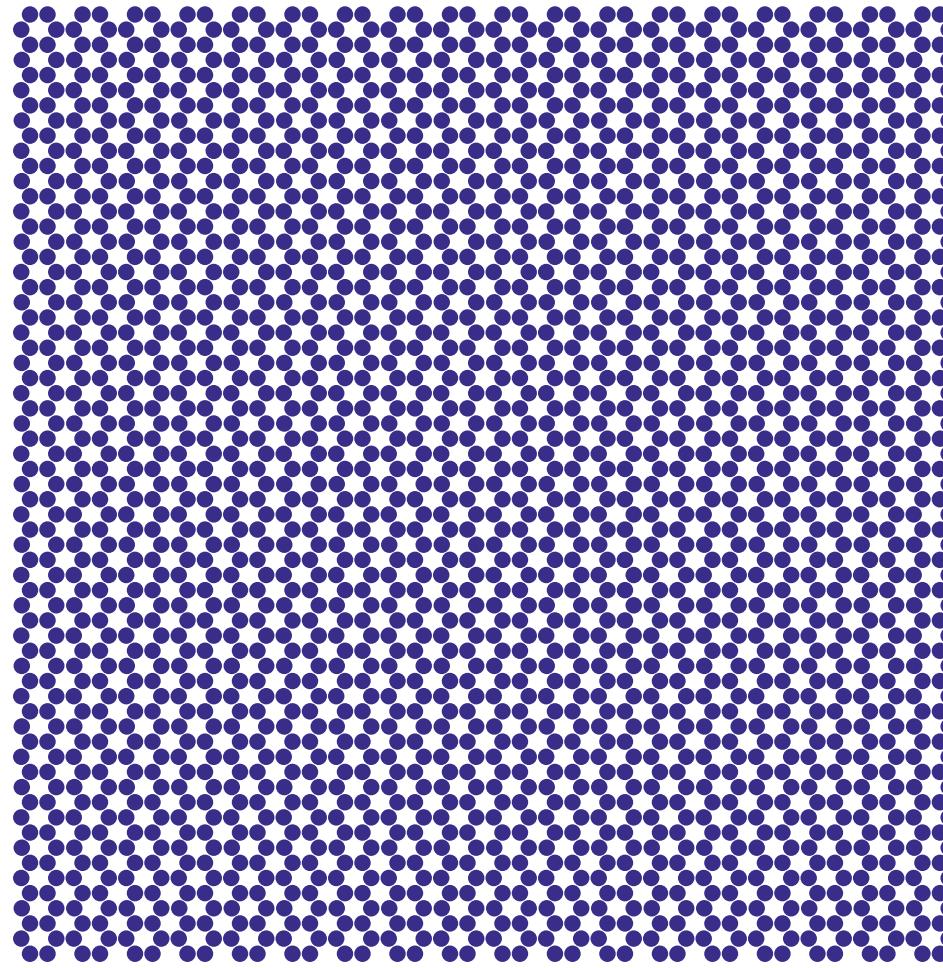


how can we localize light?

1 index

2 zero index

3 zero group velocity

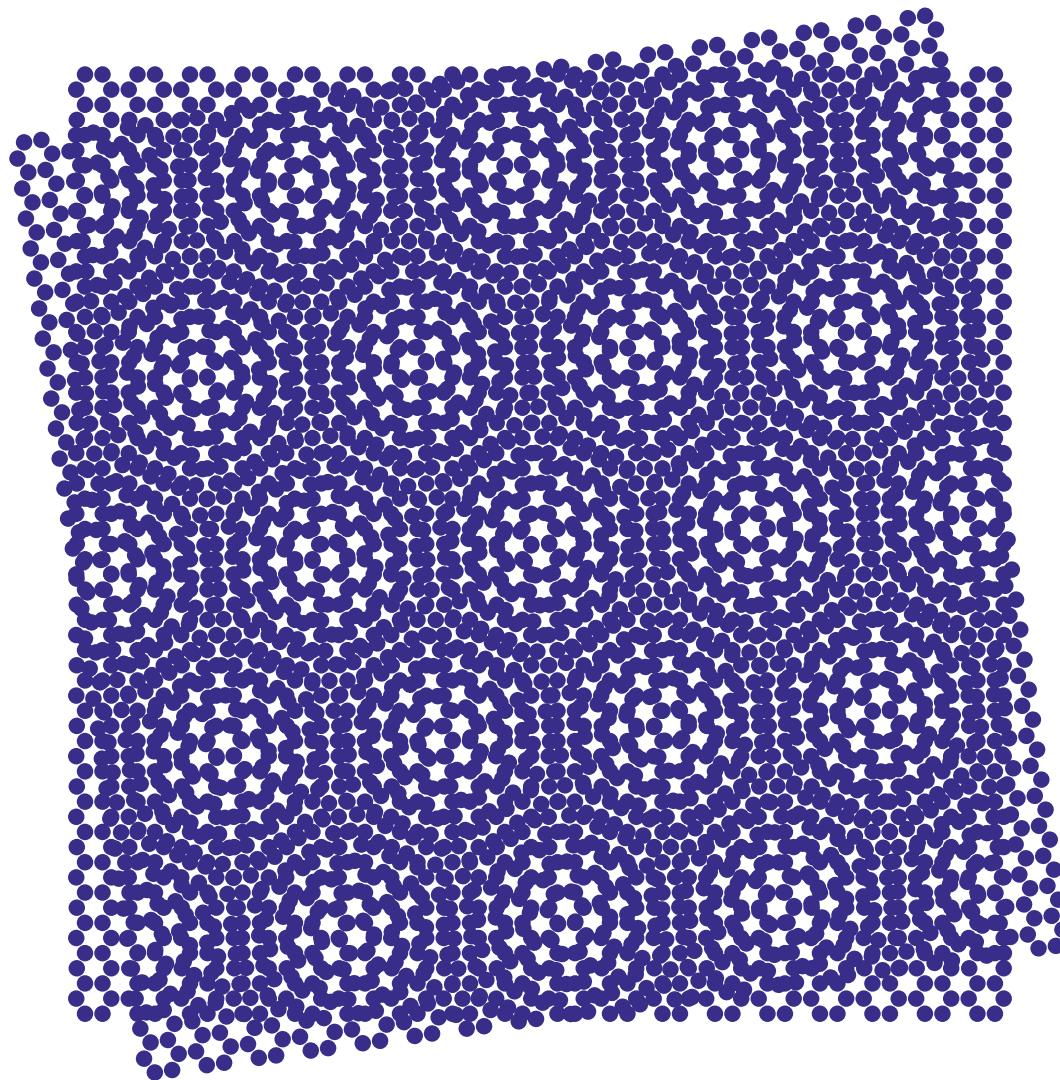


1 index

2 zero index

3 zero group velocity

8°

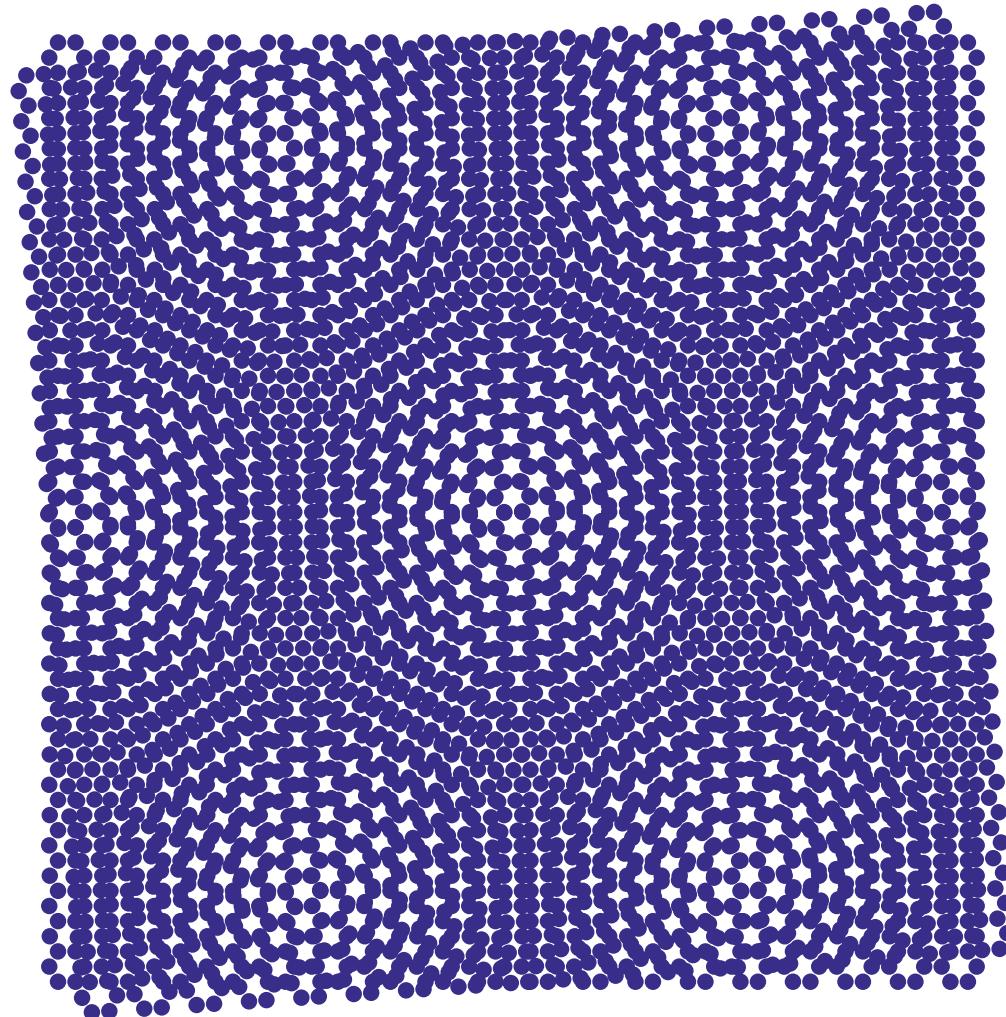


1 index

2 zero index

3 zero group velocity

4°

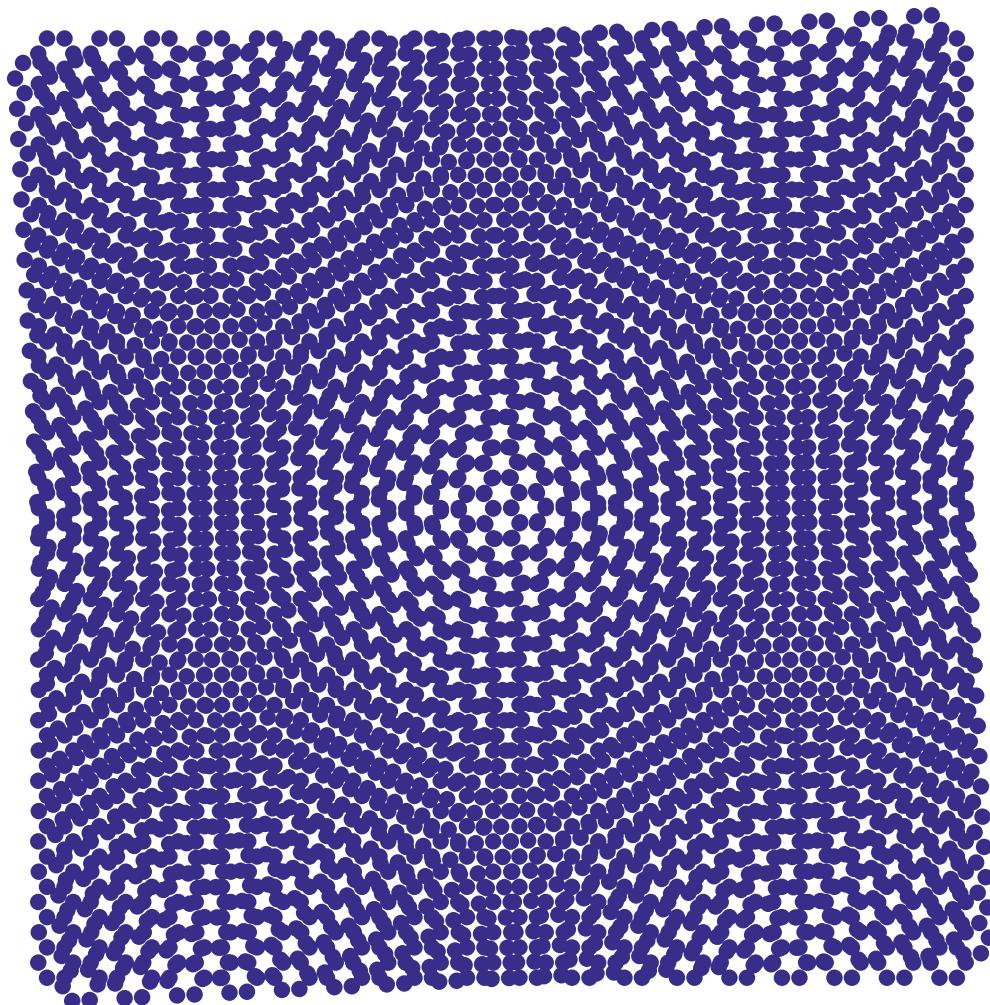


1 index

2 zero index

3 zero group velocity

3°

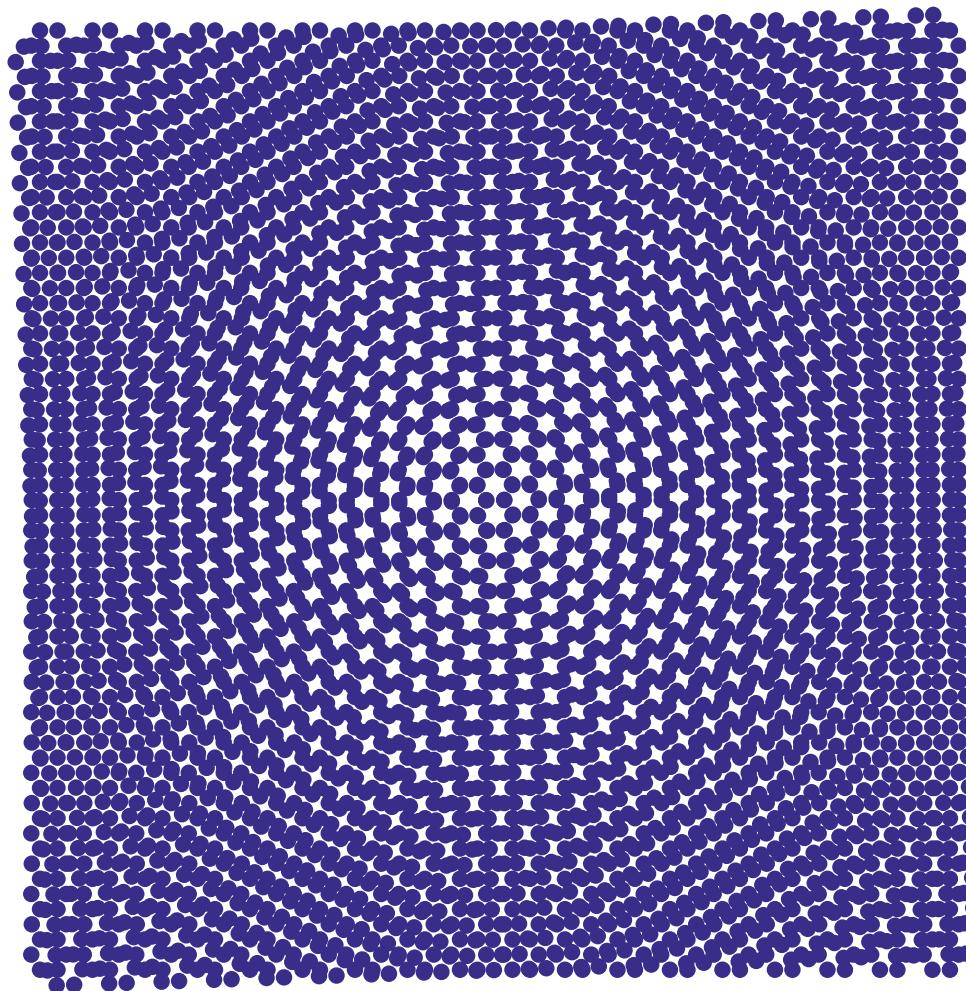


1 index

2 zero index

3 zero group velocity

2°

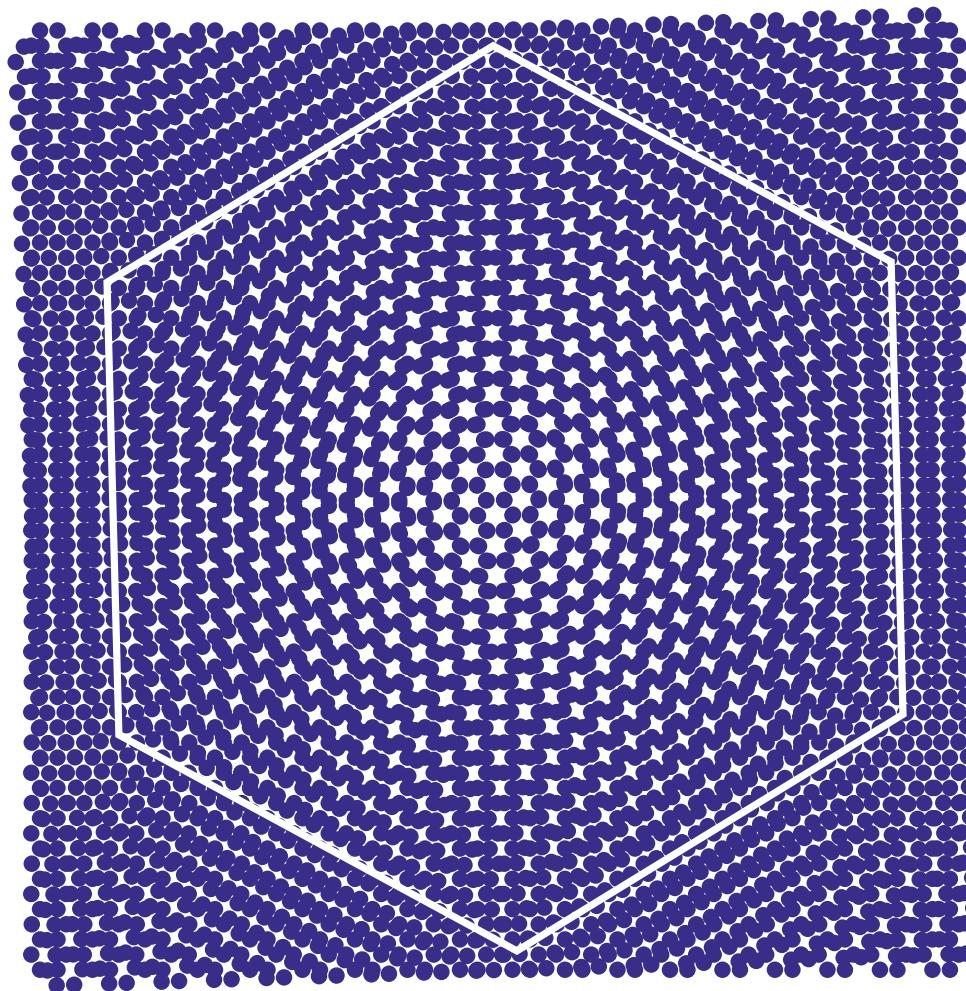


1 index

2 zero index

3 zero group velocity

2°

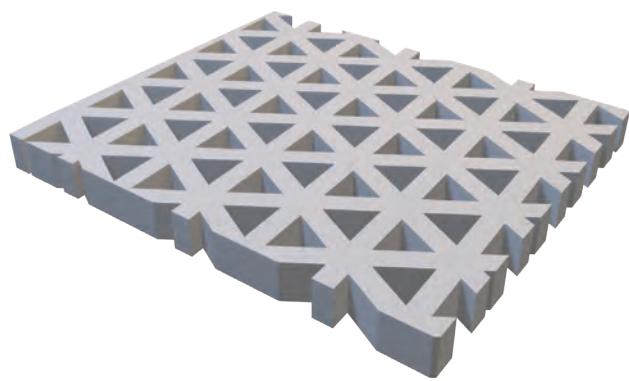


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

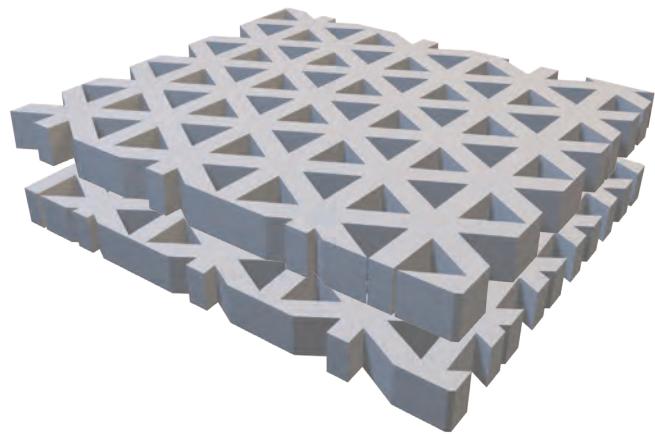


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

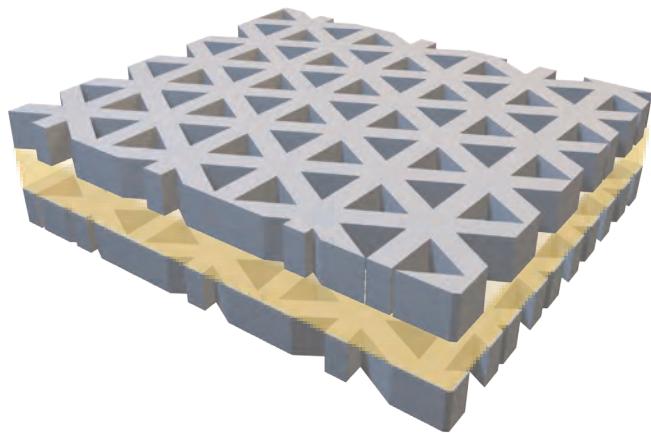


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

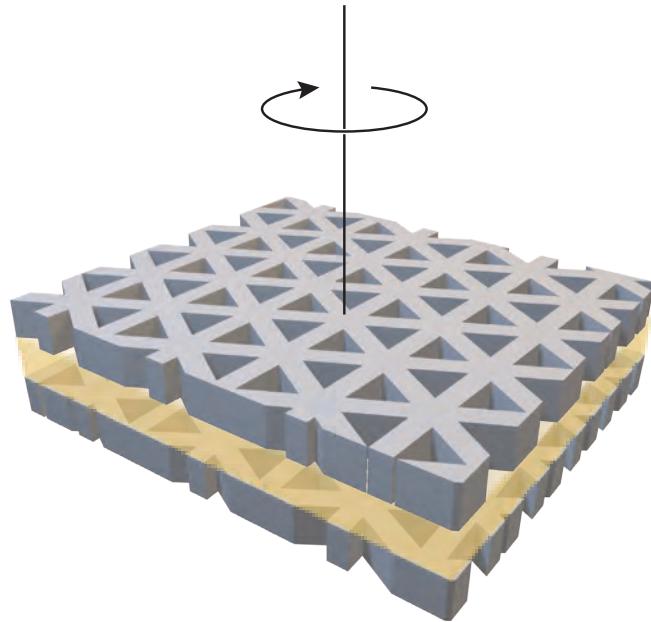


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

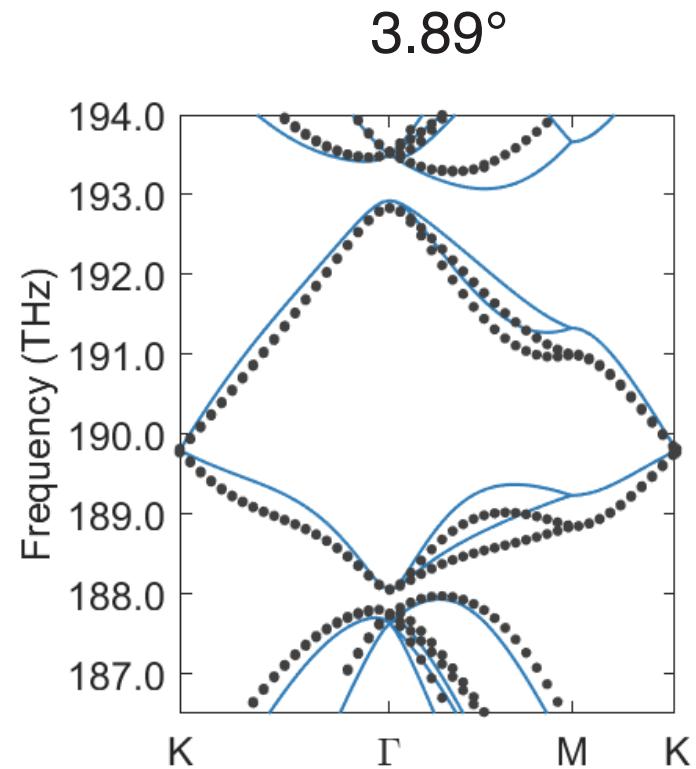
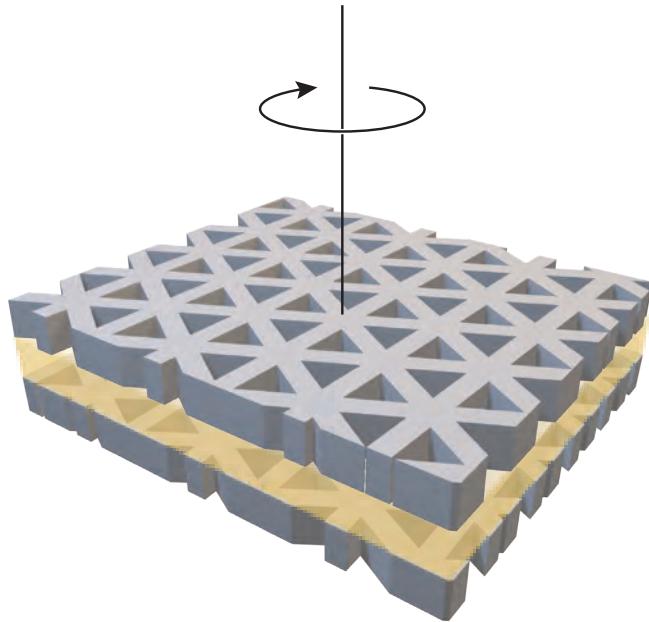


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

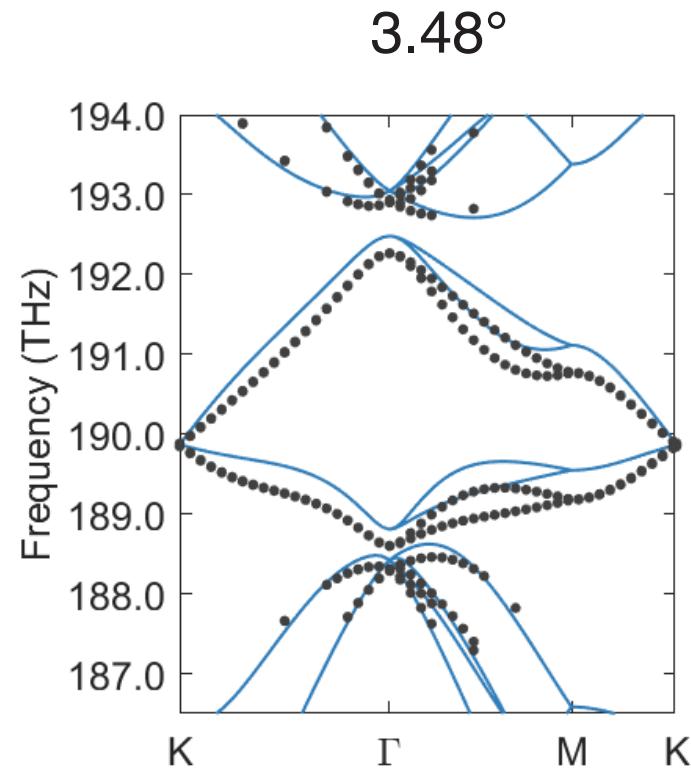
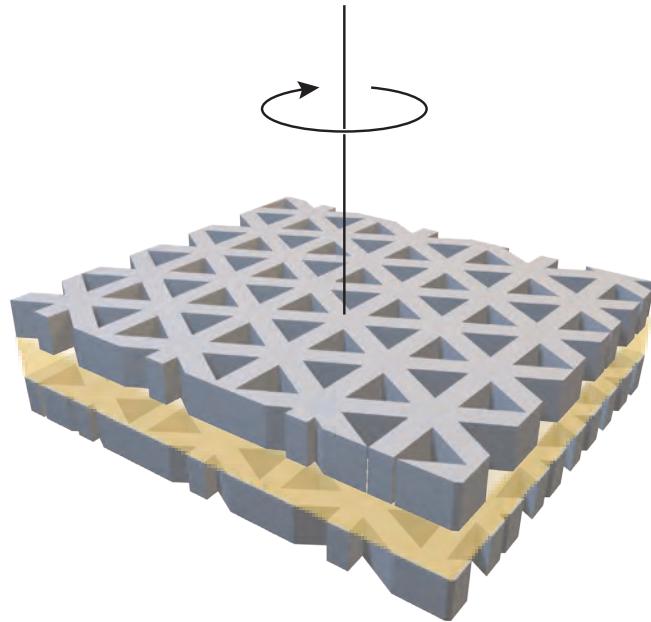


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

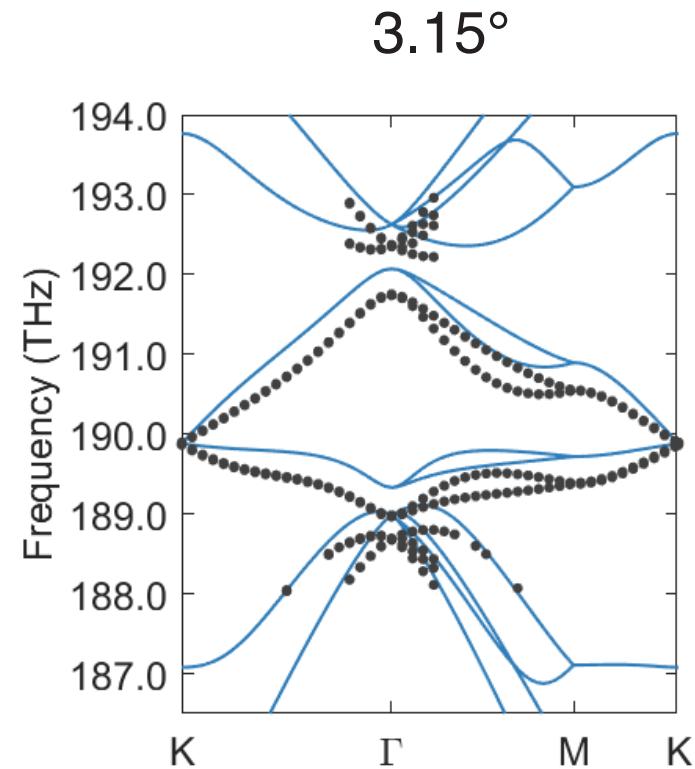
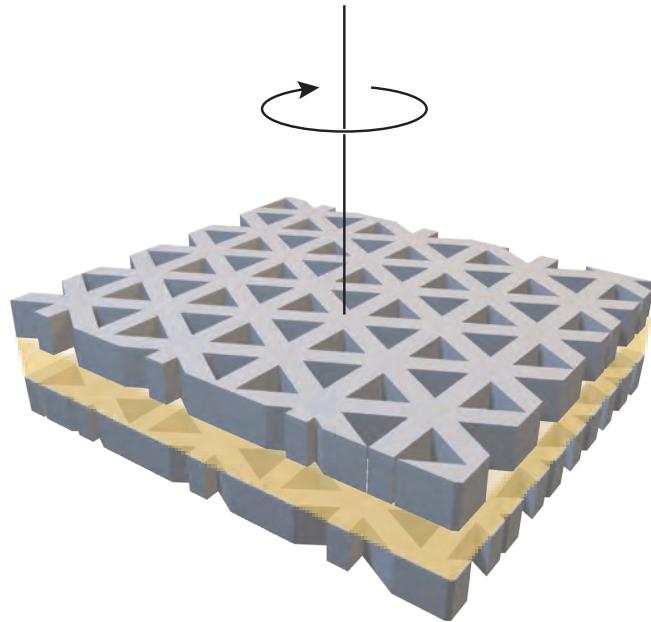


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

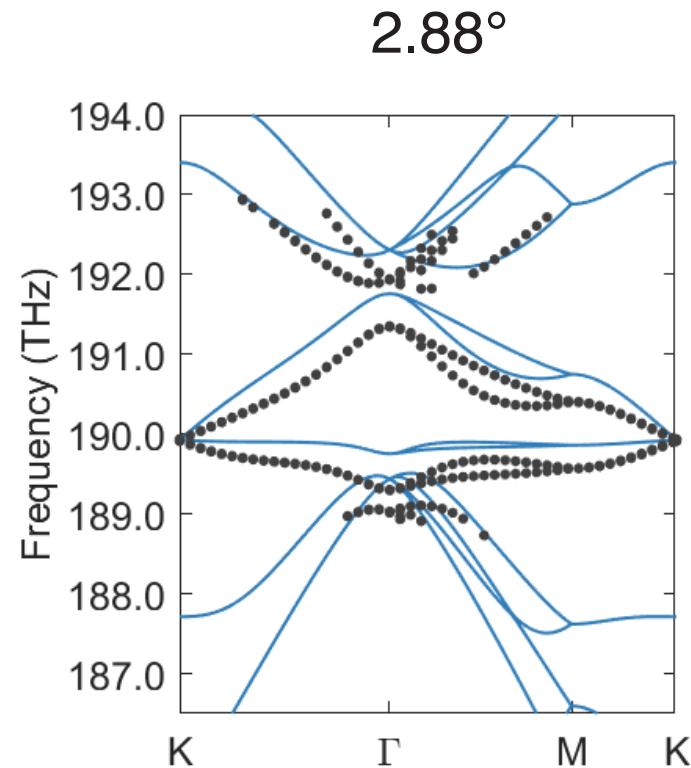
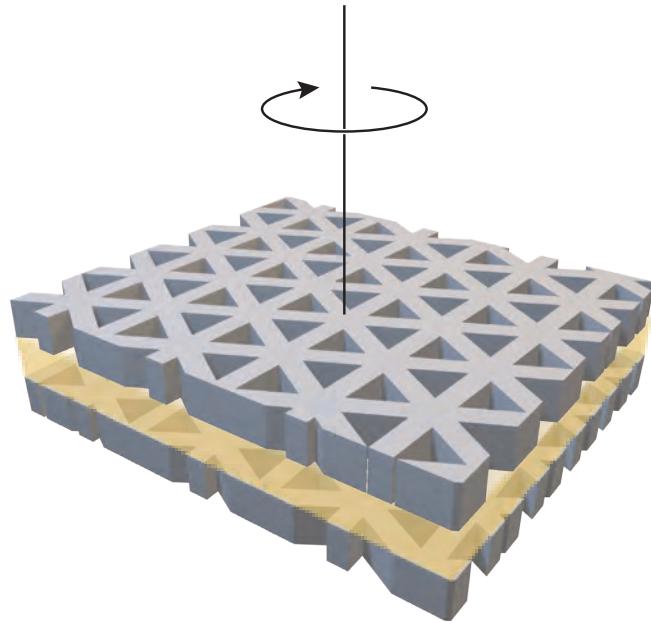


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

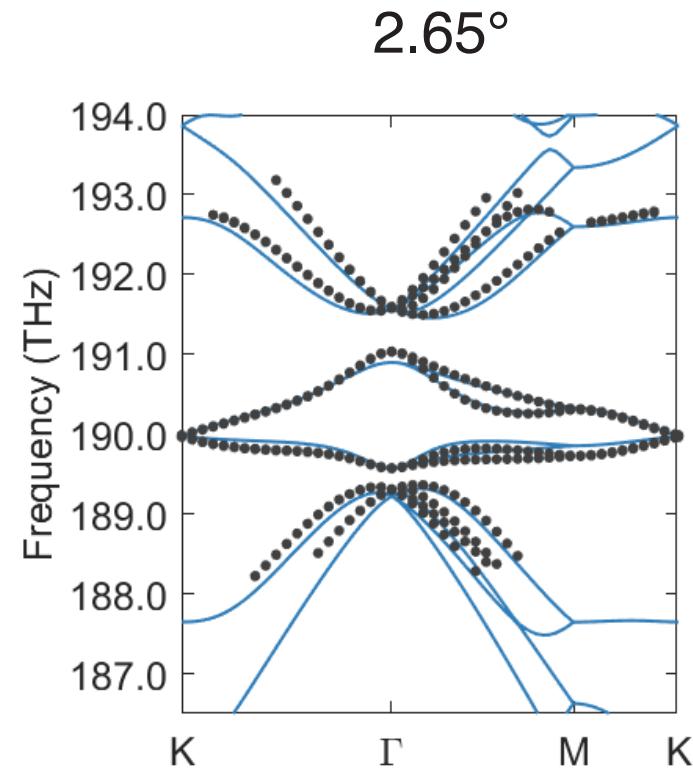
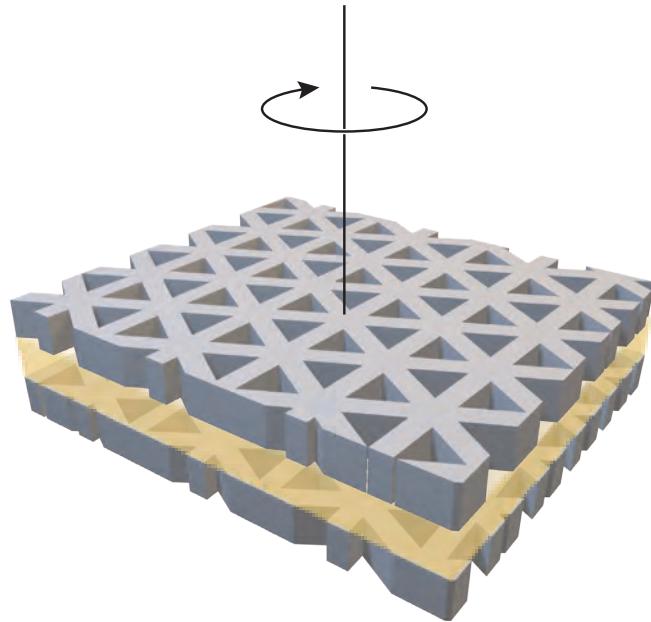


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

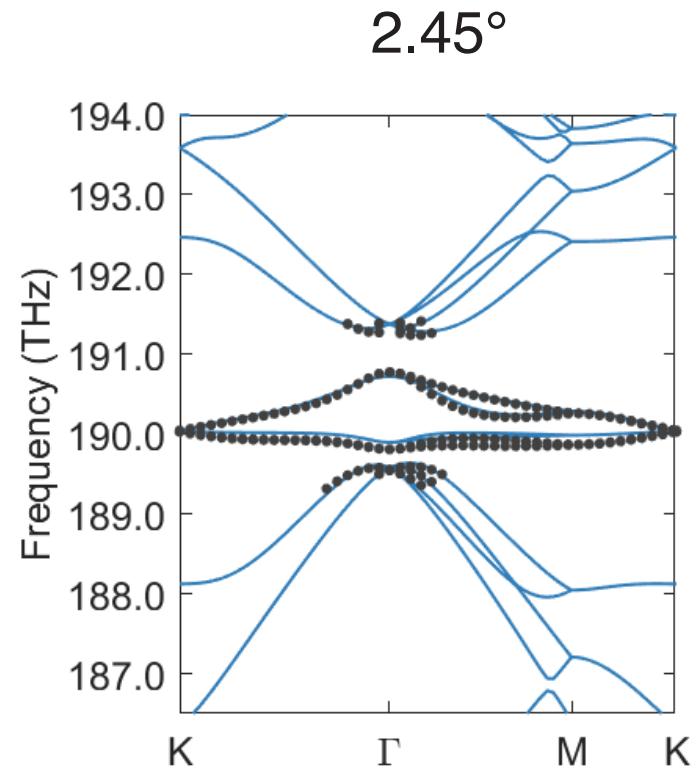
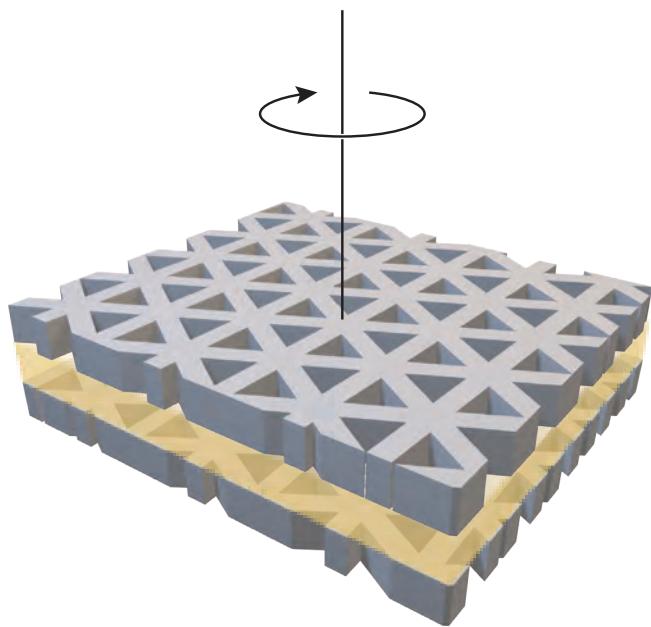


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

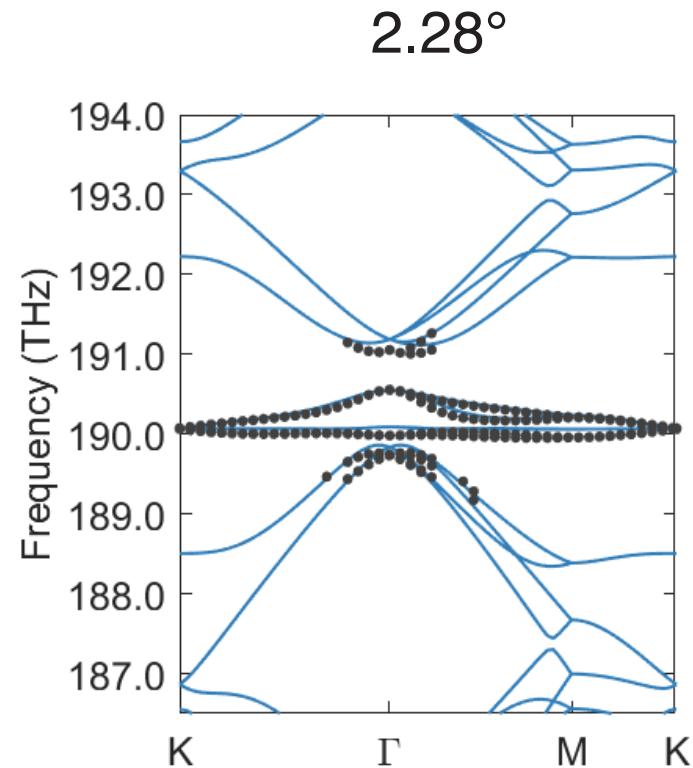
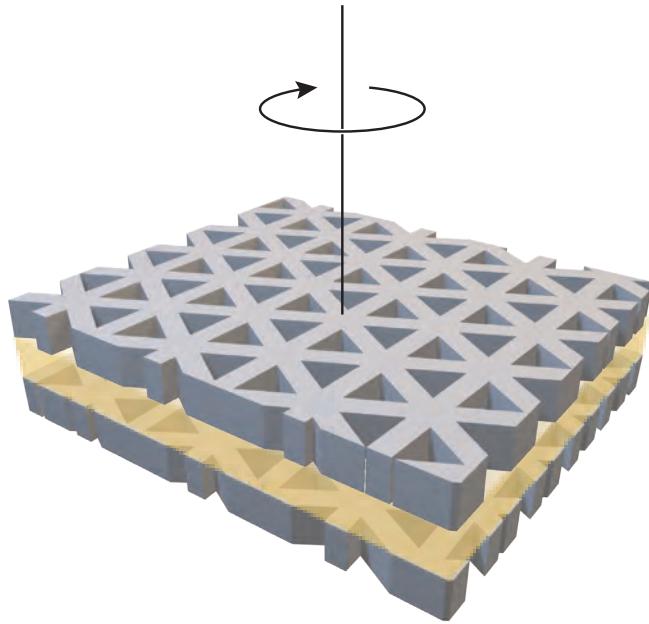


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

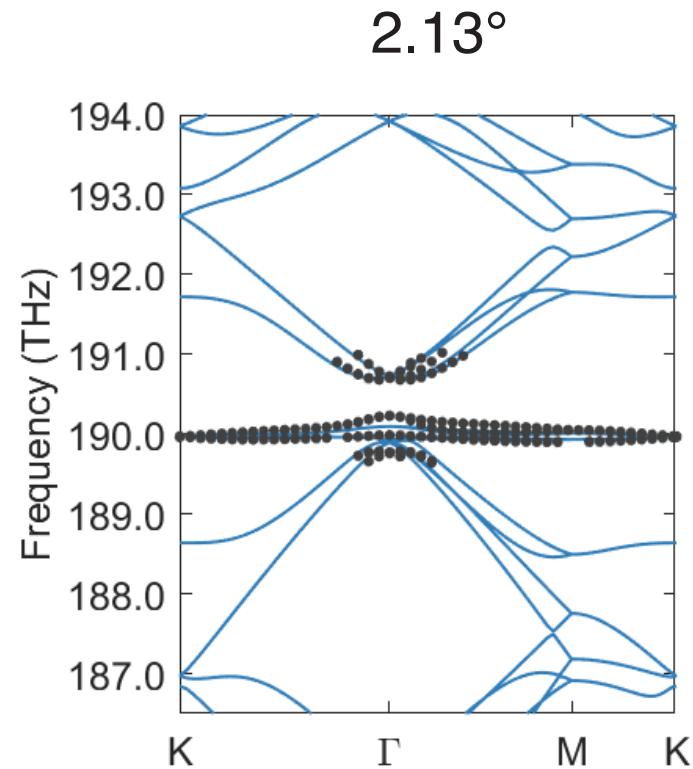
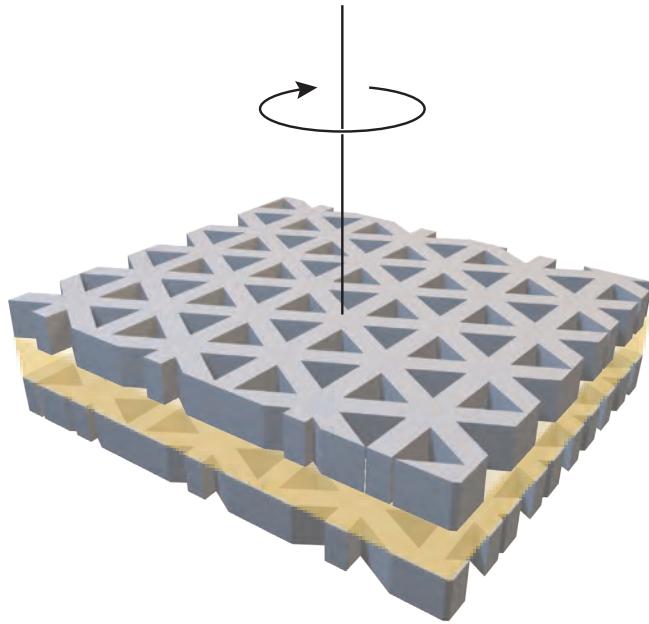


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

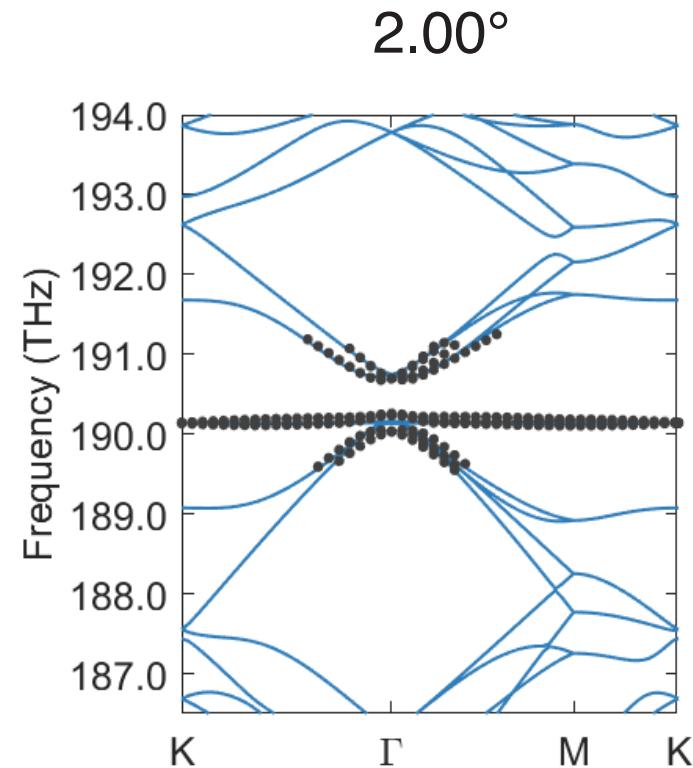
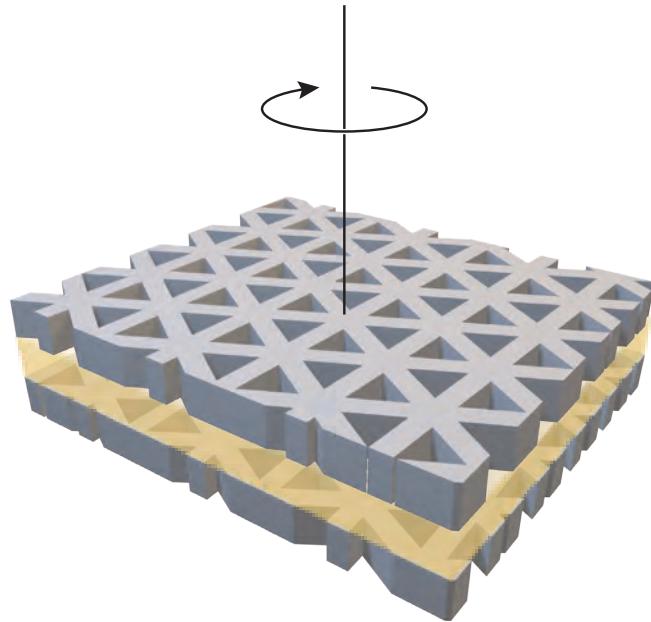


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

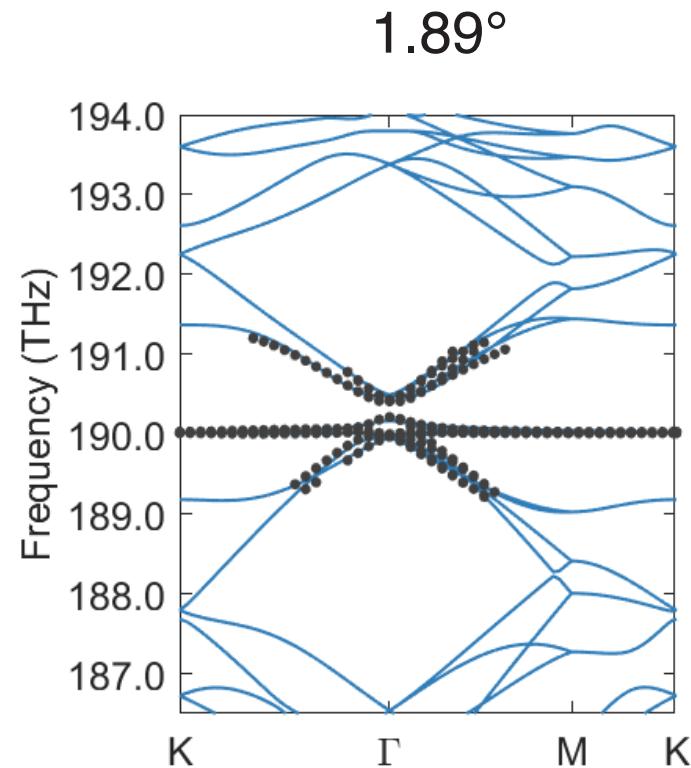
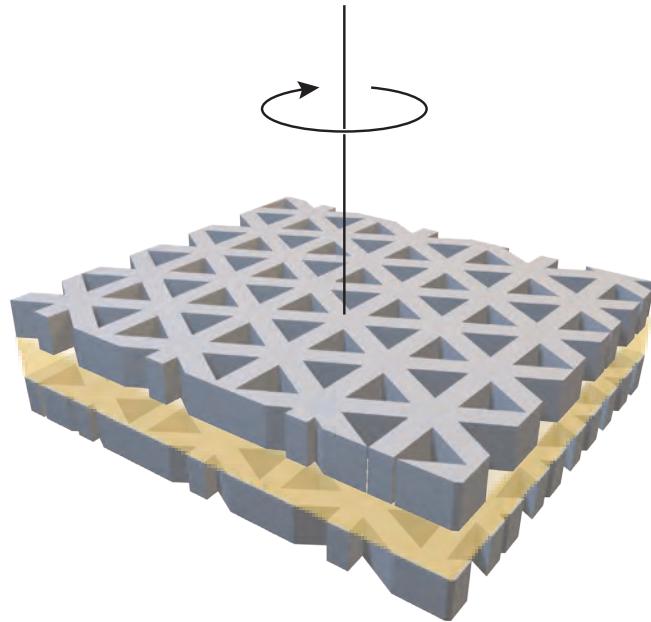


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

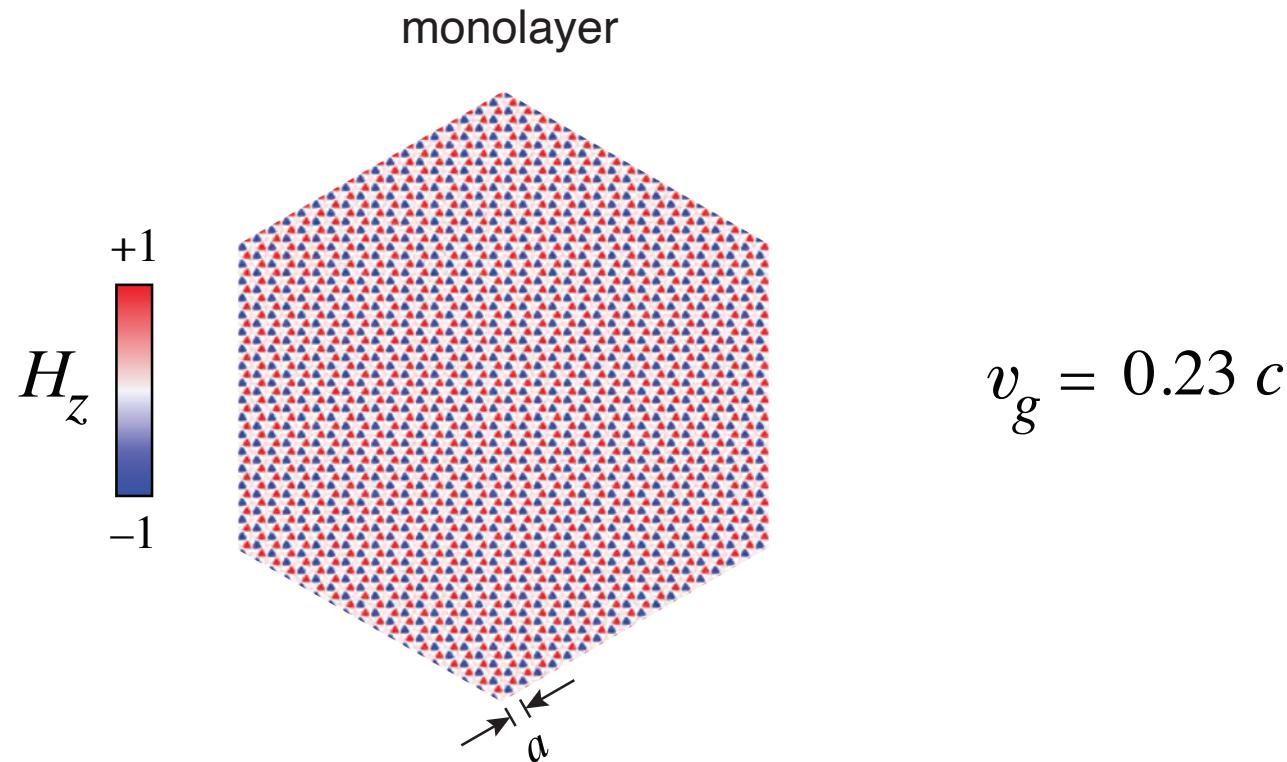


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

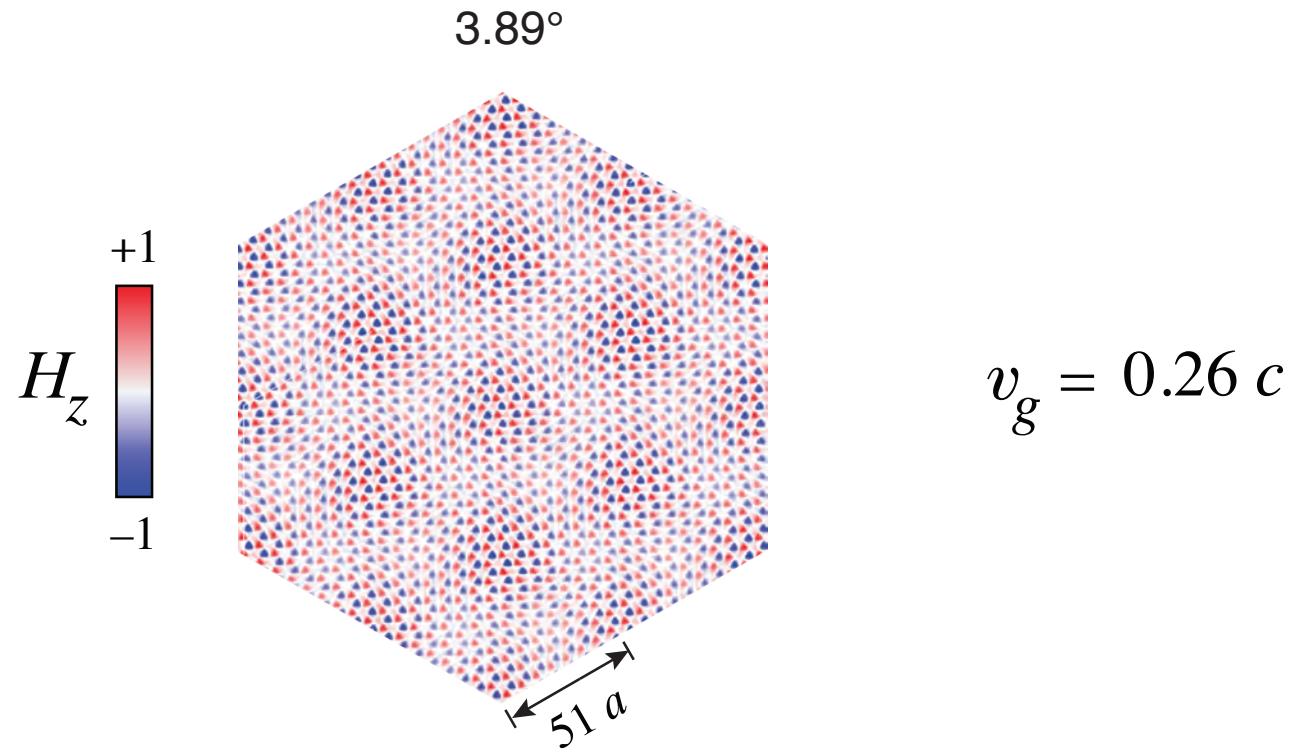


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

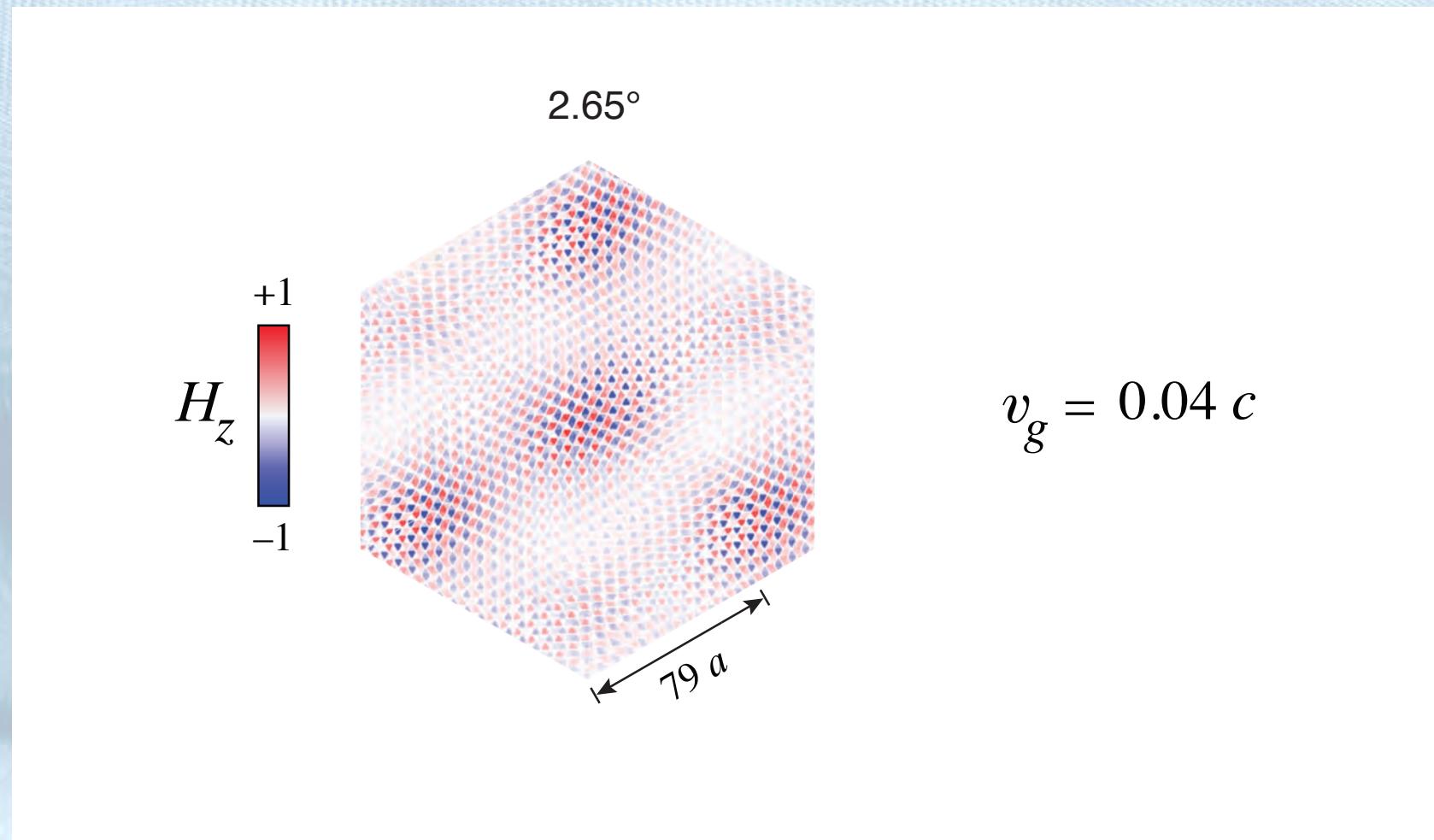


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

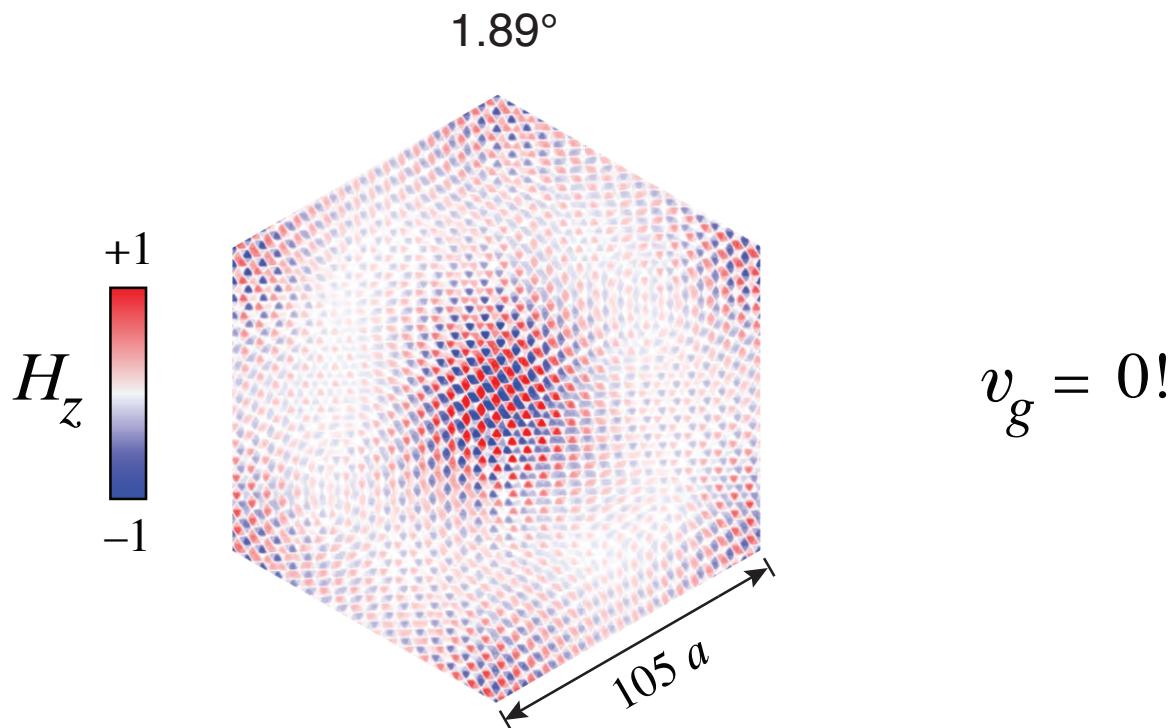


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals

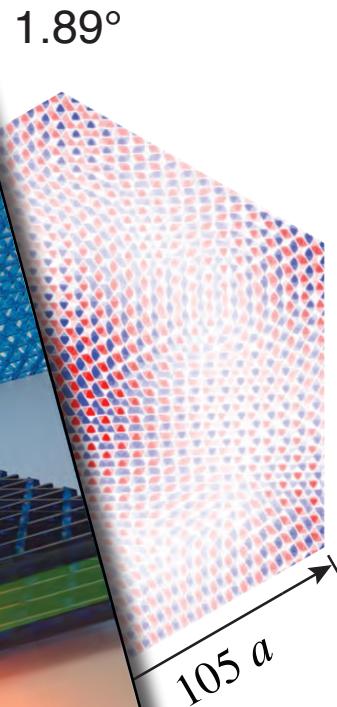


1 index

2 zero index

3 zero group velocity

twisted bilayer photonic crystals



$$v_g = 0!$$

1 index

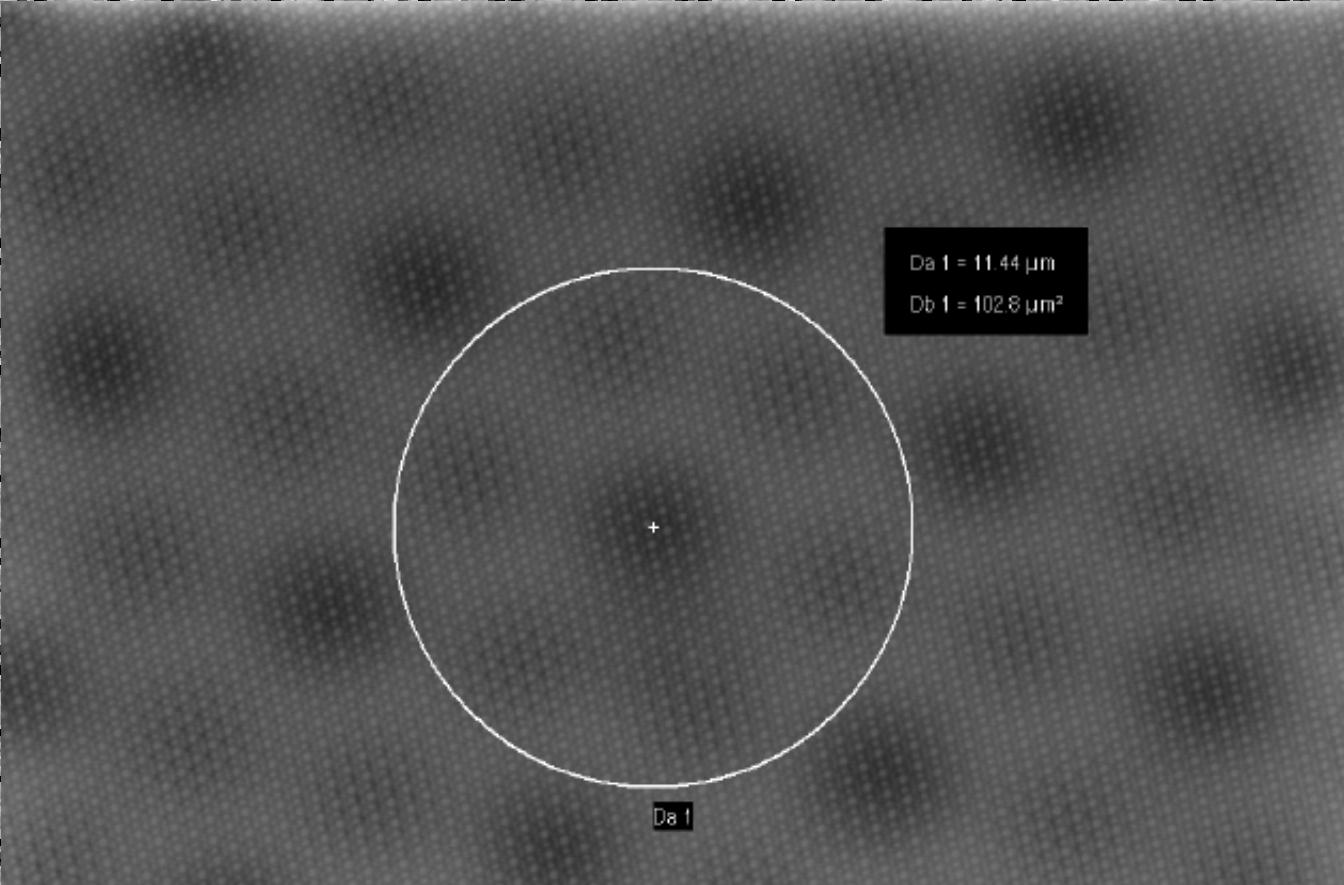
2 zero index

3 zero group velocity

1 index

2 zero index

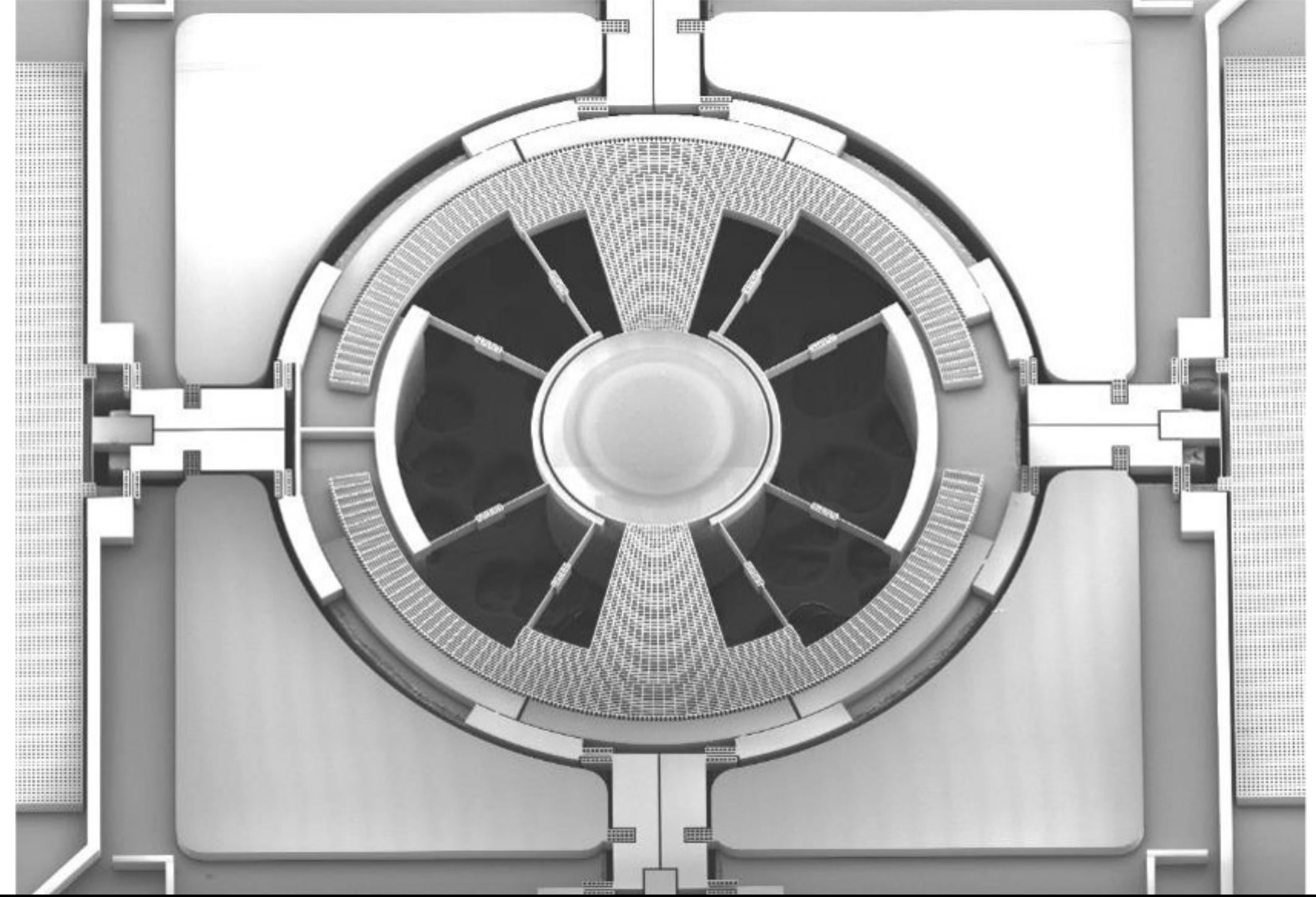
3 zero group velocity



1 index

2 zero index

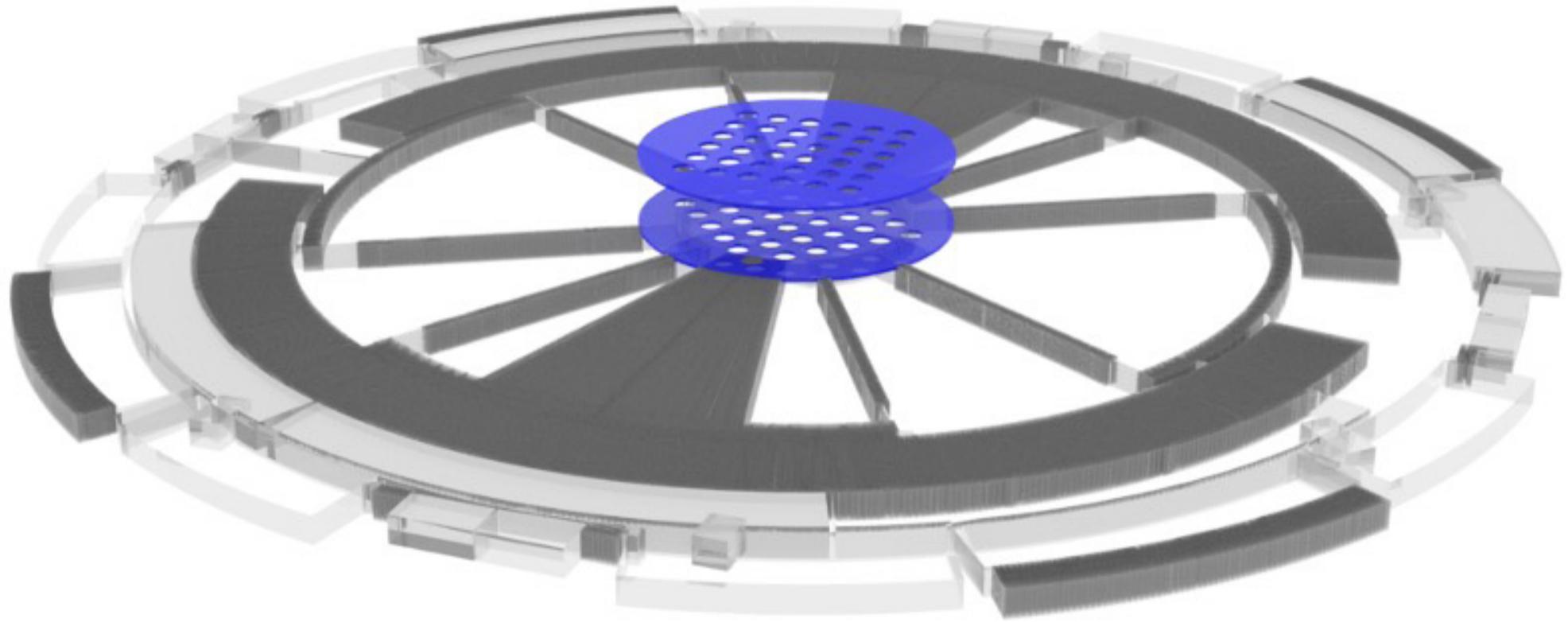
3 zero group velocity



1 index

2 zero index

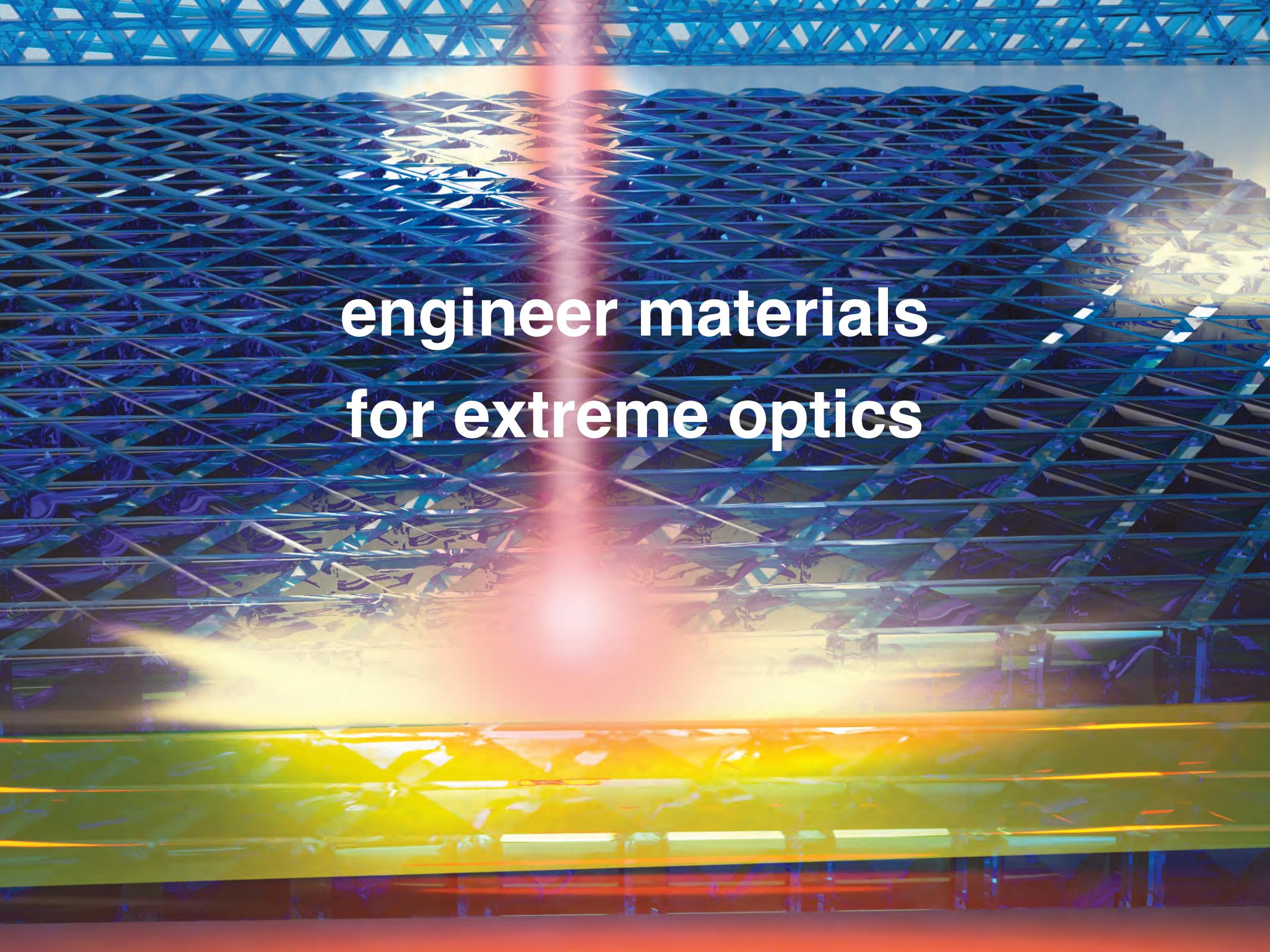
3 zero group velocity



1 index

2 zero index

3 zero group velocity



**engineer materials
for extreme optics**

Zero-Index Metamaterials:

Yang Li
Shota Kita
Phil Muñoz
Orad Reshef,
Daryl Vulis
Mei Yin
Lysander Christakis
Zin Lin, Cleaven Chia
Olivia Mello
Haoning Tang

Zero-Index Waveguide:

Orad Reshef
Justin Gagnon
Marko Loncar
Phil Muñoz
Daryl Vulis

Twisted Bilayer Photonic Crystals

Haoning Tang
Fan Du
Stephen Carr
Xueqi Ni
Huanyu Zhou
Vishantak Srikrishna
Chentong Li
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