Laser doping and texturing of silicon for advanced optoelectronic devices





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Laser doping and texturing of silicon for advanced optoelectronic devices







Renee Sher



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and also....

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> Prof. Alan Aspuru-Guzik Prof. Michael Aziz Prof. Michael Brenner Prof. Cynthia Friend Prof. Howard Stone

Dr. Martin Pralle (SiOnyx) and everyone else at SiOnyx...

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Dr. Richard Farrell (RMD) Dr. Arieh Karger (RMD) Dr. Richard Meyers (RMD)

Dr. Pat Maloney (NVSED)

Dr. Jeffrey Warrander (ARDEC)



irradiate with 100-fs 10 kJ/m² pulses







absorptance
$$(1 - R_{int} - T_{int})$$



absorptance
$$(1 - R_{int} - T_{int})$$





absorptance
$$(1 - R_{int} - T_{int})$$



absorptance
$$(1 - R_{int} - T_{int})$$



absorptance
$$(1 - R_{int} - T_{int})$$



laser treatment causes:

- surface structuring
- inclusion of dopants









































cross-sectional Transmission Electron Microscopy





M. Wall, F. Génin (LLNL)



1 µm











- 300-nm disordered surface layer
- undisturbed crystalline core

• surface layer: polycrystalline Si with 1.6% sulfur





two processes: melting and ablation




















Nature Materials 1, 217 (2002)





Nature Materials 1, 217 (2002)





Nature Materials 1, 217 (2002)



different thresholds:

melting: 1.5 kJ/m²

ablation: 3.1 kJ/m²































ероху		
laser affected region		
substrate		
100 nm		











secondary ion mass spectrometry





Things to keep in mind

- near unit absorption extending into IR
- surface structure due to ablation
- hyperdoping due to rapid melting and resolidification
- can decouple both processes





2 intermediate band

properties

gap determines optical and electronic properties



2 intermediate band

properties

shallow-level dopants control electronic properties



2 intermediate band

properties

deep-level dopants typically avoided



2 intermediate band

properties

1 part in 10⁶ sulfur introduces donor states in gap



Janzén et al., Phys. Rev. B 29, 1907 (1984)

properties

1

2 intermediate band

1 part in 10⁶ sulfur introduces donor states in gap



Janzén et al., Phys. Rev. B 29, 1907 (1984)

properties

1

2 intermediate band

at high concentration states broaden into band





properties

10⁻⁶ sulfur doping



2 intermediate band

properties

laser-doped S:Si



2 intermediate band

properties

laser-doped S:Si



2 intermediate band

properties

laser-doped S:Si



2 intermediate band

properties

device layer

buried oxide

silicon substrate



properties





properties

laser doped region buried oxide

silicon substrate



properties





properties





properties



Hall measurements







Hall measurements














intermediate band

properties





1



2 intermediate band

properties





1



properties





properties

1



impurity (donor) band centered at 310 meV





properties

PRL 108, 026401 (2012)



electronic properties, although such properties currently lack a clear microscopic explanation. Combining computational and experimental evidence, we probe the origin of sub-band-gap optical absorption and metallicity in Se-hyperdoped Si. We show that sub-band-gap absorption arises from direct defect computational and experimental evidence, we probe the origin of sub-band-gap optical absorption and metallicity in Se-hyperdoped Si. We show that sub-band-gap absorption arises from direct defect-to-conduction-band transitions rather than free carrier absorption. Density functional theory predicts the and metallicity in Se-hyperdoped Si. We show that sub-band-gap absorption arises from direct detect-to-conduction-band transitions rather than free carrier absorption. Density functional theory predicts and Se-induced insulator-to-metal transition arises from merging of defect and conduction bands. to-conduction-band transitions rather than free carrier absorption. Density functional theory predicts the order and conduction bands, at a second insulator-to-metal transition arises from merging of defect and calculations confirm the experiment. Ouantum Monte Carlo calculations confirm the experiment. Se-induced insulator-to-metal transition arises from merging of detect and conduction bands, at a concentration in excellent agreement with experiment. Quantum Monte Carlo calculations accurately, and critical concentration, demonstrate that correlation is important to describing the transition accurately. concentration in excellent agreement with experiment. Quantum Monte Carlo calculations confirm the critical concentration, demonstrate that correlation is important to describing the transition accurately, and success that it is a classic immurity-driven Mott transition. suggest that it is a classic impurity-driven Mott transition. DOI: 10.1103/PhysRevLett.108.026401 Of all the experimentally measurable physical properties of materials, electronic conductivity exhibits the largest variation, spanning a factor of 10³¹ from the best metals to variation, spanning a racion of the last century, the strongest insulators [1]. Over the last century, the puzzle of why some materials are conductors and others Pucco or why some machans are conductors and outputs insulators, and the mechanisms underlying the transformation is to the other base base base for the state of tion from one to the other, have been carefully scrutinized; yet even after such a vast body of research over such a long period, the subject remains the object of controversy. In 1056 Note introduced a model for the involvent to motel 1956, Mott introduced a model for the insulator-to-metal transition (IMT) in doped semiconductors, in which long ranged electron correlations are the driving force [2]. Hyperdoping (doping beyond the solubility limit) creates a new materials playground to explore defect-mediated a new matchais play ground to copiore acrect methated iMTs in semiconductors. In this Letter, we identify a defect-induced IMT in silicon hyperdoped with selenium to concentrations exceeding 10²⁰ cm⁻³ (compared to the subbility limit f21 of about 1016 cm⁻³) contained to the subbility limit f21 of about 1016 cm⁻³) cm⁻³ wincentrations exceeding 10° cm (compared to me ilibrium solubility limit [3] of about 10¹⁶ cm⁻³) and we deteiled extreme of the transition of the lotailed nature of the transition with both orion. We find that the IMT is an and most resembles utionally, we

Insulator-to-Metal Transition in Selenium-Hyperdoped Silicon: Observation and Origin

Elif Ertekin,^{1,*} Mark T. Winkler,^{2,†} Daniel Recht,³ Aurore J. Said,³ Michael J. Aziz,³

Tonio Buonassisi,² and Jeffrey C. Grossman^{1,24+} Tonio Buonassisi,² and Jeffrey C. Grossman^{1,24+} Massachusetts Institute of Technology, Cambridge Massachusetts 02139, USA ¹Department of Mechanical Engineering, Massachusetts Institute of Technology, Cambridge Massachusetts 02139, USA partment of Materials Science and Engineering, Massachusetts Institute of Technology, Cambridge Massachusetts 02139, USA ²Department of Mechanical Engineering, Massachusetts Institute of Technology Massachusetts 02138, USA

Hyperdoping has emerged as a promising method for designing semiconductors with unique optical and lectronic properties, although such properties currently lack a clear microscopic explanation. Combining Hyperdoping nas emerged as a promising method for designing semiconductors with unique optical and electronic properties, although such properties currently lack a clear microscopic explanation. Combining computational and experimental evidence, we probe the origin of sub-band-gan ontical absorption

> implantation followed by nanosecond pulsed-laser melting (PLM) and rapid resolidification. The PLM process enasuch as chalcogens [17]. (LLM) and rapid resolution and rule that process one bles chalcogen doping with concentrations exceeding 1% atomic; such samples exhibit unexplained optical properawme, such samples commun uncorplamen optical photons ties including broad, featureless absorption of photons ues including vival, realurcless ausviption or provide with energy lower than the band gap of silicon [9]. with energy lower than the band gap of silicon [9]. Silicon substrates (boron doped, $\rho \approx 25 \ \Omega \text{ cm}$) were ion implanted with Se to nominal doces of $3 \times 10^{15} \text{ cm}$. ion implanted with Se to nominal doses of 3×10^{15} and 1×10^{16} cm² using an ion beam energy of 176 keV The Not implained with SC to noninnal access of 3×10^{-10} . The 1×10^{16} cm² using an ion beam energy of 176 keV. (a.t. implanted samples were exposed to four laser pulses (fluimplanted samples were exposed to lour laser pulses (ilu-ences of 1.7, 1.7, 1.7 and 1.8 J cm⁻²). This fluence regiment increases of the state of the second second birds results in a slightly shallower dopant profile, and higher Les III a sugary sharrower workant province, and ingine Se concentration, than reported previously [18]. The crystalline, extends approximately 350 nm electrically isolated from the ction formed between s measured

PACS numbers: 71.30.+h, 61.72.sd, 73.61.Cw, 78.20.Bh silicon appears to justify such interest. While isolated S and Se dopants are well-established deep double donors in and se to 1/11 the entergoid entired entergies of home and be uppande are well-established user upper upper institution [3,14], the enhanced optical properties of hyper-silicon [3,14], the enhanced optical properties of hyper-Survey Logith we can ance optical properties or any per-doped silicon (in which these chalcogenic impurities are recent at much bisher concentrations) are not not and urbor announ un winch new charoserne imprimes au present at much higher concentrations) are not yet well and restord. Earther unlike the restation of anter of present at much mener concentrations) are not yet went understood. Further, unlike the prototypical system of phosphorus-doped silicon for which the INT has been extensively studied and characterized [15,16], there are very few studies of an IMT resulting from deep defects We prepared Se-doped silicon (Se:Si) samples using ion

13 JANUARY 2012





intermediate band

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ences of 1.(, 1.(, 1.) and 1.0 year (). This increase regiment results in a slightly shallower dopant profile, and higher S in a sugnery shanower uppan prome, and menor Se concentration, than reported previously [18]. The ervstalline, extends approximately 350 nm electrically isolated from the tion formed between



Double side polished p type (001) Si wafers, resistivity by under since poinsing p type (UU1) or waters, residurity of 5-25 Ω cm, were ion implanted at room temperature with of 5-25 Ω cm, $^{32}S^-$, 176 keV $^{80}Se^+$, or 245 keV $^{130}Te^+$ to either 95 keV $_{10}^{16}$ iconform² The doce of $^{32}e^-$ was varied [http://dx.doi.org/10.1063/1.4804935] either 95 KeV -5, 1/0 KeV -5e, or 245 KeV 1e to doses of 1×10^{16} ions/cm². The dose of $^{32}S^{-1}$ was varied doses of 1×10^{-100} ions/cm² and pre-amorphized by from 3×10^{14} to 1×10^{16} ions/cm² and pre-amorphized by So key S1 to doses of 5×10^{-10} tons/cm when the S dose is not greater than 1×10^{15} tons/cm². Pulsed laser melt-Silicon hyperdoped with chalcogens can be synthesized ing was performed using a XeCl exciment laser beam by pulsed laser irradiation in a sulfur bearing atmosphere, ung was producting using a new environmen laser usample (308 nm, 25 ns FWHM, 50 ns total duration). Each sample I. INTRODUCTION ion implantation followed by pulsed laser melting, ion implantation followed by pulsed laser melting, or or pulsed laser mixing. This material has attracted interest because of its sub band can abcomion and has been existed 3,4 or (JUO HILL, 2J HS F W FIVE, JU IIS WHAT AUTAHON). Each sample received three laser shots at 1.7 J/cm² followed by a fourth laser shot at 1.8 H/cm². Time received reporting of because of its sub band gap absorption and has been studied of a constitute for information (D) shots the state of the sublaser shot at 1.8 J/cm². Time-resolved reflectivity of a as a candidate for infrared (IR) photodetectors 488 nm Ar ion laser was used to measure the melt duration. The laser fluence was calibrated by comparing the melt dura-9-11 In addition, observations of carrier lifetime recovery for sufficiently high concentrations of 12,13 the laser nuclear was canonaled by comparing the men same tion with numerical solutions to the one-dimensional heat has aroused similar interest in this material. erdoping has been shown to cause an $\frac{1}{4}$ has been proposed to form an $\frac{1}{4,10,14-16}$ However,

the formation of an impurity band from isolated deep donor levels as the concentration LLC. of chalcogen atoms in metastable local configurations increases. © 2013 AIP Publishing LLC. [http://dx.doi.org/10.1063/1.4804935]

⁵Department of Physics, American University of Beirut, Beirut 1107 2020, Lebanon ⁶Institute of Electronic Structure and Laser, Foundation for Research and Technology Hellas, P.O. Box 1527, 71110 Heraklion, Greece *Department of Applied Physics, Fukuoka University, Fukuoka 814-0180, Japan ⁵Department of Physics, American University of Beirut, Beirut 1107 2020, Lebanon ⁶Institute of Electronic Structure and Laser Foundation for Research and Technolog 71110 Heraklion, Greece 7 Materials Science and Technology Department, University of Crete, P.O. Box 2208, 71003 Heraklion, Greece 7 Materials Science and Technology Department, University of Crete, P.O. Box 2208, 71003 Heraklion, Greece 8 Department of Physics, Applied Physics and Astronomy, Rensselaer Polytechnic Institute, Troy, ¹Materials Science and Technology Department, University of Crete, P.O. Box 2208, 71003 Herak ⁸Department of Physics, Applied Physics and Astronomy, Rensselaer Polytechnic Institute, Troy, New York 12180. USA New York 12180, USA 9 Harvard School of Engineering and Applied Sciences, Cambridge, Massachusetts 02138, USA (Received 9 September 2012; accepted 29 April 2013; published online 3 June 2013) We report the near through mid-infrared (MIR) optical absorption spectra, over the range 0.05-1.3 eV. of monocrystalline silicon layers hyperdoped with chalcogen atoms synthesized by We report the near through mid-infrared (MIR) optical absorption spectra, over the range 0.05–1.3 eV, of monocrystalline silicon layers hyperdoped with chalcogen atoms synthesized based ion implantation followed by pulsed laser melting. A broad mid-infrared optical absorption based at the second 0.05–1.3 eV, of monocrystalline silicon layers hyperdoped with chalcogen atoms synthesized by ion implantation followed by pulsed laser melting. A broad mid-infrared optical absorption band emerges, peaking near 0.5 eV for sulfur and selenium and 0.3 eV for tellurium hyperdoped samples. ion implantation tollowed by pulsed laser melting. A broad mid-infrared optical absorption band emerges, peaking near 0.5 eV for sulfur and selenium and 0.3 eV for tellurium hyperdoped samples. Its strength and width increase with impurity concentration. Its strength decreases markedly with emerges, peaking near U.SeV for sulfur and selenium and U.SeV for tellurium hyperdoped samples, Its strength and width increase with impurity concentration. Its strength decreases markedly with subsequent thermal annealing. The emergence of a broad MIR absorption hand is consistent with Its strength and width increase with impurity concentration. Its strength decreases markedly with subsequent thermal annealing. The emergence of a broad MIR absorption band is consistent with the formation of an impurity band from isolated deep donor levels as the concentration the formation of an impurity band from isolated deep donor levels as the concentration. subsequent thermal annealing. The emergence of a broad MIR absorption band is consistent with the formation of an impurity band from isolated deep donor levels as the concentration of chalcogen atoms in metastable local configurations increases. © 2013 AIP Publishing LLC.

Emergence of very broad infrared absorption band by hyperdoping of silicon with chalconene Ikurou Umezu,¹ Jeffrey M. Warrender,² Supakit Charnvanichborikarn,³ Atsushi Kohno,⁴ 3.8.9 James S. Williams,³ Malek Tabbal.⁵ Dimitris G. Papazodiou. Ikurou Umezu,¹ Jeffrey M. Warrender,² Supakit Charnvanichborikarn,³ Atsushi K₈, James S. Williams,³ Malek Tabbal,⁵ Dimitris G. Papazoglou,^{6,7} Xi-Cheng Zhang, and Michael J. Aziz⁹ ²U.S. Army ARDEC-Benét Laboratories, Watervliet, New York 12189, USA ³Research School of Physics and Engineering, The Australian National University, Canberra, ACT 0200, ³Australia with chalcogens and Michael J. AZIZ ¹Department of Physics, Konan University, Kobe 658-8501, Japan ¹Department of Physics, Konan University, Watervillet, New York 1218 ²U.S. Army, ARDEC, Renét Laboratories, Watervillet, New York 1218 [•]Department of Physics, Konan University, Kobe 658-8501, Japan [•]Department of Physics, Konan University, Kobe 658-8501, Japan [•]2 U.S. Army ARDEC_Benét Laboratories, Watervliet, New York 12189, USA [•]2 U.S. Army ARDEC_Benét Laboratories, The Australian National Unive [•]3 Research School of Physics and Engineering. The Australian National Unive Australia ⁴Department of Applied Physics, Fukuoka University, Fukuoka 814-0180, Japan ⁵Department of Physics American University of Reirut Reirut 1107 2020 Lebanc 5Department of Physics and Michael J. Aziz

18 The details of the sample preparation method

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intermediate band

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Things to keep in mind

properties

- IR absorption rolls off around 8 µm
- consistent evidence of intermediate band formation
- IB forms at 0.1% at. doping, broadens at higher doping
- IB merges with CB at 0.4% at. yielding metallic behavior









should have shallow junction below surface











excellent rectification (after annealing)



1 properties







1 properties







1 properties







properties

1







properties

1







properties





enhanced sensitivity

• extended IR response









www.sionyx.com

1 properties



intermediate band



SiOnyx

Combine state-of-the-art low-noise CMOS image

sensor design with enhanced quantum efficiency

US Patents: US 8,058,615; US 7,928,355; US 7,968,834

1 properties



intermediate band



SiOnyx













Resolution	pixel (µm)	noise (e/pix)	/ _{dark} (e/pix/s)	<i>P</i> (mW)
872 x 654	5.6	2.1	24	300
1280 x 720	5.6	2.1	24	360
1280 x 1024	10	2.6	83	400
	8" CIS process flow 4T pixel architecture			

SiOnyx





















2 intermediate band

properties





2 intermediate band

properties




SiOnyx







can turn b:Si absorption into carrier generation









properties

1

- can turn b:Si absorption into carrier generation
- very high responsivity in VIS and NIR





- can turn b:Si absorption into carrier generation
- very high responsivity in VIS and NIR
- disruptive improvement in Si imaging



properties



- can turn b:Si absorption into carrier generation
- very high responsivity in VIS and NIR
- disruptive improvement in Si imaging
- potential benefits in solar energy harvesting









Summary

- new doping process
- new class of material
- new types of devices







What is different about this process?







Compare femtosecond laser doping to:

- inclusion during growth
- thermal diffusion
- ion implantation









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