Nanophotonics: Linear and nonlinear optics at the nanoscale



NATO-ASI Summer school on Nano-optics Centro Ettore Majorana Erice, Italy, 10 July 2011







Outline

- propagation of pulses
- nonlinear optics
- nanoscale optics
- nonlinear optics at the nanoscale

Outline

- propagation of pulses
- nonlinear optics
- nanoscale optics
- nonlinear optics at the nanoscale

Governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

Governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

Solution:

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

Governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

Solution:

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)}$$

Governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

Solution:

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

where $\vec{E} = \vec{E}_{o}$

$$\vec{E} = \vec{E}_o \ e^{i(kx - \omega t)}$$

In nonferromagnetic media $\mu \approx 1$, and so $n \approx \sqrt{\epsilon}$.

Governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

Solution:
$$\frac{\omega}{k}$$

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$

where

$$\vec{E} = \vec{E}_o \ e^{i(kx - \omega t)}$$

In nonferromagnetic media $\mu \approx 1$, and so $n \approx \sqrt{\epsilon}$.

In dispersive media $n = n(\omega)$.

$$\boldsymbol{\epsilon} = \frac{C_d}{C_o}$$



$$\boldsymbol{\epsilon} = \frac{C_d}{C_o}$$







$$\boldsymbol{\epsilon} = \frac{C_d}{C_o}$$



Alternatively ϵ is measure of attentuation of electric field



Alternatively ϵ is measure of attentuation of electric field



In vacuum:
$$f\lambda = \frac{\omega}{k} = c \implies \omega = c k$$



In medium: $v = \frac{c}{\sqrt{\epsilon(\omega)}} = \frac{c}{n} \implies \omega = \frac{c}{\sqrt{\epsilon}} k$



Which charges participate?















Electron on a string: $F_{binding} = -m_e \omega_o^2 x$

Electron on a string:

$$F_{binding} = -m_e \omega_o^2 x$$

$$F_{damping} = -m_e \gamma \frac{dx}{dt}$$

Electron on a string:

$$F_{binding} = -m_e \omega_o^2 x$$

$$F_{damping} = -m_e \gamma \frac{dx}{dt}$$

$$F_{driving} = -eE = -eE_oe^{-i\omega t}$$

Electron on a string:

$$F_{binding} = -m_e \omega_o^2 x$$

$$F_{damping} = -m_e \gamma \frac{dx}{dt}$$

$$F_{driving} = -eE = -eE_o e^{-i\omega t}$$

Equation of motion:

$$m\frac{d^2x}{dt^2} = \sum F$$

Electron on a string:

$$F_{binding} = -m_e \omega_o^2 x$$

$$F_{damping} = -m_e \gamma \frac{dx}{dt}$$

$$F_{driving} = -eE = -eE_o e^{-i\omega t}$$

Equation of motion:

$$m\frac{d^2x}{dt^2} = \sum F$$

$$m\frac{d^2x}{dt^2} + m\gamma\frac{dx}{dt} + m\omega_o^2 x = -eE$$

Steady state: electron oscillates at driving frequency

$$x(t) = x_o e^{-i\omega t}$$

Steady state: electron oscillates at driving frequency

$$x(t) = x_o e^{-i\omega t} \qquad x_o = -\frac{e}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o$$

Steady state: electron oscillates at driving frequency

$$x(t) = x_o e^{-i\omega t} \qquad x_o = -\frac{e}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o$$

Oscillating dipole

$$p(t) = -ex(t) = \frac{e^2}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o e^{-i\omega t}$$

Steady state: electron oscillates at driving frequency

$$x(t) = x_o e^{-i\omega t} \qquad x_o = -\frac{e}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o$$

Oscillating dipole

$$p(t) = -ex(t) = \frac{e^2}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o e^{-i\omega t}$$

Polarization

$$P(t) = \left(\frac{Ne^2}{m}\right) \sum_{j} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega} E(t) \equiv \epsilon_o \chi_e E(t)$$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 + \frac{Ne^2}{\epsilon_o m} \sum_{j} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega}$$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 + \frac{Ne^2}{\epsilon_o m} \sum_{j=1}^{\infty} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega}$$

Q: For a single resonance, is the value of $\epsilon(\omega)$ at high frequency

- 1. larger than,
- 2. the same as, or
- 3. smaller than the value at low frequency?
Dielectric function

$$\boldsymbol{\epsilon}(\boldsymbol{\omega}) \equiv 1 + \chi_e = 1 + \frac{Ne^2}{\boldsymbol{\epsilon}_o m} \sum_{j}^{n} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega}$$

Q: For a single resonance, is the value of $\epsilon(\omega)$ at high frequency

- 1. larger than,
- 2. the same as, or

3. smaller than the value at low frequency?

Dielectric function

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 + \frac{Ne^2}{\epsilon_o m} \sum_{j} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega}$$



Amplitude of bound charge oscillation



Below resonance: bound charges keep up with driving field \Rightarrow field attenuated, wave propagates more slowly



At resonance: energy transfer from wave to bound charges ⇒ wave attenuates (absorption)



Above resonance: bound charges cannot keep up with driving field \Rightarrow dielectric like a vacuum



Dielectric function

$$\boldsymbol{\epsilon}(\boldsymbol{\omega}) \equiv 1 + \chi_e = 1 + \frac{Ne^2}{\boldsymbol{\epsilon}_o m} \sum_{j}^{n} \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j \omega}$$



No binding:

$$F_{binding} \approx 0$$

No binding:

$$F_{binding} \approx 0$$

Equation of motion:

$$m\frac{d^2x}{dt^2} + m\gamma\frac{dx}{dt} = -eE$$

No binding:

$$F_{binding} \approx 0$$

Equation of motion:

$$m\frac{d^2x}{dt^2} + m\gamma\frac{dx}{dt} = -eE$$

Solution:
$$x(t) = \frac{e}{m} \frac{1}{\omega^2 + i\gamma\omega} E(t)$$
 (no resonance)

No binding:

$$F_{binding} \approx 0$$

Equation of motion:

$$m\frac{d^2x}{dt^2} + m\gamma\frac{dx}{dt} = -eE$$

Solution:
$$x(t) = \frac{e}{m} \frac{1}{\omega^2 + i\gamma\omega} E(t)$$
 (no resonance)

Low frequency ($\omega \ll \gamma$) \Rightarrow current generated

$$J = -Ne \frac{dx}{dt} = \frac{Ne^2}{m} \frac{1}{\gamma - i\omega} E \approx \frac{Ne^2}{m\gamma} E \equiv \sigma E$$

 $\omega \gg \gamma$: σ complex $\Rightarrow J$ out of phase with E

 $\omega \gg \gamma$: σ complex $\Rightarrow J$ out of phase with E

Dipole:

$$p(t) = -ex(t) = -\frac{e^2}{m}\frac{1}{\omega^2 + i\gamma\omega}E(t)$$

 $\omega \gg \gamma$: σ complex $\Rightarrow J$ out of phase with E

Dipole:

$$p(t) = -ex(t) = -\frac{e^2}{m}\frac{1}{\omega^2 + i\gamma\omega}E(t)$$

Polarization:

$$P(t) = -\frac{Ne^2}{m} \frac{1}{\omega^2 + i\gamma\omega} E(t) \equiv \epsilon_o \chi_e E(t)$$

 $\omega \gg \gamma$: σ complex $\Rightarrow J$ out of phase with E

Dipole:

$$p(t) = -ex(t) = -\frac{e^2}{m}\frac{1}{\omega^2 + i\gamma\omega}E(t)$$

Polarization:

$$P(t) = -\frac{Ne^2}{m} \frac{1}{\omega^2 + i\gamma\omega} E(t) \equiv \epsilon_o \chi_e E(t)$$

Dielectric function:

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

Little damping: $\gamma \approx 0 \Rightarrow \epsilon'' = 0$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

Little damping: $\gamma \approx 0 \Rightarrow \epsilon'' = 0$

$$\epsilon'(\omega) = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2} \equiv 1 - \frac{\omega_p^2}{\omega^2}$$

$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

Little damping: $\gamma \approx 0 \Rightarrow \epsilon'' = 0$

$$\epsilon'(\omega) = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2} \equiv 1 - \frac{\omega_p^2}{\omega^2}$$



$$\epsilon(\omega) \equiv 1 + \chi_e = 1 - \frac{Ne^2}{m\epsilon_o} \frac{1}{\omega^2 + i\gamma\omega} = \epsilon'(\omega) + i\epsilon''(\omega)$$

Add damping:

$$\gamma \lesssim \omega_p$$





Plasma acts like a high-pass filter







Plasma acts like a high-pass filter

log <i>N</i> (cm⁻³)	ω_p (rad s ⁻¹)	$oldsymbol{\lambda}_p$
22	6 x 10 ¹⁵	330 nm
18	6 x 10 ¹³	33 µm
14	6 x 10 ¹¹	3.3 mm
10	6 x 10 ⁹	0.33 m



 \bigvee \bigvee $\omega < \omega_p$











Consider two propagating waves:

$$y_1 = A \sin(k_1 x - \omega_1 t)$$
 and $y_2 = A \sin(k_2 x - \omega_2 t)$

Consider two propagating waves:

$$y_1 = A \sin(k_1 x - \omega_1 t)$$
 and $y_2 = A \sin(k_2 x - \omega_2 t)$

propagating at speeds

$$v_1 = \frac{\omega_1}{k_1} = f_1 \lambda_1$$
 and $v_2 = \frac{\omega_2}{k_2} = f_2 \lambda_2$.

Consider two propagating waves:

$$y_1 = A \sin(k_1 x - \omega_1 t)$$
 and $y_2 = A \sin(k_2 x - \omega_2 t)$

propagating at speeds

$$v_1 = rac{\omega_1}{k_1} = f_1 \lambda_1$$
 and $v_2 = rac{\omega_2}{k_2} = f_2 \lambda_2$.

Superposition:

$$y = A[\sin (k_1 x - \omega_1 t) + \sin (k_2 x - \omega_2 t)]$$

Consider two propagating waves:

$$y_1 = A \sin (k_1 x - \omega_1 t)$$
 and $y_2 = A \sin (k_2 x - \omega_2 t)$

propagating at speeds

$$v_1 = rac{\omega_1}{k_1} = f_1 \lambda_1$$
 and $v_2 = rac{\omega_2}{k_2} = f_2 \lambda_2$.

Superposition:

$$y = A[\sin (k_1 x - \omega_1 t) + \sin (k_2 x - \omega_2 t)]$$

$$\sin \alpha + \sin \beta = 2 \cos \left(\frac{\alpha - \beta}{2} \right) \sin \left(\frac{\alpha + \beta}{2} \right)$$

$$y = 2A \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \sin \left[\frac{k_1 + k_2}{2}x - \frac{\omega_1 + \omega_2}{2}t\right]$$

$$y = 2A \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \sin \left[\frac{k_1 + k_2}{2}x - \frac{\omega_1 + \omega_2}{2}t\right]$$

Let: $k_1 - k_2 \equiv \Delta k$ and $\omega_1 - \omega_2 \equiv \Delta \omega$

$$y = 2A \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \sin \left[\frac{k_1 + k_2}{2}x - \frac{\omega_1 + \omega_2}{2}t\right]$$

Let: $k_1 - k_2 \equiv \Delta k$ and $\omega_1 - \omega_2 \equiv \Delta \omega$

$$\frac{k_1+k_2}{2} \equiv k$$
 and $\frac{\omega_1+\omega_2}{2} \equiv \omega$

$$y = 2A \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \sin \left[\frac{k_1 + k_2}{2}x - \frac{\omega_1 + \omega_2}{2}t\right]$$

Let: $k_1 - k_2 \equiv \Delta k$ and $\omega_1 - \omega_2 \equiv \Delta \omega$

$$\frac{k_1+k_2}{2} \equiv k$$
 and $\frac{\omega_1+\omega_2}{2} \equiv \omega$

and so:

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

$$y = 2A \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \sin \left[\frac{k_1 + k_2}{2}x - \frac{\omega_1 + \omega_2}{2}t\right]$$

Let: $k_1 - k_2 \equiv \Delta k$ and $\omega_1 - \omega_2 \equiv \Delta \omega$

$$\frac{k_1+k_2}{2} \equiv k$$
 and $\frac{\omega_1+\omega_2}{2} \equiv \omega$

and so:

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

traveling sine wave, with amplitude modulation.
$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

At
$$t = 0$$
: $y = 2A \cos \frac{1}{2}(x\Delta k) \sin (kx)$



$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

At
$$t = 0$$
:
 $y = 2A \cos \frac{1}{2}(x\Delta k) \sin (kx)$
carrier



$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

At
$$t = 0$$
:
 $y = 2A \cos \frac{1}{2}(x\Delta k) \sin (kx)$
carrier



$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

At
$$t = 0$$
:
 $y = 2A \cos \frac{1}{2}(x\Delta k) \sin (kx)$
envelope carrier



$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

$$y = 2A \cos \frac{1}{2} (x \Delta k - t \Delta \omega) \sin (kx - \omega t)$$

speed of carrier

$$v_p = \frac{\omega}{k} = f\lambda$$

$$y = 2A \cos \frac{1}{2} (x \Delta k - t \Delta \omega) \sin (kx - \omega t)$$

speed of carrier

$$v_p = \frac{\omega}{k} = f\lambda$$

speed of envelope

$$v_g = \frac{\Delta \omega}{\Delta k} = \frac{d\omega}{dk}$$

$$y = 2A \cos \frac{1}{2} (x \Delta k - t \Delta \omega) \sin (kx - \omega t)$$

speed of carrier ('phase velocity'):

$$v_p = \frac{\omega}{k} = f\lambda$$

speed of envelope ('group velocity'):

$$v_g = rac{\Delta \omega}{\Delta k} = rac{d\omega}{dk}$$

let's practice a bit!

(please complete worksheet)

For each wave, determine the wavevector k, the frequency ω , and the propagation speed v:

$$k_{1} = \frac{8.0}{k_{1}} \qquad \text{and} \qquad k_{2} = \frac{72}{0.45} = 7.6 < k_{1}$$
(B)
$$\omega_{1} = \frac{8.0}{k_{1}} \qquad \text{and} \qquad \omega_{2} = 7.2$$

$$v_{1} = \frac{\omega_{1}}{k_{1}} = 1. \qquad \text{and} \qquad v_{2} = \frac{\omega_{2}}{k_{2}} = \frac{7.2}{2.77} = 0.95$$

Does the red get ahead of the blue or the other way around? Why?

Is the dispersion in the medium through which these waves propagate normal or anomalous? Why?

What is the phase velocity of the superposition of y_1 and y_2 ?

$$V_{p} = \frac{(w_{1} + w_{2})/2}{(k_{1} + k_{2})/2} = \frac{7.6}{7.8} = 0.98$$

What is the group velocity of the superposition of y_1 and y_2 ?

$$V_g = \frac{W_1 - W_2}{k_1 - k_2} = \frac{0.8}{0.4} = 2$$

Do the crests of the carrier wave travel forward or backward through the envelope? Does your prediction agree with your observation?

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

If no dispersion

$$v_p = \frac{\omega_1}{k_1} = \frac{\omega_2}{k_2}$$

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

If no dispersion
$$v_p = \frac{\omega_1}{k_1} = \frac{\omega_2}{k_2}$$

group velocity:

$$v_g = \frac{\Delta \omega}{\Delta k} = \frac{\omega_1 - \omega_2}{k_1 - k_2} = \frac{\omega_1 / k_1 - \omega_2 / k_1}{1 - k_2 / k_1} = \frac{v_p - \omega_2 / k_1}{1 - k_2 / k_1}$$

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

If no dispersion
$$v_p = \frac{\omega_1}{k_1} = \frac{\omega_2}{k_2}$$

group velocity:

$$v_g = \frac{\Delta \omega}{\Delta k} = \frac{\omega_1 - \omega_2}{k_1 - k_2} = \frac{\omega_1 / k_1 - \omega_2 / k_1}{1 - k_2 / k_1} = \frac{v_p - \omega_2 / k_1}{1 - k_2 / k_1}$$

$$v_g = \frac{v_p - \omega_2/k_2(k_2/k_1)}{1 - k_2/k_1} = v_p \frac{1 - k_2/k_1}{1 - k_2/k_1} = v_p$$

$$y = 2A \cos \frac{1}{2}(x\Delta k - t\Delta \omega) \sin (kx - \omega t)$$

If no dispersion

$$v_p = \frac{\omega_1}{k_1} = \frac{\omega_2}{k_2}$$

group and phase velocities are the same:

$$v_g = v_p$$

and so the envelope and carrier travel together

$$y = 2A \cos \frac{1}{2} (x \Delta k - t \Delta \omega) \sin (kx - \omega t)$$

Types of dispersion:

$$\frac{dn}{d\omega} > 0$$
 normal dispersion

 $\frac{dn}{d\omega} = 0$ no dispersion

 $\frac{dn}{d\omega} < 0$ anomalous dispersion



$$y = 2A \cos \frac{1}{2} (x \Delta k - t \Delta \omega) \sin (kx - \omega t)$$

Types of dispersion:





medium causes pulse to stretch





medium causes pulse to stretch



compensate by rearranging spectral components!









How do these arrangements work?

How do these arrangements work? (please complete worksheet)







Does path length difference compensate?



grating gives low frequency longer path length!









Does path length difference compensate?



...so prism gives low frequency shorter path length!

consider traveling Gaussian pulse again:

$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$
$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$

Q: Can you tell if the medium is dispersive or not?

$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$

Q: Can you tell if the medium is dispersive or not?

- 1. Yes, it is dispersive
- 2. No, it is not dispersive (pulse shape is constant)
- 3. Cannot tell

$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$

Q: Can you tell if the medium is dispersive or not?

- 1. Yes, it is dispersive
- 2. No, it is not dispersive (pulse shape is constant)
- 3. Cannot tell 🖌

$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$

Q: Can you tell if the medium is dispersive or not?

A: Cannot tell (the medium is dispersive if $v_g \neq \frac{\omega}{k}$)

$$y(t) = \exp\left[-\frac{(x-v_g t)^2}{2\sigma_t^2}\right] \sin(kx - \omega t)$$

Q: Can you tell if the medium is dispersive or not?

A: Cannot tell (the medium is dispersive if $v_g \neq \frac{\omega}{k}$)

...but Gaussian shape of pulse is constant!













only nonlinear dispersion changes pulse shape!



Write dispersion as Taylor series:

Write dispersion as Taylor series:

$$\omega(k) = \omega_o + \left(\frac{d\omega}{dk}\right)_{k=k_o} (k-k_o) + \frac{1}{2} \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o} (k-k_o)^2$$

Write dispersion as Taylor series:

$$\omega(k) = \omega_o + \left(\frac{d\omega}{dk}\right)_{k=k_o} (k-k_o) + \frac{1}{2} \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o} (k-k_o)^2$$

let

$$u \equiv \left(\frac{d\omega}{dk}\right)_{k=k_o}$$
 and $w \equiv \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o}$.

Write dispersion as Taylor series:

$$\omega(k) = \omega_o + \left(\frac{d\omega}{dk}\right)_{k=k_o} (k-k_o) + \frac{1}{2} \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o} (k-k_o)^2$$

let

$$u \equiv \left(\frac{d\omega}{dk}\right)_{k=k_o}$$
 and $w \equiv \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o}$.

group velocity:

$$v_g = \frac{d\omega}{dk} = u + wk$$

Write dispersion as Taylor series:

$$\omega(k) = \omega_o + \left(\frac{d\omega}{dk}\right)_{k=k_o} (k-k_o) + \frac{1}{2} \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o} (k-k_o)^2$$

let

$$u \equiv \left(\frac{d\omega}{dk}\right)_{k=k_o}$$
 and $w \equiv \left(\frac{d^2\omega}{dk^2}\right)_{k=k_o}$.

group velocity:

$$v_g = \frac{d\omega}{dk} = u + wk$$

if w = 0, then group velocity and pulse shape constant!

So not path length but
$$\frac{d^2\phi}{d\omega^2}$$
 matters!

So not path length but
$$\frac{d^2\phi}{d\omega^2}$$
 matters!



So not path length but
$$\frac{d^2\phi}{d\omega^2}$$
 matters!



So not path length but
$$\frac{d^2\phi}{d\omega^2}$$
 matters!

	$rac{dl_{eff}}{d\omega}$	$\frac{d^2\phi}{d\omega^2}$
dispersion	+	+
gratings	-	-
prisms	+	_

So not path length but
$$\frac{d^2\phi}{d\omega^2}$$
 matters!



Outline

- propagation of pulses
- nonlinear optics
- nanoscale optics
- nonlinear optics at the nanoscale

Linear optics:

$$\vec{P} = \chi \vec{E}$$

Linear optics:

$$\vec{P} = \chi \vec{E}$$

Nonlinear polarization:

$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \dots$$

Linear optics:

$$\vec{P} = \chi \vec{E}$$

Nonlinear polarization:

$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \dots$$

and so:

$$P = P^{(1)} + P^{(2)} + P^{(3)} + \dots$$

Linear optics:

$$\vec{P} = \chi \vec{E}$$

Nonlinear polarization:

$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \dots$$

and so:

$$P = P^{(1)} + P^{(2)} + P^{(3)} + \dots$$
$$P^{(2)} \approx P^{(1)} \text{ when } E = E_{at} \approx \frac{e}{a} \text{, and so } \chi^{(n)} \approx \frac{\chi^{(1)}}{E_{at}^{n-1}}$$

Nonlinear polarization can drive new field:

$$\nabla^2 \vec{E} + \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{4\pi}{c^2} \frac{\partial^2 \vec{P}}{\partial t^2}$$

Nonlinear polarization can drive new field:

$$\nabla^2 \vec{E} + \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{4\pi}{c^2} \frac{\partial^2 \vec{P}}{\partial t^2}$$

But even terms disappear in media with inversion symmetry!

$$\vec{P}^{(2)} = \chi^{(2)} : \vec{E}\vec{E}$$

Nonlinear polarization can drive new field:

$$\nabla^2 \vec{E} + \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{4\pi}{c^2} \frac{\partial^2 \vec{P}}{\partial t^2}$$

But even terms disappear in media with inversion symmetry!

$$\vec{P}^{(2)} = \chi^{(2)} : \vec{E}\vec{E}$$

Invert all vectors:

$$-\vec{P}^{(2)} = \chi^{(2)}:(-\vec{E})(-\vec{E})$$

Nonlinear polarization can drive new field:

$$\nabla^2 \vec{E} + \frac{n^2}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{4\pi}{c^2} \frac{\partial^2 \vec{P}}{\partial t^2}$$

But even terms disappear in media with inversion symmetry!

$$\vec{P}^{(2)} = \chi^{(2)} : \vec{E}\vec{E}$$

Invert all vectors:

$$-\vec{P}^{(2)} = \chi^{(2)}:(-\vec{E})(-\vec{E})$$

and so $\chi^{(2)} = -\chi^{(2)} = 0$.

Consider oscillating electric field:

 $E(t) = E e^{i\omega t} + \text{c.c.}$

Consider oscillating electric field:

 $E(t) = E e^{i\omega t} + \text{c.c.}$

Second-order polarization:

$$P^{(2)}(t) = \chi^{(2)}E^2(t) = \frac{1}{2}\chi^{(2)}EE^* + \frac{1}{4}[\chi^{(2)}E^2e^{-2\omega t} + \text{c.c.}]$$

Consider oscillating electric field:

 $E(t) = E e^{i\omega t} + \text{c.c.}$

Second-order polarization:

$$P^{(2)}(t) = \chi^{(2)}E^2(t) = \frac{1}{2}\chi^{(2)}EE^* + \frac{1}{4}[\chi^{(2)}E^2e^{-2\omega t} + \text{c.c.}]$$

Consider oscillating electric field:

 $E(t) = E e^{i\omega t} + \text{c.c.}$

Second-order polarization:

$$P^{(2)}(t) = \chi^{(2)}E^2(t) = \frac{1}{2}\chi^{(2)}EE^* + \frac{1}{4}[\chi^{(2)}E^2e^{-2\omega t} + \text{c.c.}]$$

Physical interpretation:



Can also cause frequency mixing!
Can also cause frequency mixing! Let

$$E(t) = E_1 e^{-i\omega_1 t} + E_2 e^{-i\omega_2 t}$$

Can also cause frequency mixing! Let

$$E(t) = E_1 e^{-i\omega_1 t} + E_2 e^{-i\omega_2 t}$$

Second-order polarization will contain terms with

 $2\omega_1$ (SHG), $2\omega_2$ (SHG), $\omega_1 + \omega_2$ (SFG), $\omega_1 - \omega_2$ (DFG)

Can also cause frequency mixing! Let

$$E(t) = E_1 e^{-i\omega_1 t} + E_2 e^{-i\omega_2 t}$$

Second-order polarization will contain terms with

$$2\omega_1$$
 (SHG), $2\omega_2$ (SHG), $\omega_1 + \omega_2$ (SFG), $\omega_1 - \omega_2$ (DFG)

Physical interpretation:



Linear response:

$$\vec{P} = \chi \vec{E}$$



Linear response:

$$\vec{P} = \chi \vec{E}$$



Nonlinear response:

$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$





Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$





Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$





Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$





Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



٠

Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



Nonlinear response:
$$P^{(2)} = \chi^{(2)} E^2$$



- Nonlinear response: $P^{(2)} = \chi^{(2)} E^2$
- Q: Silicon atoms are arranged in this way. Does bulk silicon generate second harmonic?

- 1. Yes, silicon is not centrosymmetric (as the unit cell shows)
- 2. No, the crystal as a whole is centrosymmetric
- 3. No, any radiation at the second harmonic is absorbed
- 4. Other

- Nonlinear response: $P^{(2)} = \chi^{(2)} E^2$
- Q: Silicon atoms are arranged in this way. Does bulk silicon generate second harmonic?

- 1. Yes, silicon is not centrosymmetric (as the unit cell shows)
- 2. No, the crystal as a whole is centrosymmetric
- 3. No, any radiation at the second harmonic is absorbed
- 4. Other



Third-order polarization: $P^{(3)}(t) = \chi^{(3)}E^3(t)$

Third-order polarization: $P^{(3)}(t) = \chi^{(3)}E^3(t)$

3 frequencies, 3 terms + c.c.: complicated! In general

$$\cos^3\omega t = \frac{1}{4}\cos 3\omega t + \frac{3}{4}\cos \omega t$$

Third-order polarization: $P^{(3)}(t) = \chi^{(3)}E^3(t)$

3 frequencies, 3 terms + c.c.: complicated! In general

$$\cos^3\omega t = \frac{1}{4}\cos 3\omega t + \frac{3}{4}\cos \omega t$$

Intensity dependent term at fundamental frequency:

$$P^{(3)}(t) = \chi^{(3)}E(t)E^*(t)E(t) = \chi^{(3)}I(t)E(t)$$

Third-order polarization: $P^{(3)}(t) = \chi^{(3)}E^3(t)$

3 frequencies, 3 terms + c.c.: complicated! In general

$$\cos^3\omega t = \frac{1}{4}\cos 3\omega t + \frac{3}{4}\cos \omega t$$

Intensity dependent term at fundamental frequency:

$$P^{(3)}(t) = \chi^{(3)} E(t) E^*(t) E(t) = \chi^{(3)} I(t) E(t)$$

and so

so
$$P = P^{(1)} + P^{(3)} = (\chi^{(1)} + \chi^{(3)}I)E \equiv \chi_{eff}E$$

Third-order polarization: $P^{(3)}(t) = \chi^{(3)}E^3(t)$

3 frequencies, 3 terms + c.c.: complicated! In general

$$\cos^3\omega t = \frac{1}{4}\cos 3\omega t + \frac{3}{4}\cos \omega t$$

Intensity dependent term at fundamental frequency:

$$P^{(3)}(t) = \chi^{(3)} E(t) E^*(t) E(t) = \chi^{(3)} I(t) E(t)$$

and so

$$P = P^{(1)} + P^{(3)} = (\chi^{(1)} + \chi^{(3)}I)E \equiv \chi_{eff}E$$

$$n = \sqrt{\epsilon} = \sqrt{1 + \chi_{eff}} \approx \sqrt{1 + \chi^{(1)}} + \frac{1}{2} \frac{\chi^{(3)}I}{\sqrt{1 + \chi^{(1)}}} = n_o + n_2 I$$

$$n = n_o + n_2 I$$



$$n = n_o + n_2 I$$



$$n = n_o + n_2 I$$



$$n = n_o + n_2 I$$



$$n = n_o + n_2 I$$



Phase:

$$\frac{\phi}{2\pi} = \frac{nL}{\lambda}$$

Phase:

 $\frac{\phi}{2\pi} = \frac{nL}{\lambda}$

$$\phi = \frac{2\pi}{\lambda} L(n_o + n_2 I)$$

Phase:
$$\frac{\phi}{2\pi} = \frac{nL}{\lambda}$$
 $\phi = \frac{2\pi}{\lambda} L(n_o + n_2 I)$
 $d\phi$

Frequency change:

$$\Delta \omega = -\frac{d\phi}{dt}$$
Phase:
$$\frac{\phi}{2\pi} = \frac{nL}{\lambda}$$
 $\phi = \frac{2\pi}{\lambda} L(n_o + n_2 I)$
Frequency change: $\Delta \omega = -\frac{d\phi}{dt}$

- **Q:** Sketch the time dependence of the frequency shift for a Gaussian pulse and determine which is true (assume $n_2 > 0$):
 - 1. Leading edge is blue shifted, trailing edge red shifted
 - 2. Leading and trailing edge blue shifted, center red shifted
 - 3. Leading edge is red shifted, trailing edge blue shifted
 - Leading and trailing edge red shifted, center blue shifted
 Other

Phase:
$$\frac{\phi}{2\pi} = \frac{nL}{\lambda}$$
 $\phi = \frac{2\pi}{\lambda} L(n_o + n_2 I)$
Frequency change: $\Delta \omega = -\frac{d\phi}{dt}$

Q: Sketch the time dependence of the frequency shift for a Gaussian pulse and determine which is true (assume $n_2 > 0$):

1. Leading edge is blue shifted, trailing edge red shifted

- 2. Leading and trailing edge blue shifted, center red shifted
- 3. Leading edge is red shifted, trailing edge blue shifted 🖌

Leading and trailing edge red shifted, center blue shifted
 Other

Phase:

$$\frac{\phi}{2\pi} = \frac{nL}{\lambda}$$

$$b = \frac{2\pi}{\lambda} L(n_o + n_2 I)$$

Frequency change:

$$\Delta \omega = -\frac{d\phi}{dt} = \frac{-2\pi}{\lambda} Ln_2 \frac{dI}{dt}$$





Intensity-dependent index of refraction:

$$n = n_o + n_2 I$$



Intensity-dependent index of refraction:

$$n = n_o + n_2 I$$



self-focusing



but susceptibility is complex!

susceptibility	real part	imaginary part
linear	refraction	absorption
nonlinear	SHG, SFG, DFG, THG,	multiphoton absorption

$$\alpha = \alpha_o + \beta I + \gamma I^2 + \dots$$

Outline

- propagation of pulses
- nonlinear optics
- nanoscale optics
- nonlinear optics at the nanoscale



two crossed planar waves...



... cause an interference pattern



E = 0 on the nodal lines



...satisfying boundary conditions for planar-mirror waveguide





transverse standing wave, traveling along axis





transverse standing wave, traveling along axis













































































boundary conditions only satisfied for certain θ



standing wave in y-direction, traveling in z-direction



consider wave incident at angle θ





twice-reflected wave


self consistency:

$$AC - AB = 2d \sin \theta = m\lambda \quad (m = 1, 2,)$$



self consistency:

$$AC - AB = 2d \sin\theta = m\lambda$$
 $(m = 1, 2,)$
 $\sin\theta_m = m \frac{\lambda}{2d}$



self consistency:

$$AC - AB = 2d \sin\theta = m\lambda$$
 $(m = 1, 2,)$
 $\sin\theta_m = m \frac{\lambda}{2d}$



self consistency:

$$AC - AB = 2d \sin\theta = m\lambda$$
 (m = 1, 2,)
 $\sin\theta_m = m \frac{\lambda}{2d}$



self consistency:

$$AC - AB = 2d \sin \theta = m\lambda$$
 $(m = 1, 2,)$
 $\sin \theta_m = m \frac{\lambda}{2d}$



number of modes:

$$M = \frac{2d}{\lambda}$$





now consider a planar dielectric waveguide



rays incident at angle $\theta > \pi/2 - \theta_c$ are unguided



rays incident at angle $\theta < \pi/2 - \theta_c$ are guided



rays incident at angle $\theta < \pi/2 - \theta_c$ are guided



self consistency:

$$AC - AB = 2d \sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$



self consistency:

$$AC - AB = 2d \sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

SO:

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$

1 10





self consistency:

$$AC - AB = 2d\sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$



self consistency:

$$AC - AB = 2d\sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$



self consistency:

$$AC - AB = 2d\sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$



self consistency:

$$AC - AB = 2d\sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$



self consistency:

$$AC - AB = 2d\sin\theta - \frac{\varphi_r}{\pi}\lambda = m\lambda \quad (m = 0, 1, 2...)$$

$$\tan\left(\frac{\pi d}{\lambda}\sin\theta - m\frac{\pi}{2}\right) = \left(\frac{\sin^2(\pi/2 - \theta_c)}{\sin^2\theta} - 1\right)^{1/2}$$



number of modes:

$$M \doteq \frac{\sin(\pi/2 - \theta_c)}{\lambda/2d}$$



number of modes:

$$M \doteq \frac{\sin(\pi/2 - \theta_c)}{\lambda/2d}$$

or:

$$M \doteq 2 \frac{d}{\lambda} (n_1^2 - n_2^2)^{1/2}$$



propagation constant of guided wave:

$$\beta_m^2 = k^2 - k_y^2 = k^2 - \frac{m^2 \pi^2}{d^2}$$

group velocity:

$$v_m = c \cos \theta_m$$



single mode condition for 600-nm light:

planar mirror
$$M = \frac{2d}{\lambda}$$
 $300 < d < 600 \text{ nm}$

dielectric
$$M \doteq 2 \frac{d}{\lambda} (n_1^2 - n_2^2)^{1/2}$$
 $d < 268 \text{ nm}$



single mode condition for 600-nm light:

planar mirror
$$M = \frac{2d}{\lambda}$$
 $300 < d < 600 \text{ nm}$

dielectric
$$M \doteq 2 \frac{d}{\lambda} (n_1^2 - n_2^2)^{1/2}$$
 $d < 268 \text{ nm}$

can make *d* larger by making $n_1 - n_2$ smaller!



$$\nabla^2 \vec{A} + \omega^2 \mu_o \epsilon \vec{A} = -i\omega\mu_o \nabla \epsilon \Phi$$



$$\nabla^2 \vec{A} + \omega^2 \mu_o \epsilon \vec{A} = 0$$



$$\nabla^2 \vec{A} + \omega^2 \mu_o \epsilon \vec{A} = 0$$

Substituting

$$\vec{A} = \hat{y}u(x,y)e^{-i\beta z}$$



$$\nabla^2 \vec{A} + \omega^2 \mu_o \epsilon \vec{A} = 0$$

Substituting

$$\vec{A} = \hat{y}u(x,y)e^{-i\beta z}$$

yields:

$$\nabla_T^2 u + \left[-\beta^2 + \omega^2 \mu \epsilon(r)\right] u = 0$$



$$\nabla^2 \vec{A} + \omega^2 \mu_o \epsilon \vec{A} = 0$$

Substituting

$$\vec{A} = \hat{y}u(x,y)e^{-i\beta z}$$

yields:

$$\nabla_T^2 u + \left[-\beta^2 + \omega^2 \mu \epsilon(r)\right] u = 0$$

Compare to time-independent Schrödinger equation:

$$\nabla^2 \psi + \frac{2m}{\hbar^2} [E - V(r)] \psi = 0$$





















single mode condition for 600-nm light:

$$M \doteq 2 \frac{d}{\lambda} (n_1^2 - n_2^2)^{1/2}$$

without cladding:

d < 268 nm
Waveguiding

single mode condition for 600-nm light:

$$M \doteq 2 \frac{d}{\lambda} (n_1^2 - n_2^2)^{1/2}$$

without cladding: d < 268 nm

Add cladding with 0.4% index difference:

 $d < 5 \ \mu m$

Waveguiding

commercial single-mode fiber (Corning Titan[®])



operating wavelength: $\lambda = 1310 \text{ nm}/1550 \text{ nm}$



drawbacks of clad fibers:

- weak confinement
- no tight bending
- coupling requires splicing

standard fiber	



















Nature, 426, 816 (2003)







States and





















Waveguiding

Specifications

diameter D:	down to 20 nm
length L:	up to 90 mm
aspect ratio <i>D/L</i> :	up to 10 ⁶
diameter uniformity $\Delta D/L$:	2 x 10 ⁻⁶

Nature, 426, 816 (2003)



240-nm wire



RMS roughness < 0.5 nm











Optical properties

coupling light into nanowires



Optical properties

coupling light into nanowires



Optical properties

coupling light into nanowires


280-nm nanowire

360 nm

(all flight and a standard of the standard of the standard and a second and a second standard standard for the standard for the standard standard standard for the standard s







Poynting vector profile for 800-nm nanowire



Poynting vector profile for 800-nm nanowire



Poynting vector profile for 800-nm nanowire



Poynting vector profile for 600-nm nanowire



Poynting vector profile for 500-nm nanowire



Poynting vector profile for 400-nm nanowire



Poynting vector profile for 300-nm nanowire



Poynting vector profile for 200-nm nanowire



Waveguiding

fraction of power carried in core





















"tunneling" of light



"tunneling" of light



intensity distribution









minimum bending radius: 5.6 μm



virtually no loss through 5 µm corner!





aerogel

420 nm

420 nm

Nanoletters, 5, 259 (2005)



in

out

Nanoletters, 5, 259 (2005)



Nanoletters, 5, 259 (2005)

use tapered fibers to couple light to nanoscale objects

ZnO:non-toxic, wide bandgap semiconductor

vapor transport grown ZnO nanowires





80–400 nm diameter, up to 80 µm long

best of both worlds

ZnO	silica
bottom-up	top-down
semiconductor	glass
active photonic devices	passive waveguides
electrical operation	link to macroworld


















FDTD simulation



ab-initio.mit.edu/wiki/index/Meep

coupling efficiency



coupling efficiency







transmission spectrum



Nano Lett., 7, 3675 (2007)

transmission spectrum



transmission spectrum



transmission spectrum







large diameter: multimode



small diameter: single mode

Points to keep in mind:

- low-loss guiding
- convenient evanescent coupling
- attached to ordinary fiber

Outline

- propagation of pulses
- nonlinear optics
- nanoscale optics
- nonlinear optics at the nanoscale

strong confinement — high intensity





mode field diameter (λ = 800 nm)



M.A. Foster, et al., Optics Express, 12, 2880 (2004)

mode field diameter (λ = 800 nm)



M.A. Foster, et al., Optics Express, 12, 2880 (2004)

nonlinear parameter



M.A. Foster, et al., Optics Express, 12, 2880 (2004)

dispersion important!

dispersion:

- modal dispersion
- material dispersion
- waveguide dispersion
- nonlinear dispersion

waveguide dispersion



waveguide dispersion



waveguide dispersion



waveguide dispersion



waveguide dispersion


waveguide dispersion



waveguide dispersion



waveguide dispersion



waveguide dispersion



waveguide dispersion



nanowire continuum generation



nanowire continuum generation



nanowire continuum generation



nanowire continuum generation



nanowire continuum generation



nanowire continuum generation



nanowire continuum generation



energy in nanowire \approx 1 nJ!

- nanojoule nonlinear optics
- optimum diameter for silica 500–600 nm
- low dispersion











output = transmitted cw + ccw power



input electric field amplitude E_{in}



coupling parameter: ρ



phase accumulation over path length of loop L



coupling parameter: ρ



output is sum of transmitted cw and ccw



accumulated phase:

$$\phi = k_o n L$$

accumulated phase:

$$\phi = k_o n L$$

nonlinear index:

$$n = n_o + n_2 I = n_o + n_2 \frac{P_i}{A_{eff}}$$

accumulated phase:

$$\phi = k_o n L$$

nonlinear index:

$$n = n_o + n_2 I = n_o + n_2 \frac{P_i}{A_{eff}}$$

nonlinear parameter:

$$\gamma = n_2 \frac{k_o}{A_{eff}}$$

power-dependent output:

$$\frac{E_{out}^2}{E_{in}^2} = 1 - 2\rho(1-\rho)\{1 + \cos[(1-2\rho)\gamma P_o L]\}$$

power-dependent output:

$$\frac{E_{out}^2}{E_{in}^2} = 1 - 2\rho(1-\rho)\{1 + \cos[(1-2\rho)\gamma P_o L]\}$$

Q: What happens for a 50-50 coupler?

- **1. All the light is transmitted.**
- 2. Half the light is transmitted.
- 3. No light is transmitted.
- 4. The transmission depends on the input power P_{o} .
- 5. Other

power-dependent output:

$$\frac{E_{out}^2}{E_{in}^2} = 1 - 2\rho(1-\rho)\{1 + \cos[(1-2\rho)\gamma P_o L]\}$$

Q: What happens for a 50-50 coupler?

- **1. All the light is transmitted.**
- 2. Half the light is transmitted.
- 3. No light is transmitted.
- 4. The transmission depends on the input power P_{o} .
- 5. Other

power-dependent output:

$$\frac{E_{out}^2}{E_{in}^2} = 1 - 2\rho(1-\rho)\{1 + \cos[(1-2\rho)\gamma P_o L]\}$$

for 50-50 coupler:

$$\rho = 0.5$$

power-dependent output:

$$\frac{E_{out}^2}{E_{in}^2} = 1 - 2\rho(1-\rho)\{1 + \cos[(1-2\rho)\gamma P_o L]\}$$

for 50-50 coupler:

$$\rho = 0.5$$

no transmission:

$$\frac{E_{out}^2}{E_{in}^2} = 0$$

when $\rho \neq 0.5$:



nonlinear nanogate



nonlinear nanogate






























for NAND gate need ouput with no input



for NAND gate need ouput with no input



for NAND gate need ouput with no input



universal NAND gate



universal NAND gate



universal NAND gate



mesoporous silica

Sagnac loop



output

mesoporous silica

Sagnac loop

very preliminary data



light-by-light modulation!

very preliminary data



very preliminary data













- several nanodevices demonstrated
- large γ permits miniature Sagnac loops
- switching energy \approx 100 pJ







Funding:

Harvard Center for Imaging and Mesoscopic Structures National Science Foundation National Natural Science Foundation of China

for a copy of this presentation:

http://mazur.harvard.edu



eric_mazur



Google Search I'm Feeling Luck



mazur		
	mazur	

Google Search Thi reeling	Lucky
---------------------------	-------



mazur		




mazur		



Funding:

Harvard Center for Imaging and Mesoscopic Structures National Science Foundation National Natural Science Foundation of China

for a copy of this presentation:

http://mazur.harvard.edu



eric_mazur