Less is more: Extreme optics with zero refractive index





Pontificia Universidad Católica de Chile Santiago, Chile, 12 January 2016

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governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$



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governed by wave equation

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where

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$



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and
$$n=\sqrt{\epsilon\mu}$$
 .



governed by wave equation

$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

 $\frac{1}{-c}$

n

Solution:
$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)}$$

 (\mathbf{i})

where
$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c =$$

and
$$n = \sqrt{\epsilon \mu}$$
 .

In dispersive media $n = n(\omega)$.



Lorentz oscillator

































for a strong (dielectric) resonance ε can become negative





valence electrons in dielectric then behave like a plasma





with plasma frequency above the resonance





(and far below the UV region)





Index also determined by magnetic response

$$n = \sqrt{\boldsymbol{\omega}}$$



Index also determined by magnetic response

$$n = \sqrt{\boldsymbol{\omega}}$$

and magnetic response shows similar resonances





Magnetic response





but magnetic resonances occur below optical frequencies





Magnetic response

so, in optical regime, $\mu \approx 1$





Index of refraction

$$n = \sqrt{\epsilon \mu}$$

Both ϵ and μ are complex and their real parts can be negative.



Index of refraction

$$n = \sqrt{\epsilon \mu}$$

Both ϵ and μ are complex and their real parts can be negative.

What happens when $\operatorname{Re}\epsilon$ and/or $\operatorname{Re}\mu$ is negative?



$$\varepsilon = |\varepsilon| e^{i\theta} \qquad \mu = |\mu| e^{i\phi}$$



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$$\varepsilon = |\varepsilon| e^{i\theta} \qquad \mu = |\mu| e^{i\phi}$$

Index

$$n = \sqrt{|\varepsilon| |\mu|} e^{i\frac{\theta + \phi}{2}}$$





$$\varepsilon = |\varepsilon| e^{i\theta} \qquad \mu = |\mu| e^{i\phi}$$

Index

$$n = \sqrt{|\varepsilon| |\mu|} e^{i\frac{\theta + \phi}{2}}$$




Q: Is this the only possible solution?

- **1. yes**
- 2. no, there's one more
- 3. there are many more
- 4. it depends









Can add 2π to exponent

$$e^{+i(\theta+\phi)} = e^{+i[\theta+\phi+2\pi]}$$





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and so

$$n = \sqrt{|\varepsilon||\mu|} e^{i\left[\frac{\theta+\phi}{2}+\pi\right]}$$





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and so

$$n = \sqrt{|\mathcal{E}||\mu|} e^{i\left[\frac{\theta+\phi}{2}+\pi\right]}$$

but...

$$k = \frac{2\pi n}{\lambda_o} = \frac{2\pi (n' + in'')}{\lambda_o} = k' + ik''$$





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$$E = E_o e^{i(kx - \omega t)} = E_o e^{-k''x} e^{i(k'x - \omega t)}$$





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Im(*n*) n ϵ μ $\operatorname{Re}(n)$ n must lie here for passive material

$$E = E_o e^{i(kx - \omega t)} = E_o e^{-k''x} e^{i(k'x - \omega t)}$$



Can add $2\pi\, {\rm to} \ {\rm exponent}$

$$e^{+i(\theta+\phi)} = e^{+i[\theta+\phi+2\pi]}$$

and so

$$n = \sqrt{\left| \mathcal{E} \right| \left| \mu \right|} e^{i \left[\frac{\theta + \phi}{2} + \pi \right]}$$

but...

$$k = \frac{2\pi n}{\lambda_o} = \frac{2\pi (n' + in'')}{\lambda_o} = k' + ik''$$

Im(*n*) E μ $\operatorname{Re}(n)$ must lie here for passive material

$$E = E_o e^{i(kx - \omega t)} = E_o e^{-k''x} e^{i(k'x - \omega t)}$$



Q: Is this the only possible solution?







To find *n* (passive materials):

- 1. Draw line that bisects ϵ and μ
- 2. Choose upper branch







For certain values of ϵ and μ we can get a *negative* $\operatorname{Re}(n)!$



Q: Must both $\operatorname{Re}\epsilon < 0$ and $\operatorname{Re}\mu < 0$

to get a negative index?

1. yes

2. no







1. yes

2. no 🖌





Note: need magnetic response

to achieve $n \le 0$!





Now remember

$$k = \frac{2\pi n}{\lambda_o} = \frac{2\pi (n' + in'')}{\lambda_o} = k' + ik''$$



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Spatial and temporal dependence of wave component

$$E = E_o e^{i(kx - \omega t)} = E_o e^{-k''x} e^{i(k'x - \omega t)}$$



Now remember

$$k = \frac{2\pi n}{\lambda_o} = \frac{2\pi (n' + in'')}{\lambda_o} = k' + ik''$$

Spatial and temporal dependence of wave component

$$E = E_o e^{i(kx - \omega t)} = E_o e^{-k''x} e^{i(k'x - \omega t)}$$



































































What happens on the axes?




what if we let $\varepsilon = 0$?







what if we let $\varepsilon = 0$?







Q: If n = 0, which of the following is true?

- 1. the frequency goes to zero.
- 2. the phase velocity becomes infinite.
- 3. both of the above.
- 4. neither of the above.





$$\nabla^2 \vec{E} - \frac{\mu \epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

solution

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solution

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)} \longrightarrow \vec{E} = \vec{E}_o e^{-i\omega t}$$

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c$$







solution

$$\vec{E} = \vec{E}_o e^{i(kx - \omega t)} \longrightarrow \vec{E} = \vec{E}_o e^{-i\omega t}$$

$$\frac{\omega}{k} = \frac{1}{\sqrt{\epsilon\mu}} c = \frac{1}{n} c \longrightarrow \infty$$





Q: If n = 0, which of the following is true?

1. the frequency goes to zero.

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- 3. both of the above.
- 4. neither of the above.





















































"tunneling with infinite decay length"







$$n = \sqrt{\varepsilon \mu}$$





$$n = \sqrt{\epsilon \mu}$$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$





$$n = \sqrt{\epsilon \mu}$$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}}$$





$$\varepsilon \to 0$$
 $n = \sqrt{\varepsilon \mu} \to 0$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}}$$





$$\varepsilon \to 0$$
 $n = \sqrt{\varepsilon \mu} \to 0$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}} \quad \longrightarrow \infty$$





$$\varepsilon \to 0$$
 $n = \sqrt{\varepsilon \mu} \to 0$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1} \longrightarrow 1$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}} \quad \longrightarrow \infty$$





$$\mu \to 0 \qquad \qquad n = \sqrt{\varepsilon \mu} \to 0$$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}}$$





$$\mu \to 0 \qquad \qquad n = \sqrt{\varepsilon \mu} \to 0$$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}} \longrightarrow 0$$





$$\mu \to 0 \qquad \qquad n = \sqrt{\varepsilon \mu} \to 0$$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1} \rightarrow -1$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}} \longrightarrow 0$$





$$\varepsilon, \mu \to 0$$
 $n = \sqrt{\varepsilon \mu} \to 0$

but ϵ and μ also determine reflectivity

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

$$Z = \sqrt{\frac{\mu}{\varepsilon}} \qquad \text{finite!}$$





but $\mu \neq 1$ requires a magnetic response!













use array of dielectric rods









incident electromagnetic wave ($\lambda_{eff} \approx a$)







produces an electric response...









... but different electric fields front and back...









...induce different polarizations on opposite sides...









...causing a current loop...









...which, in turn, produces an induced magnetic field








adjust design so electrical and magnetic resonances coincide

























































How to fabricate?



































































































































































































SU8 slab waveguide



Si waveguide







SU8 slab waveguide

prism



Si waveguide















at design wavelength (1590 nm)








below design wavelength (1530 nm)









above design wavelength (1650 nm)









On-chip zero-index prism









































































Wavelength dependence of index









Wavelength dependence of index









More extreme optics

- suppressing losses
- beam steering & supercoupling
- nonlinear optics
- quantum optics









on-chip zero-index material

uniform field inside material (infinite wavelength)

many exciting applications ahead!

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Zero-index metamaterials

PHASE-CHANGE MATERIALS Multi-level memory

MID-INFRARED SOURCES Powerful pulse train

OPTICAL COMPUTING Analog approach

More info: download paper!

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for a copy of this presentation:

http://ericmazur.com

