Reflections





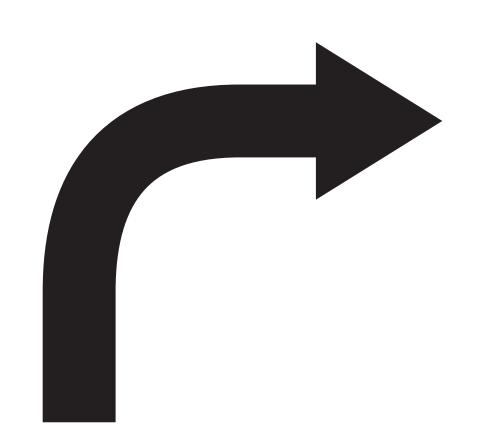
Reflections







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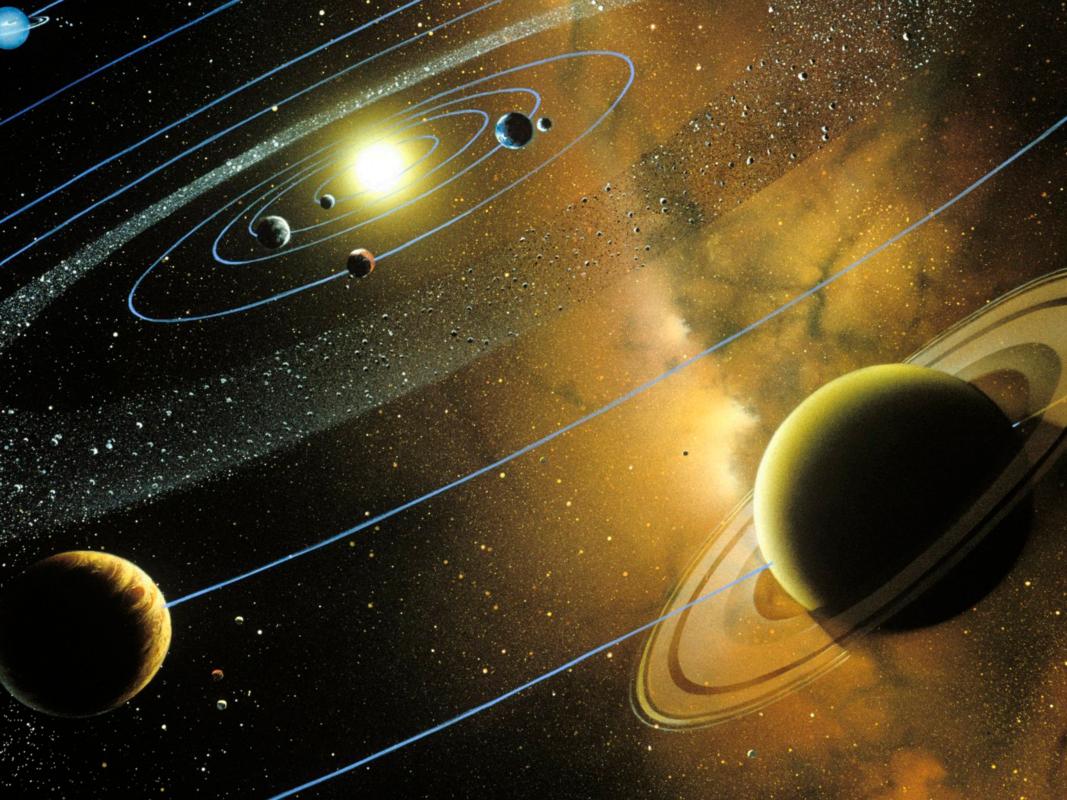












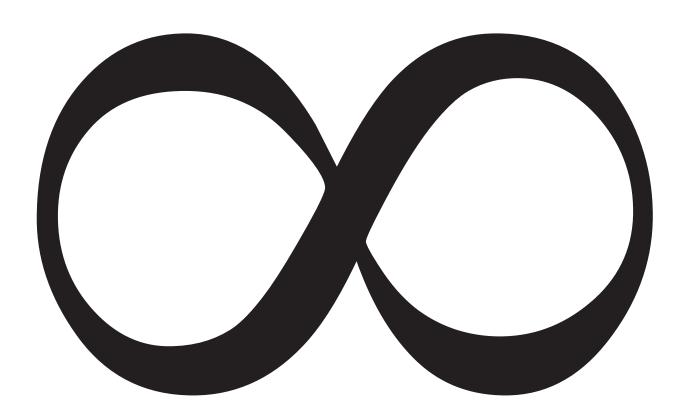
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Why am 11?









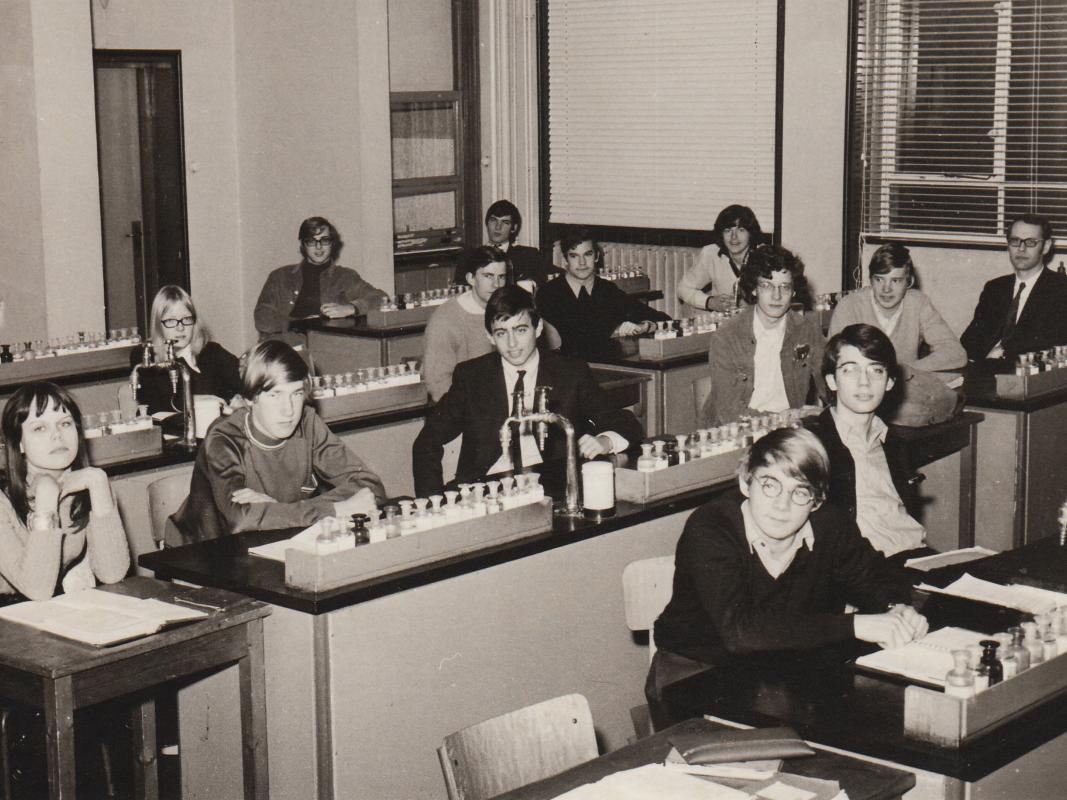


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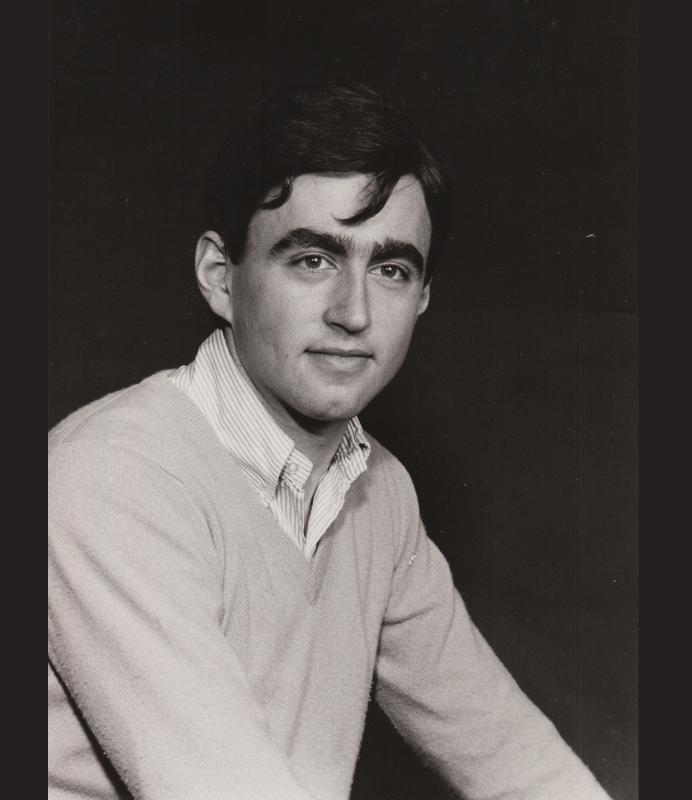


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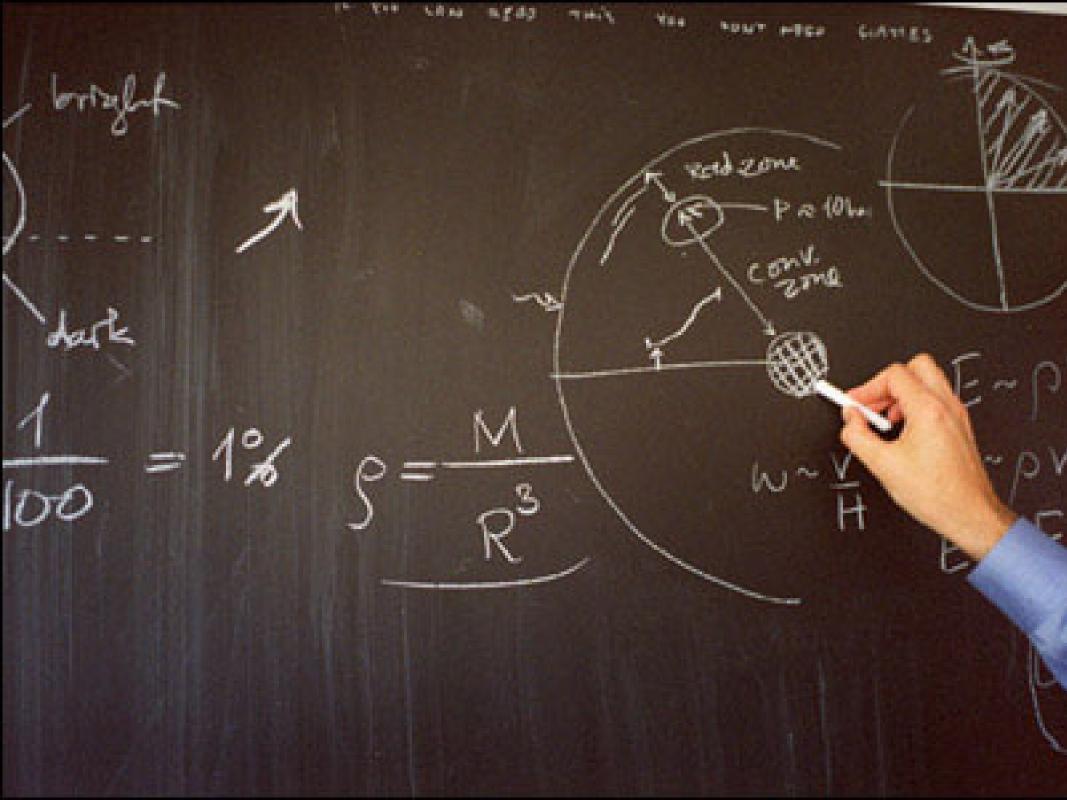
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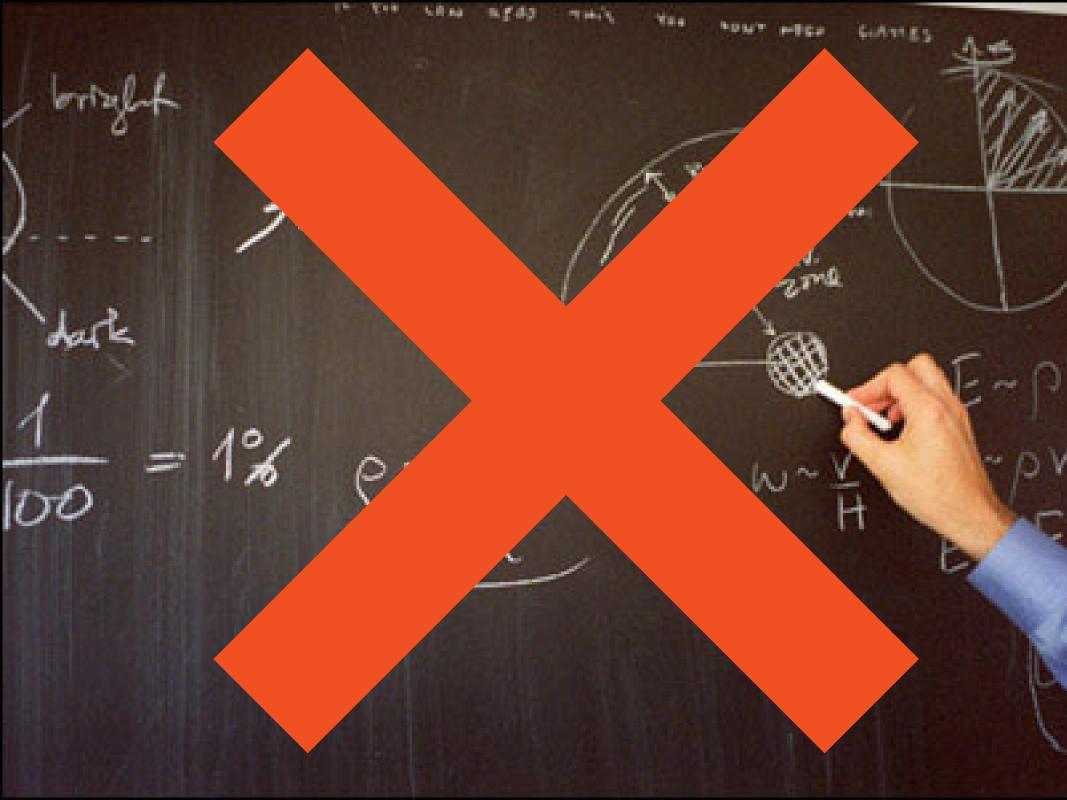


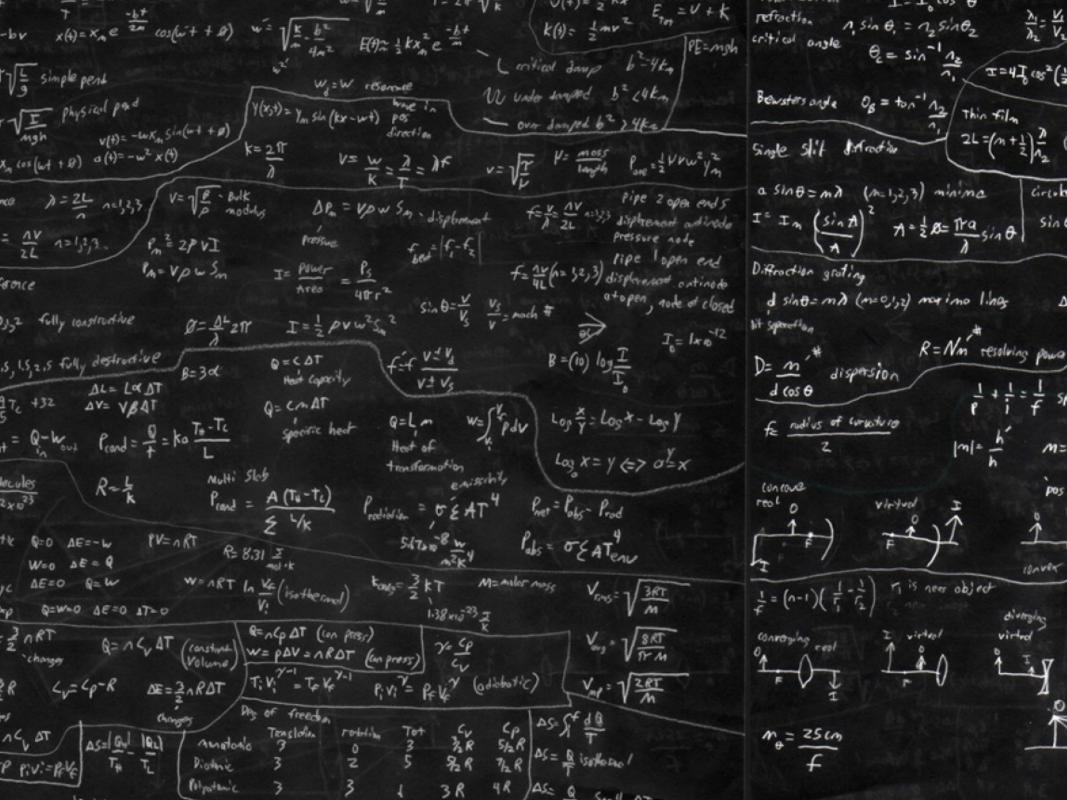


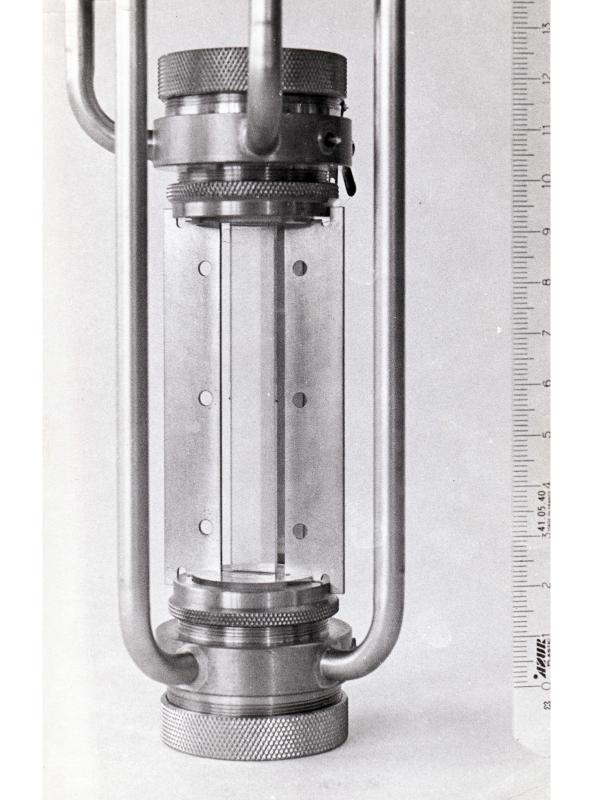


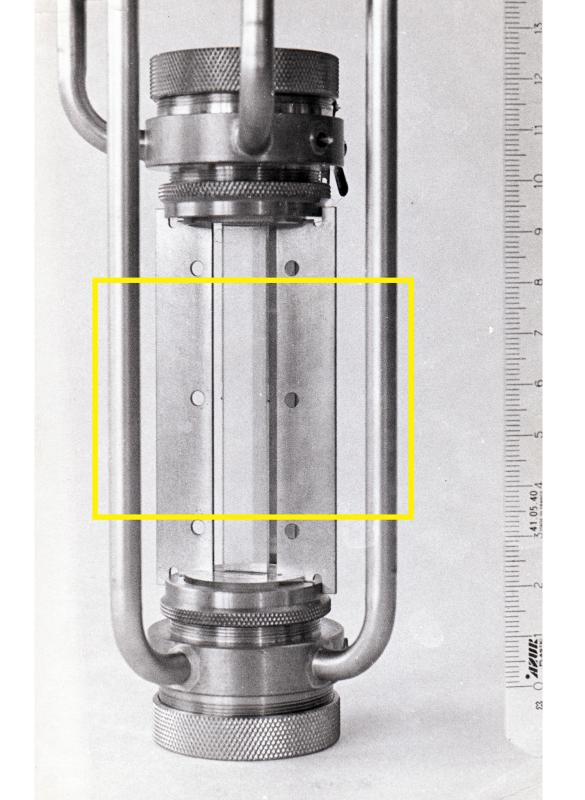


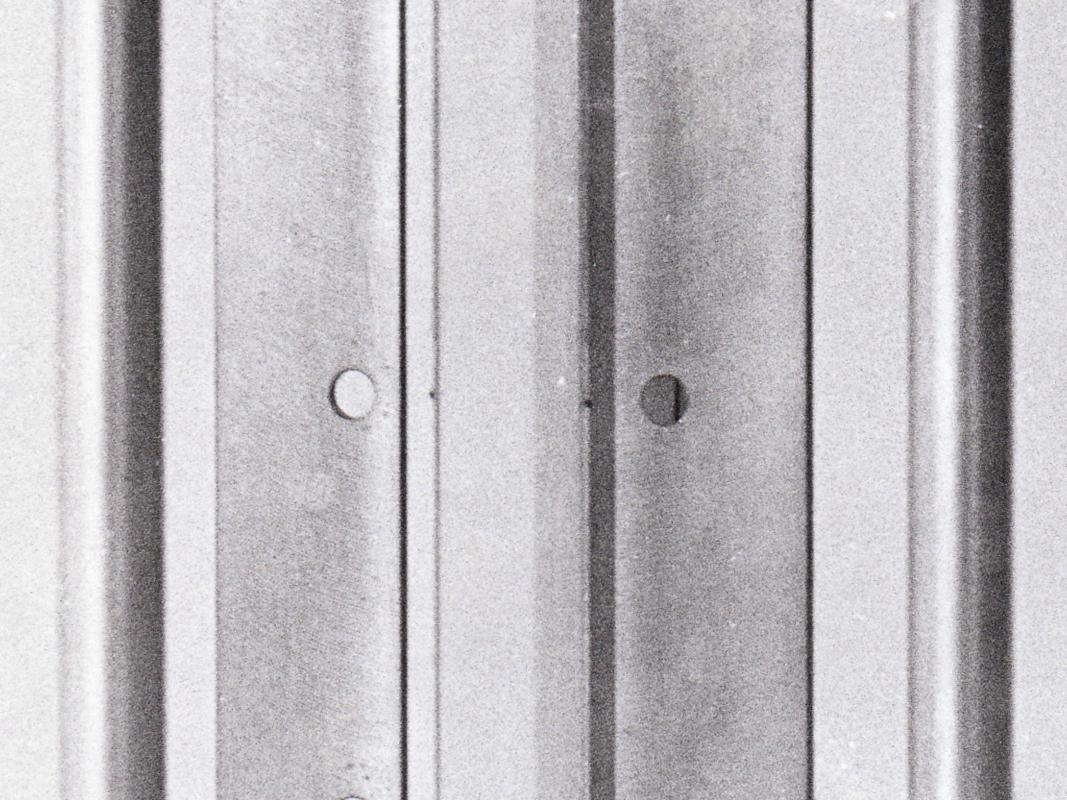


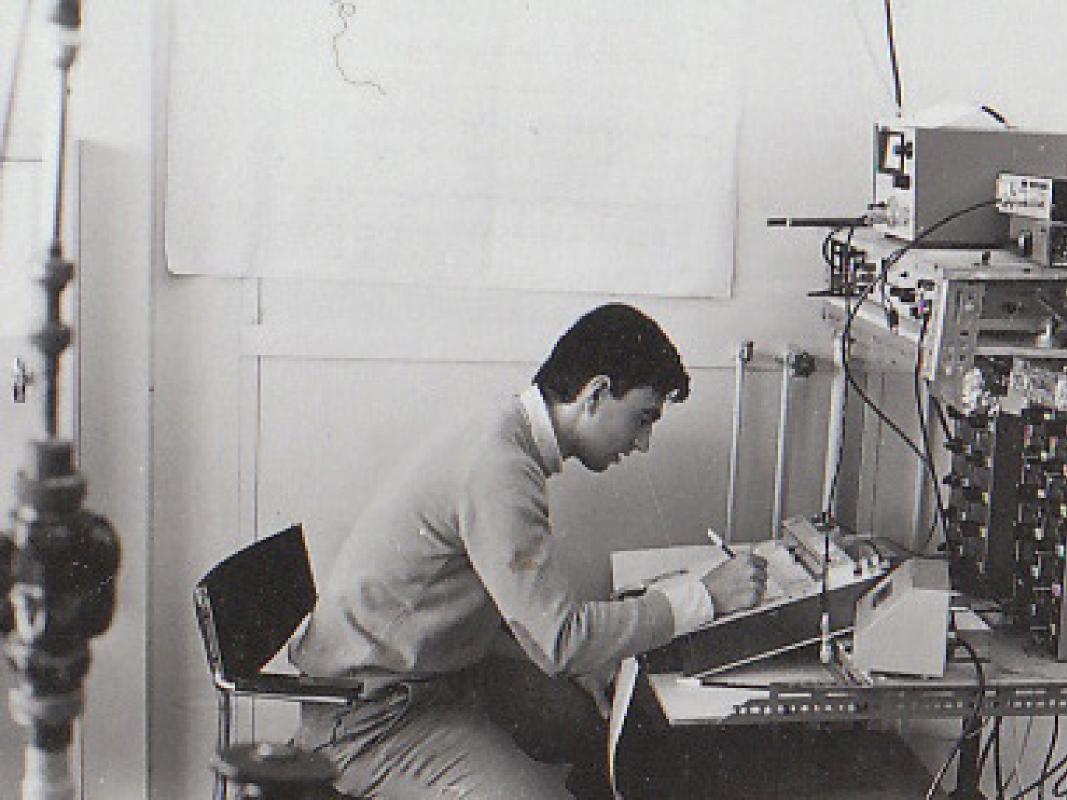


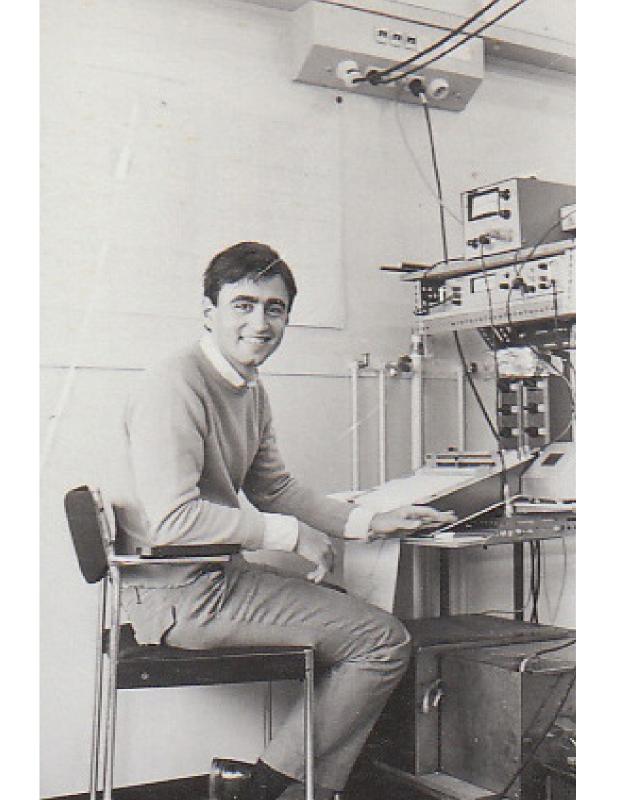










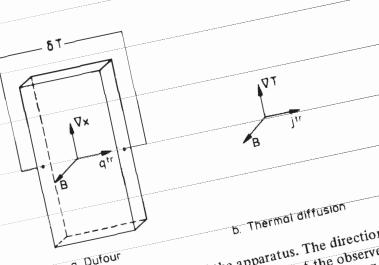


EXPERIMENTS ON THE INFLUENCE OF A MAGNETIC FIELD ON THE DUFOUR-EFFECT IN POLYATOMIC GASES: CONFIRMATION OF AN ONSAGER RELATION

E. MAZUR, G.W. 't HOOFT and L.J.F. HERMANS Huygens Laboratorium der Rijksuniversiteit, Leiden, The Netherlands

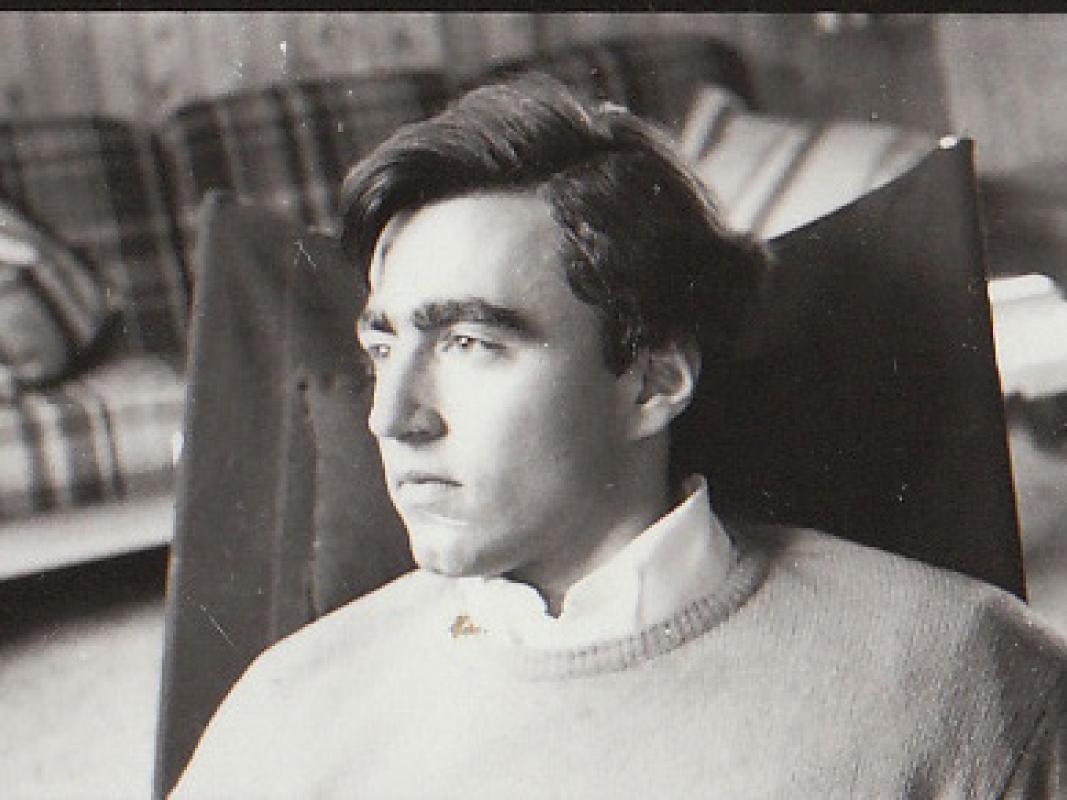
Experimental data are reported on the influence of a magnetic field on the Dufour-effect, the reciprocal phenomon of thermal diffusion, in an equimolar No-Ar mixture at room temperature. An Oneager relation in the nrecence Experimental data are reported on the influence of a magnetic field on the Dutour-effect, the reciprocal phenomenon of thermal diffusion, in an equimolar N2-Ar mixture at room temperature. An Onsager relation in the presence of a magnetic field is confirmed. of a magnetic field is confirmed.

In the absence of a magnetic field the Onsager relation between the Dufour and the thermal diffusion coefficient has been experimentally confirmed by Waldmann thirty years ago [1]. For transport phenomena occurring in polyatomic gases under the influence of an external field (Senftleben-Beenakker effects) [2], such a relation has not been verified to date. Recently [3], experiments were performed on the influence of a magnetic field on thermal diffusion: transverse thermal diffusion was measured and preliminary results for the system N2-Ar were reported. will be described on exture under the



28 November 1977

Fig. 1. (a) Schematic diagram of the apparatus. The directions of the applied N2 concentration gradient and of the observed heat flux are indicated. (b) Direction of the observed N2 flux in the thermal diffusion experiment (ref. [6]). the narrow walls of the . Tr and odd



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De heer E. Mazur 2352 BW Leiderdorp

afd. dept. abt./ref. zeichen personeelzaken TT/MR

in-dialling durchwahl doorkiesnummer

(040) 7 55850 accès intern dir.

datum, date

24 november 1981

onderw. re. conc. betr.

Van Prof. Knaap vernamen wij dat u mogelijkerwijs geïn-teresseerd hent in een bezoek aan het natuurkundig laho-Van Prof. Knaap vernamen wij dat u mogelijkerwijs gein-teresseerd bent in een bezoek aan het natuurkundig labo-In verband hiermede nodigen wij u uit op 15 januari 1981
om 0.00 uur bij de afdeling nersoneelzaken. Willemstraat Geachte heer Mazur, In verband hiermede nodigen wij u uit op 15 januari traat om 9.00 uur bij de afdeling personeelzaken, naar onder ratorium.

V.UU uur bij de ardeling personeelzaken, Willemstraat naar ondernlattegrondje bijgeslolion door ons wor-

HARVARD UNIVERSITY

DIVISION OF APPLIED SCIENCES

PIERCE HALL, CAMBRIDGE, MASSACHUSETTS 02138

12 March 1981

Mr. Eric Mazur Huygens Laboratorium Wassenaarseweg 78 2300 RA LEIDEN The Netherlands

Dear Mr. Mazur:

Referring to our correspondence in the fall of 1980, I am now pleased to inform you that we could offer you a postdoctoral position as Research Fellow, with an annual stipend of \$18,600. The appointment could start any time in the fall of 1981, preferably September 15 or October 1, 1981, for the period of one year. Usually the appointment is renewable for a second year, as a two or three-year stay is preferable for the completion of an experimental project.

The expectation is that you would participate in our research on infrared excitation of molecules with short CO₂ laser pulses, and in our work on collisional effects in four-wave light mixing with dye lasers. I am sending you some preprints of our most recent work in this area, under separate cover. Please let me know whether you are interested in the position, for which your doctor's degree from Leiden University is a prerequisite. If your answer is positive, I should like to get an indication when you will defend your thesis, and when you could start here.





Pierce Hall / 8 Aug 8/ Cambridge, MA 02138

Den Heer E. Mazur Leiden.

Waarde Hazur.

Van 5-11 September ben ik, vol. voor familieberoek, in Nederland. Graag zou ik op dofg September met je willen praten over je a.s. bezoek. Er is geen tijd om teng te sekrijven. Ik bel wel (7 of 8 Sept) voor de finitieve afspraak. Zou je ook je ouders willen inlichten over nign komst?

Met vriende gle groeten

N. Bloemberg

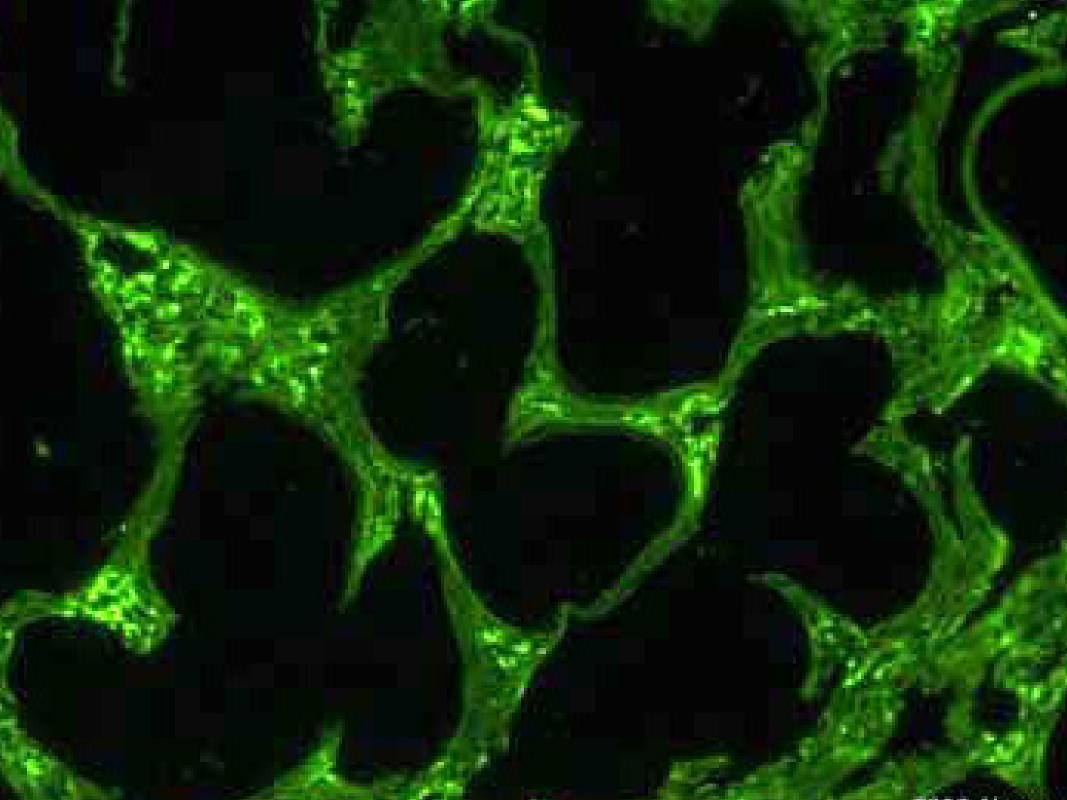


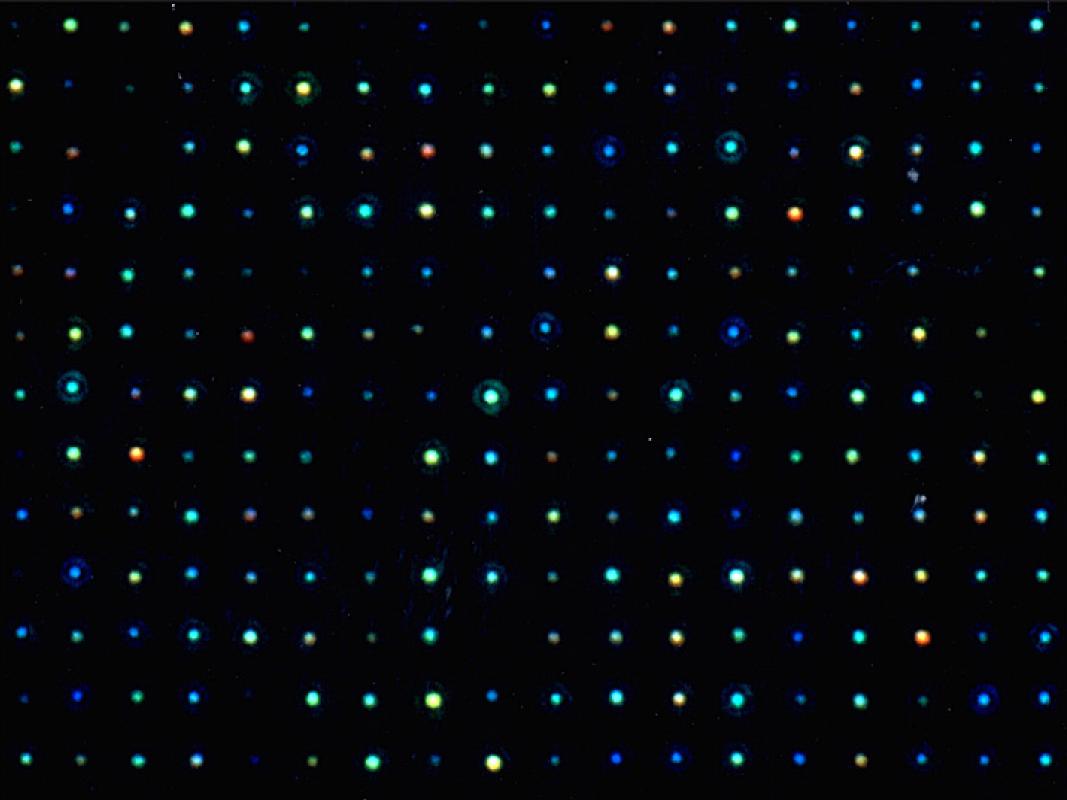


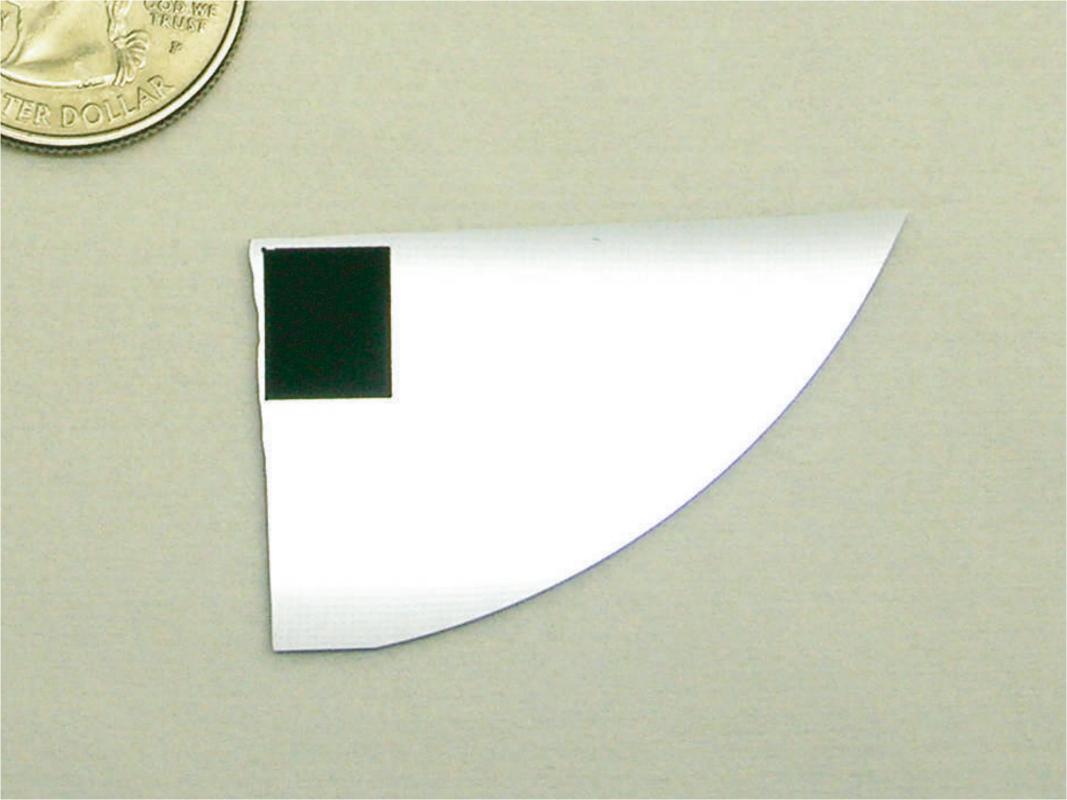




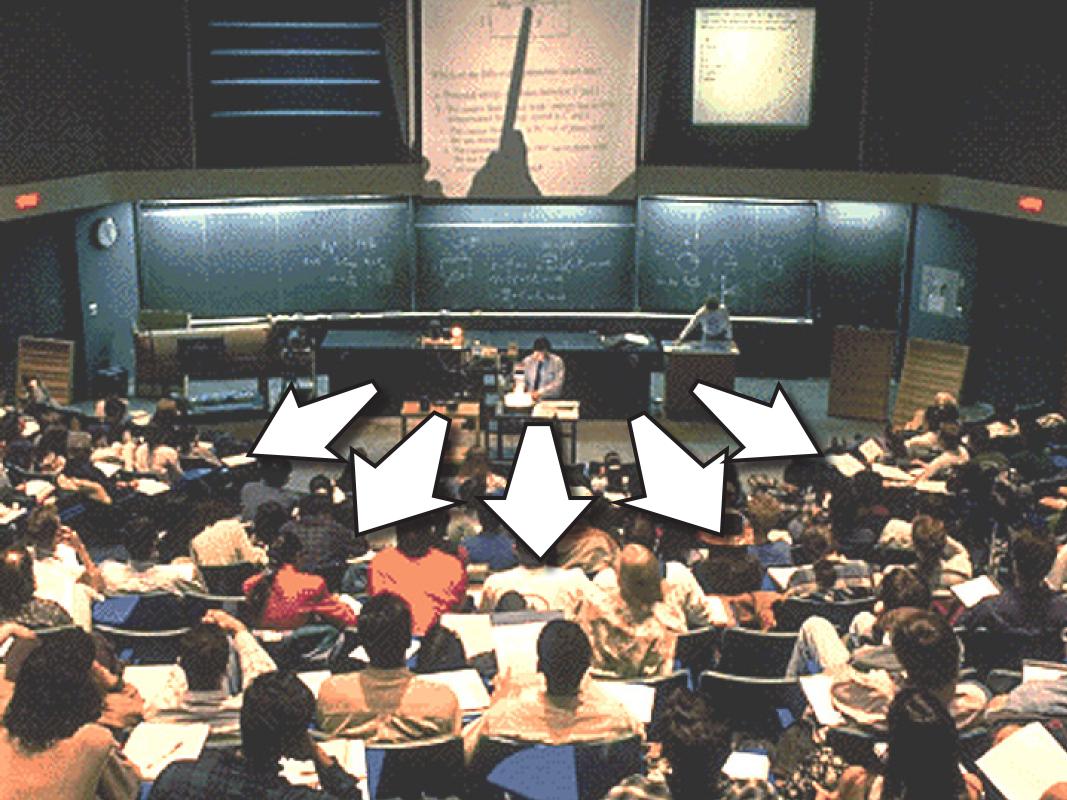
serendipity











Introducing Macintosh. For the rest of us.

In the olden days, before 1984, not very many people used computers. For a very good reason.



Not very many people knew how. And not very many people wanted

to learn.

After all, in those days, it meant listening to your stomach growl through computer seminars. Falling asleep over computer manuals. And staying awake nights to memorize commands so

to understand them.

Then, on a particularly bright day

in Cupertino, California, some particularly bright engineers had a particularly bright idea: since computers are so smart, wouldn't it make more sense

computers?

So it was that those very engineers worked long days and late nights and a few legal holidays, teaching tiny silicon chips all about people. How they make mistakes and change their minds. How they refer to file folders and save old phone numbers. How they labor for their livelihoods, and doodle in their spare time.

For the first time in recorded computer history, hardware engineers

complicated you'd have to be a computer actually talked to software engineers in moderate tones of voice, and both were united by a common goal: to build the most powerful, most portable, most flexible, most versatile computer not-verymuch-money could buy

And when the engineers were finally finished, they introduced us to to teach computers about a personal computer so personable, people, instead of teaching people about it can practically shake hands.

And so easy to use, most people already know how.

They didn't call it the QZ190, or the Zipchip 5000.

They called it Macintosh." And now we'd like to introduce it to you.



VOLUME 27, NUMBER 5

Fluctuating hydrodynamics and capillary waves

Department of Physics, University of Toronto, Toronto, Ontario, Canada M5S 1A7

We derive an equation of motion for the instantaneous position of a planar liquid surface For the fluid constrained to be in a using muctuating nyarodynamics. For the mud constrained to be in a monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient, we then obtain the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state by a small temperature gradient of the dynamic monequilibrium steady state of the dynamic monequ structure factor $S_{ss}(q,\omega)$, which represents the spectrum of thermal ripplons (capillary structure factor $S_{ss}(q,\omega)$, which represents the spectrum of the liquid curfoce. Structure 1actor $D_{88}(q,\omega)$, which represents the spectrum of thermal hypnoris (capitals).

Waves) on the liquid surface. Our analysis is done for both classical and quantum of the heights of the true similar mode in S. (2011). Waves) on the inquire surface. Our analysis is uone for bouncies and quantum finds. We find that there is an asymmetry in the heights of the two ripplon peaks in $S_{ss}(q,\omega)$ due to broken time reversal summatry. Our summatrial estimates of the effect for liquid believed to broken time reversal summatry. to broken time-reversal symmetry. Our numerical estimates of the effect for liquid helium. to broken time-reversal symmetry. Our numerical estimates of the check for inquire nemanity of $S_{ss}(q,\omega)$; II, suggest that it could be measurable. We have undertaken a moment analysis of $S_{ss}(q,\omega)$; 11, suggest that it could be measurable. We have undertaken a moment analysis of $\mathfrak{I}_{ss}(q,\omega)$; the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, though it vanishes in the first frequency moment is long ranged $(\sim 1/q^2)$ in the steady state, the first frequency moment is $(\sim 1/q^2)$ in the steady state, the first frequency moment is $(\sim 1/q^2)$ in the steady state, the first frequency moment is $(\sim 1/q^2)$ in the steady state, the first frequency moment is $(\sim 1/q^2)$ in the steady state, the first frequency moment is $(\sim 1/q^2)$. equilibrium. From the zeroth frequency moment we recover the familiar result for the instance of the moment we recover the familiar result for the moment with the moment we recover the familiar result for the moment with the moment we recover the familiar result for the moment with the moment we recover the familiar result for the moment with the moment we recover the moment with the m frared logarithmic divergence of the mean-square displacement of the interface; the analysis explicitly involves the coupling of the surface to the bulk fluid. We modify the so-called capillary-wave model by including the surface-bulk coupling, and thereby make it consistent with the capillary-wave dispersion.

I. INTRODUCTION

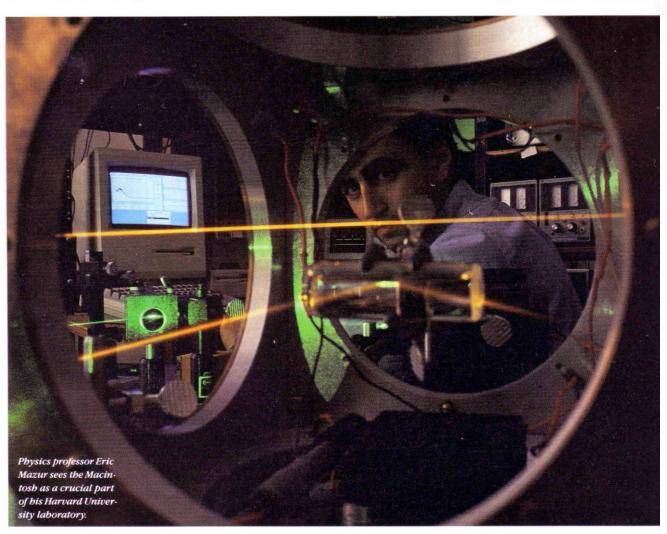
In this paper we present new results for the frequency spectrum of excitations on a liquid surface. We derive an equation of motion for the instantaneous position of the interface from linear fluctuating hydrodynamics. This equation is used to analyze fluctuations about equilibrium, and fluctuations steady state characterized by The analysis is made

In the following we present related contributions to the problems outlined in the preceding two paragraphs. The equilibrium, and nonequilibrium steady state, spectrum of capillary waves is studied through fluctuating hydrodynamics. The format of the pa-In Sec. II we use the Landau-Lifshitz equations of

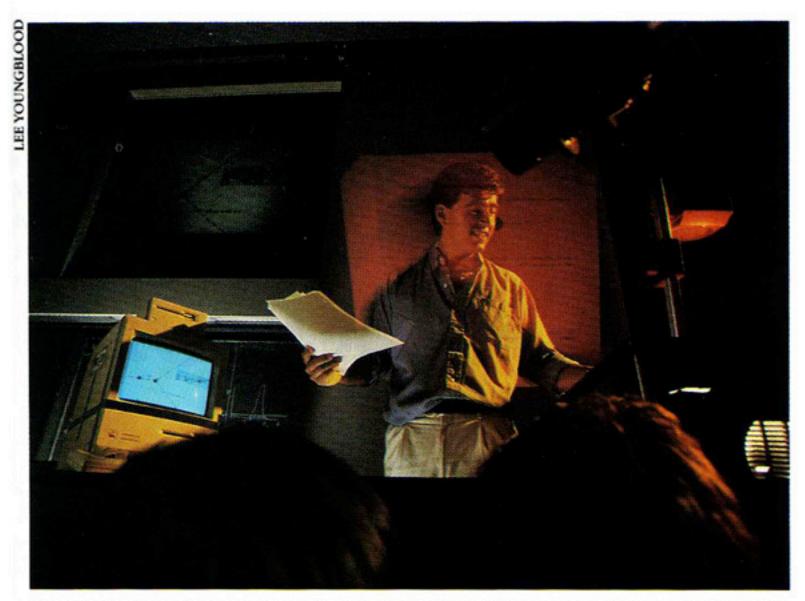
linear fluctuating hydrodynamics to obtain an equaper is as follows. tion of motion for the instantaneous position of a planar liquid-vapor interface 5, valid for both classical and quantum systems. We find that the thermal fluctuations which drive the surface ripples are due the explicit counling between the surface and of motion for 5, Eq. (2.14) heen given in the

The Labtop Macintosh

Dan McNeill and Paul Freiberger



Macworld News



Many instructors are using the Macintosh as a teaching assistant. Harvard physics professor Eric Mazur animates the collision of two protons with VideoWorks.

DOING PHYSICS WITH COMPUTERS

Eric Mazur Harvard University

re the days of watching analog me-Introduction ters, taking notes in thick lab books, and plotting data points on graph paper gone forever? Is research in the physical sciences becoming so complex that one can no longer do research without computers? A superficial survey of the current research in physics might lead one to give an affirmative answer to these questions. It is therefore interesting to note that the award of this year's Nobel Prize in physics to Alex Müller and Georg Bednorz for their work on superconductivity was hailed as a victory for relatively simple, small-scale research. No computer was needed to show that their compound of copper, oxygen, lanthanum, and barium becomes superconducting. It shows that one can still make major breakthroughs with very simple means without computers. Computers excel at performing tedious routine tasks. Physics research, on the other hand, seldom entails routine work. This holds true, in particular in theoretical physics: indeed, no . Jony that the

mediate group, needless to say, is roughly equally distributed between those who do and those who do not use computers. This simple survey shows that, not surprisingly, the distribution of computers among physicists seems to be strongly related to age. Therefore, one might expect that in just a few more years computers will be used in one way or another by the entire physics community.

The use of computers in physics is certainly not restricted to the "computer Research revolution" of the last decade. Large computers have been used for numerical

the research in about 95% of the experimental papers involves computers. For theoretical papers this percentage is 40%. The use of computers is more common in experimental than in theoretical groups, in part because more affordable, small personal computers are usually sufficient for the needs of an experimentalist. A more detailed study shows that the

use of computers in physics can be divided into several, fairly different categories. An overview of these categories is presented in Table 1. The first category, laboratory equipment control, is basically a hardware application of computers to physics. The next two categories, numeri-

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AN INTERACTIVE STUDY COMPANION by ERIC MAZUR Harvard University

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